

Allen-Cahn Solutions with Triple Junction Structure at Infinity

Étienne Sandier¹, Peter Sternberg²

¹) Université Paris-Est Créteil
sandier@u-pec.fr

²) Department of Mathematics, Indiana University, Bloomington
sternber@iu.edu

Abstract

We construct an entire solution $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ to the elliptic system

$$\Delta U = \nabla_u W(U),$$

where $W : \mathbb{R}^2 \rightarrow [0, \infty)$ is a ‘triple-well’ potential. This solution is a local minimizer of the associated energy

$$\int \frac{1}{2} |\nabla U|^2 + W(U) dx$$

in the sense that U minimizes the energy on any compact set among competitors agreeing with U outside that set. Furthermore, we show that along subsequences, the ‘blowdowns’ of U given by $U_R(x) := U(Rx)$ approach a minimal triple junction as $R \rightarrow \infty$. Previous results had assumed various levels of symmetry for the potential and had not established local minimality, but here we make no such symmetry assumptions.

1 Introduction

We will construct an entire solution $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ to the system

$$\Delta U = \nabla_u W(U), \tag{1.1}$$

which is minimizing on compact sets with respect to the associated energy

$$E(u) = \int \frac{1}{2} |\nabla u|^2 + W(u) dx,$$

where $W : \mathbb{R}^2 \rightarrow [0, \infty)$ is a C^2 ‘triple-well’ potential. That is, we assume that

$$\{p \in \mathbb{R}^2 : W(p) = 0\} = P := \{p_1, p_2, p_3\},$$

and we assume non-degeneracy of the potential wells in the sense that

$$D^2W(p_\ell) \geq bI \quad \text{for } \ell = 1, 2, 3 \text{ for some } b > 0, \text{ where } I \text{ is the } 2 \times 2 \text{ identity matrix.} \quad (1.2)$$

Additionally, we assume that for some $M > 0$,

$$p \cdot \nabla W(p) \geq 0 \quad \text{for } |p| \geq M. \quad (1.3)$$

As in many studies of vector Allen-Cahn, we will make extensive use of the following degenerate Riemannian metric on \mathbb{R}^2 :

$$d(p, q) := \inf \left\{ \sqrt{2} \int_0^1 W^{1/2}(\gamma(t)) |\gamma'(t)| dt : \gamma \in C^1([0, 1], \mathbb{R}^2), \gamma(0) = p, \gamma(1) = q \right\}, \quad (1.4)$$

and we denote by $c_{ij} := d(p_i, p_j)$ for $i \neq j$. We will assume that the strict triangle inequality holds between the wells p_1, p_2 and p_3 :

$$c_{12} < c_{13} + c_{23}, \quad c_{13} < c_{12} + c_{23} \quad \text{and} \quad c_{23} < c_{13} + c_{12}. \quad (1.5)$$

Under these assumptions, for $1 \leq i < j \leq 3$ there exists at least one length-minimizing geodesic ζ_{ij} joining p_i to p_j , see e.g. [3, 21, 29]. We will make the generic assumption that there is a *unique* such geodesic for each $i, j \in 1, 2, 3$, $i \neq j$, though perhaps this can be relaxed.

We also note that an equivalent variational description of the c_{ij} 's is given by

$$c_{ij} = \inf \left\{ \int_{-\infty}^{\infty} W(f(t)) + \frac{1}{2} |f'(t)|^2 dt : f \in H_{loc}^1(\mathbb{R}, \mathbb{R}^2), f(-\infty) = p_i, f(\infty) = p_j \right\}. \quad (1.6)$$

Under an appropriate parametrization, we then find that each $\zeta_{ij} : \mathbb{R} \rightarrow \mathbb{R}^2$ satisfies the system

$$\zeta_{ij}''(t) = \nabla_u W(\zeta_{ij}(t)) \quad \text{for } -\infty < t < \infty, \zeta_{ij}(-\infty) = p_i, \zeta_{ij}(\infty) = p_j. \quad (1.7)$$

From the perspective of ODE's, these geodesics ζ_{ij} represent heteroclinic connections between the potential wells.

We now denote by \mathcal{A} the set of all functions $u^* : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ taking the form

$$u^*(x) = \begin{cases} p_1 & \text{on } S_1 \\ p_2 & \text{on } S_2 \\ p_3 & \text{on } S_3, \end{cases} \quad (1.8)$$

where for $\ell = 1, 2$ and 3 , S_ℓ is a single (infinite) sector emanating from the origin with the three opening angles α_ℓ given by

$$\frac{\sin(\alpha_1)}{c_{23}} = \frac{\sin(\alpha_2)}{c_{13}} = \frac{\sin(\alpha_3)}{c_{12}}. \quad (1.9)$$

See Figure 1.

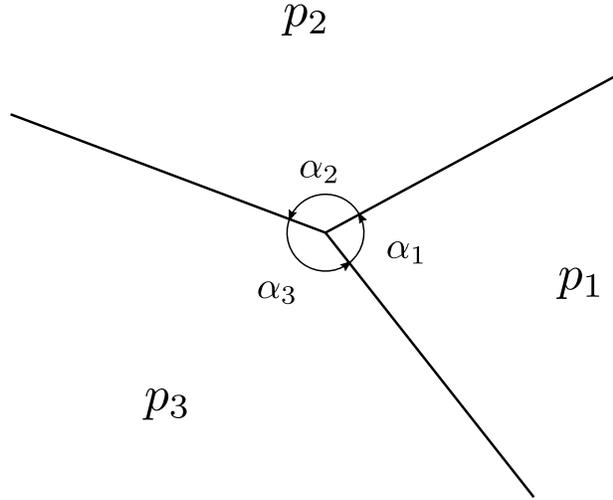


Figure 1: A locally minimizing partition of \mathbb{R}^2 with a triple junction.

The partition $\{S_1, S_2, S_3\}$ represents a locally minimizing partition of \mathbb{R}^2 with respect to the weighted perimeter functional

$$\{S_1, S_2, S_3\} \mapsto \sum_{1 \leq i < j \leq 3} c_{ij} \mathcal{H}^1(\partial S_i \cap \partial S_j), \quad (1.10)$$

where \mathcal{H}^1 refers to one-dimensional Hausdorff measure, and the condition (1.9) naturally arises as a criticality condition. As we will recall in Section 2, this partitioning problem represents the Γ -limit of a scaled version of the energy E , namely $E_R(u, \Omega)$, defined for any planar domain Ω , any $R > 0$ and any $u \in H^1(\Omega; \mathbb{R}^2)$, via

$$E_R(u, \Omega) = \int_{\Omega} RW(u) + \frac{1}{2R} |\nabla u|^2 \, dx. \quad (1.11)$$

We will write simply $E(u, \Omega)$ when referring to E_1 (i.e. $R = 1$).

We will establish a connection between the structure at infinity of our entire solution U to (1.1) and the triple junction partitions given by (1.8) by studying the asymptotic behavior of the blowdowns of U .

Our main result is the following:

Theorem 1.1. *There exists an entire solution $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ to*

$$\Delta U = \nabla_u W(U) \quad (1.12)$$

which is a local minimizer of energy in the sense that for every compact set $K \subset \mathbb{R}^2$ and for every $v \in H_{loc}^1(\mathbb{R}^2; \mathbb{R}^2)$ satisfying $v = U$ on $\mathbb{R}^2 \setminus K$ one has

$$E(U, K) \leq E(v, K). \quad (1.13)$$

Furthermore, defining U_R as the blowdown of U via

$$U_R(x) := U(Rx), \quad (1.14)$$

we have that on any compact set $K \subset \mathbb{R}^2$:

$$\text{dist}_{L^1(K; \mathbb{R}^2)}(U_R, \mathcal{A}) \rightarrow 0 \text{ as } R \rightarrow \infty. \quad (1.15)$$

That is,

$$\lim_{R \rightarrow \infty} \left(\inf_{u^* \in \mathcal{A}} \|U_R - u^*\|_{L^1(K; \mathbb{R}^2)} \right) = 0.$$

Remark 1.2. We believe that a stronger conclusion holds, namely that there exists a $u^* \in \mathcal{A}$ such that

$$\lim_{R \rightarrow \infty} \|U_R - u^*\|_{L^1(K; \mathbb{R}^2)} = 0.$$

A step in the proof of the above theorem is the following result, of independent interest.

Theorem 1.3. Assume $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is an entire solution to (1.12) which is a local minimizer of energy such that for some sequence $R_j \rightarrow +\infty$, the sequence U_{R_j} converges locally in L^1 to the function

$$u_0(x_1, x_2) = \begin{cases} p_i & \text{if } x_2 < 0. \\ p_j & \text{if } x_2 > 0. \end{cases}, \quad (1.16)$$

for some pair $p_i \neq p_j$. Then $U(x_1, x_2) = \zeta_{ij}(x_2 + \Delta)$, for some $\Delta \in \mathbb{R}$.

To place these results in context, we note that there is a large, and growing, collection of work on the general topic of finding entire solutions $u : \mathbb{R}^n \rightarrow \mathbb{R}^m$ to the vector Allen-Cahn system under various assumptions on the potential $W : \mathbb{R}^m \rightarrow \mathbb{R}$, on n and on m . See, for example [1, 4, 9, 10, 13, 18]. A source for a number of these results is the book [5]. Most of these results, however, include some form of symmetry assumption on W . We also mention the recent work [11] addressing concentration of general vector-valued critical points of Allen-Cahn in the plane.

Regarding the case under consideration here, namely $n = m = 2$ and W a triple well potential, an important first result on entire solutions appears in [13],

where the authors assume the potential is equivariant by the symmetry group of the equilateral triangle. The convergence to the minimal triple junction partition (1.8)-(1.9) they achieve under these symmetry assumptions (with necessarily each $\alpha_\ell = 2\pi/3$) is much stronger than (1.15). In particular, they show that

$$\lim_{t \rightarrow \infty} U \left(t \frac{x}{|x|} \right) = p_\ell \text{ for } x \in S_\ell \text{ off of the three rays } \partial S_k \cap \partial S_\ell \text{ for } 1 \leq k < \ell \leq 3.$$

On the other hand, since they work within the class of equivariant competitors, there is no claim of stability with respect to general perturbations. In a more recent contribution to this problem [15], the symmetry assumptions on W are weakened to include only the rotation subgroup of the full symmetry group of the equilateral triangle, thus relaxing the assumption of reflectional symmetry.

Our proof of Theorem 1.1 proceeds by first appealing to [28] to construct a sequence of L^1 -local minimizers of E_R on a particular nonconvex bounded domain, cf. Theorem 2.2. The candidate for our entire solution arises through a blow-up of this sequence, but care is needed here to execute the blow-up about a point where the local minimizers take a value far from the three heteroclinics. This analysis is carried out in Section 2, culminating in Proposition 2.4, where the blow-up limit U is shown to be an entire, locally minimizing solution to (1.1) that avoids the three heteroclinics at the origin.

The next step involves an analysis of the blowdowns of any local minimizer U in the sense of (1.13). Here we invoke the machinery of Γ -convergence, including an identification of the Γ -limit for vector Allen-Cahn subject to a Dirichlet condition carried out recently in [16]. We argue in Proposition 3.1 that, up to passing to subsequences, these blowdowns converge to an L^1 -local minimizer u_0 of the Γ -limit given in (2.5), which takes the form of a partitioning problem .

The crucial estimate in our blowdown analysis comes in the form of an asymptotic equipartition of energy of any local minimizer, namely

$$\int_{B_R} \left(W(U) + \frac{1}{2} |\nabla U|^2 - \sqrt{2} \sqrt{W(U)} |\nabla U| \right) dx < C_2 R^{1-\alpha} \quad \text{for } R \gg 1, \quad (1.17)$$

where B_R is the disc of radius R centered at the origin, and $C_2 > 0$ and $\alpha \in (0, 1)$ are constants independent of R . This is established in Proposition 3.3. The proof utilizes the regularity theory for the partitioning of a ball into three sets subject to a Dirichlet condition to obtain an upper bound on the energy, as well as a comparison between the infimum of such a partitioning problem and a related, less standard, partitioning problem described below (1.19) to obtain a matching lower bound. We appeal to the regularity theory for both problems as recently presented in [24]. We note that in [11] there appear other results

on the asymptotic behavior of the ‘discrepancy measure,’ that is, the integrand of (1.17), but these have a different nature given that they are derived only for critical points, not local minimizers, of the energy E .

From (1.17) and a Pohozaev identity, we are able to establish the convergence of the blowdowns to a minimal cone via Lemmas 3.4 and 3.5 and Theorem 3.6. It is then simple to conclude that one of three limits must arise: either (i) the minimal cone is \mathbb{R}^2 , i.e. $u_0 = p_\ell$ for some $p_\ell \in P$, (ii) the minimal cone is a half-space, i.e. $u_0 = p_i$ and p_j for $i \neq j$ on either side of a line, or (iii) the minimal cone is given by three sectors satisfying (1.9) so that u_0 is given by (1.8), cf. Proposition 3.7.

Eliminating possibility (i) is easy, but eliminating (ii)—which roughly corresponds to arguing that at infinity, the entire solution U does not look like a heteroclinic—is much more delicate. This is the content of Section 4. The proof is by contradiction. We first obtain an upper bound for the energy that corresponds to the cost of a heteroclinic. Then we obtain a contradictory lower bound using crucially that U was constructed in such a way that $U(0)$ is far from the three heteroclinics.

In Section 5, we compare the two partitioning problems that emerge in our proof of (1.17). As one is somewhat non-standard, we hope this section will be of independent interest. The first problem involves the minimization of the partitioning functional

$$(S_1, S_2, S_3) \mapsto t_1 \mathcal{H}^1(\partial S_1 \cap B) + t_2 \mathcal{H}^1(\partial S_2 \cap B) + t_3 \mathcal{H}^1(\partial S_3 \cap B), \quad (1.18)$$

where $B \subset \mathbb{R}^2$ is a ball, t_1, t_2 and t_3 are positive numbers, and the admissible competitors (S_1, S_2, S_3) are all partitions of B satisfying a Dirichlet condition

$$\partial S_\ell \cap \partial B = f^{-1}(p_\ell) \quad \text{for } \ell = 1, 2, 3 \quad \text{and } f \in BV(\partial B; P). \quad (1.19)$$

(In the present context of Allen-Cahn, the coefficients t_ℓ are related to the constants c_{ij} via (3.18), making (1.18) equivalent to (1.10).) For the second problem, one fixes any number $\delta > 0$ and then minimizes the same partitioning functional (1.18) among triples (S_1, S_2, S_3) of disjoint subsets of B again subject to (1.19), but now under the more relaxed condition that $|B \setminus \cup_\ell S_\ell| \leq \delta$. In other words, the competitors only need to “almost partition” the ball. In Theorem 5.4, we prove that the infimum of the second, more relaxed problem cannot lie more than $O(\delta^{1/2})$ below the infimum of the first problem.

In a personal communication in October of 2023, Nick Alikakos brought to our attention that he and Zhiyuan Geng were working on the same type of result. Their efforts eventually led to [6] and [7]. They obtain the same conclusion as that of our Theorem 1.1, along with information about the proximity of the

entire solution to the three potential wells along sequences of points going to infinity. The methods are quite different, with their result on convergence of blowdowns relying on a characterization of minimizing planar partitions into three sets, see Remark 3.8. As described above, our approach involves a new result on asymptotic equipartition of energy for local minimizers, along with the analysis of the rather novel geometry problem of “almost partitions.”

2 Construction of a candidate for the entire solution

Throughout this article, we will denote by $B_r(x)$ the ball in \mathbb{R}^2 of radius r and center x , unless the center is the origin, in which case we will simply write B_r .

2.1 Γ -convergence results

Our approach in this article will at times invoke Γ -convergence results relating the energy $E_R(u, \Omega)$ from (1.11) for a bounded domain $\Omega \subset \mathbb{R}^2$ to the functional

$$E_0(u, \Omega) := \sum_{1 \leq i < j \leq 3} c_{ij} \mathcal{H}^1(\partial^* S_i \cap \partial^* S_j \cap \Omega), \quad (2.1)$$

where $S_j := u^{-1}(p_j)$ for $j = 1, 2, 3$, and $\partial^* S$ refers to the reduced boundary of a set S of finite perimeter, cf. [17].

Building on previous Γ -convergence results for vector Modica-Mortola in the double-well case, e.g. [14, 26, 27], the Γ -convergence of $\{E_R(\cdot, \Omega)\}$ to $E_0(\cdot, \Omega)$ for bounded $\Omega \subset \mathbb{R}^n$ in the setting of a multi-well potential and in the topology $L^1(\Omega; \mathbb{R}^n)$ is established in [8].

We will also require a generalization of this Γ -convergence result to the situation where a Dirichlet condition is specified on $\partial\Omega$. Modica-Mortola type results that accommodate a Dirichlet condition appear in [25] in the scalar setting and in [23] in the context of the closely related Landau-deGennes energy. For our setting, however, we point to the recent result in [16]. For our purposes, it will suffice to state it for any bounded planar domain Ω with smooth boundary and for Dirichlet data taking values in the potential wells, though it holds more generally. To this end, let $h \in BV(\partial\Omega; P)$ and consider any sequence $\{h^R\} \subset H^1(\partial\Omega; \mathbb{R}^2)$ such that

$$|\partial_s h^R| \leq CR \text{ for some } C > 0 \quad \text{and} \quad h^R \rightarrow h \text{ in } L^1(\partial\Omega; \mathbb{R}^2) \text{ as } R \rightarrow \infty, \quad (2.2)$$

and such that

$$\int_{\partial\Omega} RW(h^R) + \frac{1}{2R} |\partial_s h^R|^2 \, d\mathcal{H}^1 < C \quad \text{for some constant } C \text{ independent of } R. \quad (2.3)$$

Next define

$$\tilde{E}_R(u, \Omega) := \begin{cases} E_R(u, \Omega) & \text{if } u \in H^1(\Omega; \mathbb{R}^2), u = h^R \text{ on } \partial\Omega, \\ +\infty & \text{otherwise,} \end{cases} \quad (2.4)$$

and define

$$E_0^h(u, \Omega) := \begin{cases} \sqrt{2} E_0(u, \Omega) + \sum_{j=1}^3 \int_{S_j \cap \partial\Omega} d(p_j, h(x)) d\mathcal{H}^1 & \text{if } u \in BV(\Omega; P), \\ +\infty & \text{otherwise,} \end{cases} \quad (2.5)$$

where E_0 is defined in (2.1) and $d(\cdot, \cdot)$ is given by (1.4). Then we have

Theorem 2.1. ([16]) *Assume $\{h^R\}$ satisfies (2.2) and (2.3). Then, as $R \rightarrow \infty$, the sequence $\{\tilde{E}_R(\cdot, \Omega)\}$ has the $L^1 - \Gamma$ -limit $E_0^h(\cdot, \Omega)$. That is, for every $u \in L^1(\Omega; \mathbb{R}^2)$ we have the following two conditions:*

(i) (Lower-semi-continuity) *If $\{v_R\} \subset L^1(\Omega; \mathbb{R}^2)$ is any sequence converging to u in L^1 then*

$$\liminf_{R \rightarrow \infty} \tilde{E}_R(v_R, \Omega) \geq E_0^h(u), \quad (2.6)$$

and

(ii) (Recovery sequence) *There exists a sequence $\{V_R\} \subset L^1(\Omega; \mathbb{R}^2)$ converging to u in L^1 such that*

$$\lim_{R \rightarrow \infty} \tilde{E}_R(V_R, \Omega) = E_0^h(u). \quad (2.7)$$

2.2 Construction of the entire solution via blow-up

Our candidate for an entire solution satisfying Theorem 1.1 will be constructed through a blow-up process, starting from an L^1 -local minimizer of $E_R(\cdot, \Omega)$ for a particular choice of Ω . This local minimizer is, in turn, constructed in [28] using Γ -convergence techniques.

To place ourselves in the setting of [28], we fix any $u^* \in \mathcal{A}$ given by (1.8) and let x_1, x_2 and x_3 be the three points on ∂B_1 where the three phase boundaries hit the unit circle. Then let $\Omega \subset \mathbb{R}^2$ be any bounded, simply connected open set containing B_1 such that $\partial\Omega$ is smooth and $\partial\Omega \cap \partial B_1 = \{x_1, x_2, x_3\}$. Finally, assume that $\partial\Omega$ is strictly concave at these three points. See Figure 2

Under these assumptions on Ω , the following theorem is proven in [28], utilizing the local minimizer property associated with Γ -convergence established in [19].

Theorem 2.2. For $\Omega \subset \mathbb{R}^2$ as described above there exists a number $\delta_0 > 0$ such that for all R sufficiently large, there exists an L^1 -local minimizer u_R of $E_R(\cdot, \Omega)$ in the sense that

$$E_R(u_R, \Omega) \leq E_R(v, \Omega) \quad \text{provided } \|v - u_R\|_{L^1(\Omega)} \leq \delta_0. \quad (2.8)$$

Furthermore,

$$u_R \rightarrow u^* \text{ in } L^1(\Omega) \quad (2.9)$$

and

$$E_R(u_R, \Omega) \rightarrow E_0(u^*, \Omega). \quad (2.10)$$

Necessarily, such a local minimizer satisfies the Euler-Lagrange equation associated with E_R , namely

$$\frac{1}{R^2} \Delta u_R = \nabla_u W(u_R) \quad \text{in } \Omega, \quad (2.11)$$

along with homogeneous Neumann boundary conditions on $\partial\Omega$.

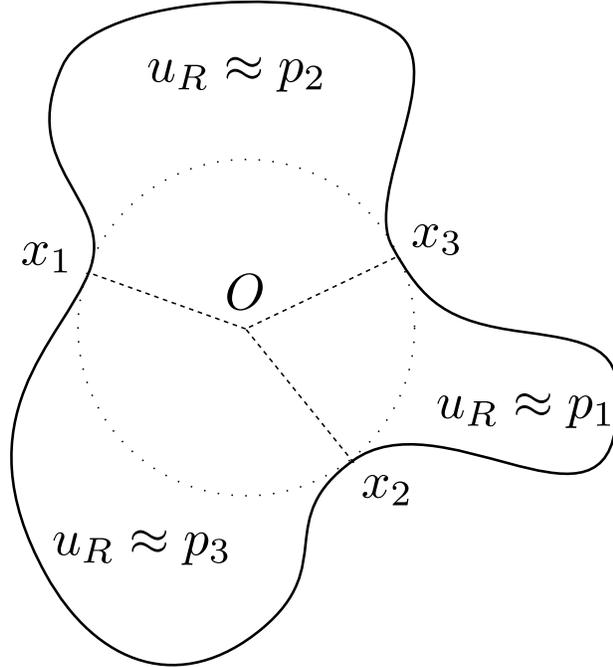


Figure 2: The L^1 -local minimizer u_R .

Referring back to the three geodesics ζ_{ij} defined below (1.5), we note that each is a simple curve (i.e. no self-intersections) and furthermore, any two of

them, after including their endpoints, only intersect at one of their endpoints, e.g. ζ_{12} and ζ_{13} only intersect at p_1 . This is because any transversal crossing would necessarily create a non- C^1 geodesic, violating regularity theory and any tangential intersection would violate the uniqueness of solutions to (1.7) subject to given initial conditions. As such, if we define Λ as the union of the closure of the images of these three geodesics, that is,

$$\Lambda := P \cup \zeta_{12}(\mathbb{R}) \cup \zeta_{23}(\mathbb{R}) \cup \zeta_{13}(\mathbb{R}), \quad (2.12)$$

then we can identify Λ as a simple, closed curve in \mathbb{R}^2 passing through p_1, p_2 and p_3 which is smooth except at these 3 points. See Figure 3.

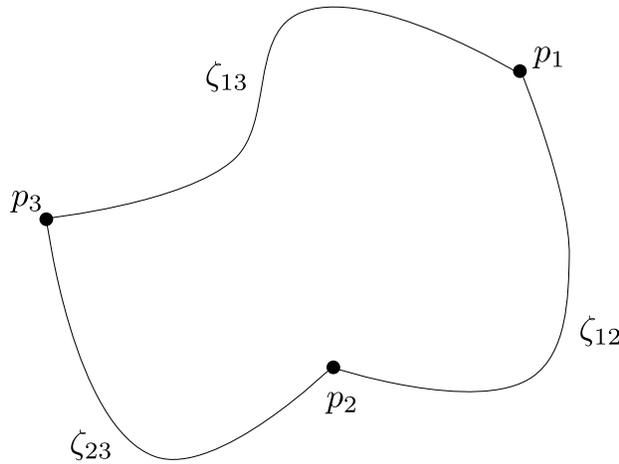


Figure 3: The closed curve Λ consisting of the three heteroclinics.

An important property of the local minimizers constructed in Theorem 2.2 is the following.

Lemma 2.3. *Let $\{u_R\}$ be the sequence of L^1 local minimizers established in Theorem 2.2. Then there exists a ball B' compactly contained in B_1 , a point p inside Λ and a sequence of points $\{x_R\} \subset B'$ such that $u_R(x_R) = p$. In particular, there is a value $a_0 > 0$ such that*

$$\text{dist}(u_R(x_R), \Lambda) > a_0.$$

Proof. We have that $\|u_R - u^*\|_{L^1(B_1)}$ tends to zero as $R \rightarrow +\infty$. Hence, by Fatou's Lemma,

$$0 \geq \int_0^1 \liminf_{R \rightarrow +\infty} \|u_R - u^*\|_{L^1(\partial B_r)} dr \geq 0. \quad (2.13)$$

Similarly, for almost every $r \in (0, 1)$, it holds that

$$\liminf_{R \rightarrow +\infty} E_R(u_R, \partial B_r) < +\infty. \quad (2.14)$$

It follows that there exists $s \in (1/4, 1/3)$ and $t \in (1/2, 2/3)$, and a subsequence still denoted $\{u_R\}$, such that $u_R \rightarrow u^*$ in $L^1(\partial \mathcal{A})$, where $\mathcal{A} = B_t \setminus B_s$ and (2.14) holds for $r = s, t$.

Since it also follows from standard elliptic estimates that $|\nabla u_R| \leq CR$, we have all the hypotheses of Theorem 2.1 satisfied on $\Omega = \mathcal{A}$ and so we can assert the existence of a recovery sequence, say $\{\tilde{u}_R\}$, associated with u^* and the boundary values of u_R on $\partial \mathcal{A}$. It then follows from the L^1 -local minimality of u_R (2.8) that

$$\begin{aligned} (t-s)(c_{12} + c_{23} + c_{13}) &\geq \lim_{R \rightarrow +\infty} \int_s^t E_R(\tilde{u}_R, \partial B_r) dr \\ &\geq \lim_{R \rightarrow +\infty} \int_s^t E_R(u_R, \partial B_r) dr \geq (t-s)(c_{12} + c_{23} + c_{13}). \end{aligned}$$

Using (2.13) again, we deduce the existence of some $r \in (1/4, 2/3)$ such that $u_R \rightarrow u^*$ in $L^1(\partial B_r)$ and

$$E_R(u_R, \partial B_r) \rightarrow c_{12} + c_{23} + c_{13}. \quad (2.15)$$

From the convergence of $\{u_R\}$ in $L^1(\partial B_r)$, we deduce the existence of three angles θ_i , $i = 1, 2, 3$, such that going to a further subsequence $u_R(re^{i\theta_i})$ converges to the well p_i for each i . It then follows from (2.15) that u_R is a minimizing sequence for the one dimensional energy E_R on each of the arcs A_{ij} , where A_{ij} is the portion of ∂B_r between the angles θ_i and θ_j .

In light of our assumption of uniqueness for the three heteroclinic connections, we may assert that for R sufficiently large one has

$$u_R(A_{ij}) \text{ is uniformly close to } \zeta_{ij}(\mathbb{R}) \text{ for } 1 \leq i < j \leq 3.$$

Consequently, the closed curve $u_R(\partial B_r)$ is uniformly close to the simple, closed curve Λ .

However, since Λ is a Jordan curve, it partitions \mathbb{R}^2 into an inside, say U , and an outside unbounded set. For any $p \in U$, the index of Λ with respect to p is equal to ± 1 and the same must be true for the curve $u_R(\partial B_r)$ since it is uniformly close to Λ for R large enough. Therefore, the latter curve cannot be homotopic to a constant in $\mathbb{R}^2 \setminus \{p\}$, and so $p \in u_R(B(0, r))$ for any R large enough. Selecting any $p \in U$ and any $x_R \in B(0, r)$ such that $u_R(x_R) = p$, the result follows. □

We now introduce our candidate for the entire solution of Theorem 1.1 by taking a limit of blow-ups of $\{u_R\}$.

Proposition 2.4. *Let $\{u_R\}$ be the sequence of local minimizers established in Theorem 2.2. Let $\Omega^R := \{x : \frac{x}{R} + x_R \in \Omega\}$ where $\{x_R\}$ is the sequence introduced in Lemma 2.3. Also define $V^R(x) : \Omega^R \rightarrow \mathbb{R}^2$ via $V^R(x) := u_R(\frac{x}{R} + x_R)$. Then there exists a subsequence $\{R_j\} \rightarrow \infty$ and a function $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ such that*

$$V^{R_j} \rightarrow U \text{ in } C^2 \text{ on compact subsets of } \mathbb{R}^2 \quad (2.16)$$

where U solves (1.1). Furthermore, U is a local minimizer of E in the sense of (1.13). Finally, we have

$$\text{dist}(U(0), \Lambda) > 0. \quad (2.17)$$

Remark 2.5. *For the remainder of the paper, when we say that a function U is a local minimizer of E , we will always mean in the sense of (1.13).*

Proof. Assumption (1.3) implies through the maximum principle applied to $|u_R|^2$ that $\|u_R\|_{L^\infty(\Omega)} \leq M$ and so the same is true of $\{V_R\}$. In light of (2.11) we observe that V^R satisfies (1.1) on Ω^R . Then standard elliptic estimates and bootstrapping leads, in particular, to uniform $C^{2,\alpha}$ bounds on compact sets for $\{V^R\}$. The conclusion (2.16) follows as does the assertion that U solves (1.1).

To establish the local minimality of U , fix any compact set K and let $\tilde{v} : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ be any smooth function supported in K . Let $\tilde{v}_R(x) := \tilde{v}(R(x - x_R))$ so that \tilde{v}_R is supported in $x_R + \frac{1}{R}K$. Then we have

$$\int_K |\tilde{v}_R| \, dx \leq \left(\max_K |\tilde{v}| \right) \frac{1}{R^2} |K| \quad \text{where } |K| = \text{Lebesgue measure of } K.$$

Now taking R large enough so that $(\max_K |\tilde{v}|) \frac{1}{R^2} |K| < \frac{\delta_0}{2}$ we can invoke (2.8) to conclude that

$$\begin{aligned} 0 &\leq E_R(u_R + \tilde{v}_R, \Omega) - E_R(u_R, \Omega) \\ &= E_R(u_R + \tilde{v}_R, x_R + \frac{1}{R}K) - E_R(u_R, x_R + \frac{1}{R}K) \\ &= \frac{1}{R} (E(V^R + \tilde{v}, K) - E(V^R, K)). \end{aligned}$$

Passing to the limit $R_j \rightarrow \infty$ in the inequality $E(V^{R_j}, K) \leq E(V^{R_j} + \tilde{v}, K)$ we obtain (1.13).

Property (2.17) follows from Lemma 2.3 in light of the uniform convergence of $V^{R_j} \rightarrow U$. \square

We conclude this section with a simple but crucial estimate on U given in the following:

Lemma 2.6. *There exists a constant $C_1 = C_1(W)$ such that for every $R > 0$ one has*

$$E(U, B_R) \leq C_1 R. \quad (2.18)$$

Proof. We may as well assume $R > 1$. Then we appeal to the local minimality of U , namely (1.13), with v chosen to equal, say, p_1 on B_{R-1} and then v smoothly interpolating between p_1 and U on the annulus $B_R \setminus B_{R-1}$. Since $E(v, B_{R-1}) = 0$ and U and ∇U are uniformly bounded in terms of W on the annulus, the result follows. \square

3 Blowdown Analysis

In this section, we will characterize the limits of the blowdowns of any local minimizer of E . For this purpose we will need the following compactness result associated with local minimizers of E .

Proposition 3.1. *Let $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ be a local minimizer of E . We have:*

(i) *Let $\{R_j\} \rightarrow \infty$ be any sequence. Then there exists a subsequence $\{R_{j_k}\}$ and a function $u_0 \in BV_{\text{loc}}(\mathbb{R}^2; P)$ such that the blowdowns $\{U_{R_{j_k}}\}$ of U satisfy*

$$U_{R_{j_k}} \rightarrow u_0 \quad \text{in } L^1_{\text{loc}}(\mathbb{R}^2; \mathbb{R}^2). \quad (3.1)$$

(ii) *After perhaps passing to a further subsequence (still denoted by $\{R_{j_k}\}$), one has for every $\ell \in \mathbb{Z}^+$ there exists a radius $\lambda_\ell \in [\ell, \ell + 1]$ such that*

$$\sup_k \left(\int_{\partial B_{\lambda_\ell}} R_{j_k} W(U_{R_{j_k}}) + \frac{1}{2R_{j_k}} |\nabla U_{R_{j_k}}|^2 d\mathcal{H}^1 \right) \leq 3C_1, \quad (3.2)$$

where C_1 is the constant appearing in Lemma 2.6.

Furthermore,

$$U_{R_{j_k}} \rightarrow \text{trace of } u_0 \quad \text{in } L^1(\partial B_{\lambda_\ell}; \mathbb{R}^2). \quad (3.3)$$

Lastly, u_0 is a local minimizer of $E_0(\cdot, \mathbb{R}^2)$ given in (2.1) in the sense that

$$E_0(u_0, K) \leq E_0(v, K) \quad (3.4)$$

for every compact $K \subset \mathbb{R}^2$ and every $v \in BV_{\text{loc}}(\mathbb{R}^2; P)$ such that $v = u_0$ on $\mathbb{R}^2 \setminus K$.

Proof. Since $E(U, B_{\lambda R}) = R E_R(U_R, B_\lambda)$, it then follows from Lemma 2.6 that

$$E_R(U_R, B_\lambda) \leq \lambda C_1 \quad \text{for any } \lambda > 0. \quad (3.5)$$

Hence, the sequence $\{U_R\}$ has uniformly bounded energy on any ball B_λ and so the proof of (i) follows from [8], Prop. 4.1 using a diagonalization procedure. To prove (ii), we note that from (3.5), in particular, it follows that

$$\int_1^2 \int_{\partial B_\lambda} R_{j_k} W(U_{R_{j_k}}) + \frac{1}{2R_{j_k}} \left| \nabla U_{R_{j_k}} \right|^2 d\mathcal{H}^1 d\lambda < 2C_1 \quad \text{for all } k. \quad (3.6)$$

Letting

$$f_k(\lambda) := \int_{\partial B_\lambda} R_{j_k} W(U_{R_{j_k}}) + \frac{1}{2R_{j_k}} \left| \nabla U_{R_{j_k}} \right|^2 d\mathcal{H}^1,$$

let us suppose that (3.2) is false for $\ell = 1$. Then necessarily, for every $\lambda \in [1, 2]$ it would hold that

$$\liminf_{k \rightarrow \infty} f_k(\lambda) \geq 3C_1. \quad (3.7)$$

Then, by Fatou's Lemma,

$$2C_1 \geq \liminf_{k \rightarrow \infty} \int_1^2 f_k(\lambda) d\lambda \geq \int_1^2 \liminf_{k \rightarrow \infty} f_k(\lambda) d\lambda \geq 3C_1,$$

and a contradiction is reached. Passing to a further subsequence, the existence of a function $h \in BV(\partial B_{\lambda_1}; P)$ such that $U_{R_{j_k}} \rightarrow h$ in $L^1(\partial B_{\lambda_1}; \mathbb{R}^2)$ follows from (3.2) using the same compactness argument from [8], applied now to the energy restricted to the circle ∂B_{λ_1} with the full gradient replaced by the tangential gradient.

To establish (3.3) and the local minimality of u_0 we observe from the Γ -convergence result Theorem 2.1 that u_0 is the limit of minimizers of $\tilde{E}_{R_{j_k}}(\cdot, B_{\lambda_1})$. Hence, u_0 must necessarily minimize $E_0^h(\cdot, B_{\lambda_1})$. Indeed, for any $v \in BV(B_{\lambda_1}; P)$ one has

$$E_0^h(v, B_\lambda) = \lim_{k \rightarrow \infty} \tilde{E}_{R_{j_k}}(V_k, B_{\lambda_1}) \geq \liminf_{k \rightarrow \infty} \tilde{E}_{R_{j_k}}(U_{R_{j_k}}, B_{\lambda_1}) \geq E_0^h(u_0, B_{\lambda_1}),$$

where $\{V_k\}$ is the recovery sequence associated with v guaranteed to exist by Theorem 2.1. We note that the first inequality above follows from the local minimality of $U_{R_{j_k}}$ in the sense of (1.13), since by construction, $V_k = U_{R_{j_k}}$ on ∂B_{λ_1} .

It follows that in fact, $u_0 = h$ on ∂B_{λ_1} , since, for example, if $h = p_1$ along some arc $\gamma \subset \partial B_{\lambda_1}$, while the trace of $u_0 = p_2$ on γ , then one could produce a

lower energy competitor v for the energy $E_0^h(\cdot, B_{\lambda_1})$ by setting $v = p_1$ inside the slice of B_{λ_1} bounded by γ and the secant line L connecting the endpoints of γ . Then

$$E_0^h(u_0, B_{\lambda_1}) - E_0^h(v, B_{\lambda_1}) = d(p_1, p_2) (\mathcal{H}^1(\gamma) - \mathcal{H}^1(L)) > 0,$$

in light of the strict convexity of B_{λ_1} , thus contradicting the minimality of u_0 . Necessarily then, u_0 is also a minimizer of E_0 among competitors agreeing with u_0 on ∂B_{λ_1} .

We conclude the proof by noting that the same logic allows us to select a value $\lambda_\ell \in [\ell, \ell + 1]$ for every $\ell \in \mathbb{Z}^+$ and thus to conclude that u_0 minimizes E_0 in every ball B_{λ_ℓ} among all competitors that agree with u_0 on $\partial B_{\lambda_\ell}$. Hence, (3.4) holds. \square

3.1 Pohozaev and asymptotic equipartition of energy

With an eye towards utilizing a Pohozaev identity, we next introduce the stress-energy tensor associated with a solution $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ to (1.1):

$$T_{ij} = U_{x_i} U_{x_j} - \delta_{ij} \left(\frac{1}{2} |\nabla U|^2 + W(U) \right).$$

A standard calculation yields that T is divergence-free. From this fact we get the Pohozaev identity on the ball B_R :

$$\int_{B_R} (x_i T_{ij})_{x_j} dx = \int_{B_R} \delta_{ij} T_{ij} + x_i (T_{ij})_{x_j} dx = \int_{B_R} \text{tr } T dx.$$

Applying the divergence theorem leads to

$$R \int_{\partial B_R} \nu_i T_{ij} \nu_j d\mathcal{H}^1 = -2 \int_{B_R} W(U) dx,$$

where $\nu = x/R$ is the outer unit normal to B_R . Using the definition of T this can be written as

$$\frac{1}{2} \int_{\partial B_R} \frac{1}{2} |U_\nu|^2 - \frac{1}{2} |U_s|^2 - W(U) d\mathcal{H}^1 = -\frac{1}{R} \int_{B_R} W(U) dx, \quad (3.8)$$

where U_s denotes the tangential derivative of U along ∂B_R . Through (3.8) we immediately obtain the following identity.

Proposition 3.2. *Any entire solution U to (1.1) satisfies*

$$\frac{d}{dR} \left(\frac{1}{R} \int_{B_R} W(U) dx \right) = \frac{1}{2R} \int_{\partial B_R} \frac{1}{2} |U_\nu|^2 - \frac{1}{2} |U_s|^2 + W(U) d\mathcal{H}^1 \quad (3.9)$$

for all $R > 0$.

Our aim is to obtain a kind of asymptotic monotonicity result. To this end, we define

$$\tilde{W}(R) := \frac{1}{R} \int_{B_R} W(U) dx. \quad (3.10)$$

Then for any two values $0 < R_1 < R_2$ we integrate (3.9) to find that

$$\begin{aligned} \tilde{W}(R_2) - \tilde{W}(R_1) &= \int_{R_1}^{R_2} \frac{1}{2r} \int_{\partial B_r} \frac{1}{2} |U_\nu|^2 - \frac{1}{2} |U_s|^2 + W(U) d\mathcal{H}^1 dr \\ &= \int_{B_{R_2} \setminus B_{R_1}} \frac{1}{2|x|} \left(\frac{1}{2} |U_\nu|^2 - \frac{1}{2} |U_s|^2 + W(U) \right) dx \\ &= \int_{B_{R_2} \setminus B_{R_1}} \frac{1}{2|x|} \left(W(U) - \frac{1}{2} |\nabla U|^2 + |U_\nu|^2 \right) dx. \end{aligned} \quad (3.11)$$

Hence,

$$\begin{aligned} \tilde{W}(R_2) - \tilde{W}(R_1) &\geq - \int_{B_{R_2} \setminus B_{R_1}} \frac{1}{2|x|} \left| W(U) - \frac{1}{2} |\nabla U|^2 \right| dx \\ &\geq - \frac{1}{2R_1} \int_{B_{R_2} \setminus B_{R_1}} \left(\sqrt{W(U)} - \frac{1}{\sqrt{2}} |\nabla U| \right) \left(\sqrt{W(U)} + \frac{1}{\sqrt{2}} |\nabla U| \right) dx \\ &\geq - \frac{1}{2R_1} \left\{ \int_{B_{R_2} \setminus B_{R_1}} \left(\sqrt{W(U)} - \frac{1}{\sqrt{2}} |\nabla U| \right)^2 dx \right\}^{1/2} \left\{ \int_{B_{R_2} \setminus B_{R_1}} \left(\sqrt{W(U)} + \frac{1}{\sqrt{2}} |\nabla U| \right)^2 dx \right\}^{1/2}. \end{aligned} \quad (3.12)$$

Now in light of (2.18), we have that

$$\begin{aligned} \left\{ \int_{B_{R_2} \setminus B_{R_1}} \left(\sqrt{W(U)} + \frac{1}{\sqrt{2}} |\nabla U| \right)^2 dx \right\}^{1/2} &\leq \sqrt{2} \left\{ \int_{B_{R_2} \setminus B_{R_1}} \left(W(U) + \frac{1}{2} |\nabla U|^2 \right) dx \right\}^{1/2} \\ &\leq \sqrt{2C_1} R_2^{1/2}, \end{aligned} \quad (3.13)$$

so that (3.12) implies

$$\tilde{W}(R_2) - \tilde{W}(R_1) \geq - \sqrt{\frac{C_1}{2}} \frac{R_2^{1/2}}{R_1} \left\{ \int_{B_{R_2} \setminus B_{R_1}} \left(\sqrt{W(U)} - \frac{1}{\sqrt{2}} |\nabla U| \right)^2 dx \right\}^{1/2}. \quad (3.14)$$

Inequality (3.14) shows that we can achieve an asymptotic monotonicity-type formula provided we can establish a sufficiently sharp measure of equipartition of energy.

The key estimate we will show is:

Proposition 3.3. *There exist constants $C_2 > 0$ and $\alpha \in (0, 1)$, such that for any local minimizer U of E and any R sufficiently large one has the estimate*

$$\int_{B_R} \left(\sqrt{W(U)} - \frac{1}{\sqrt{2}} |\nabla U| \right)^2 dx < C_2 R^{1-\alpha}. \quad (3.15)$$

Our proof of Proposition 3.3 will involve the construction of a recovery sequence with a quantitative error bound, corresponding to a minimizer of $E_0(\cdot, B)$ in a ball B subject to a general Dirichlet condition $h \in BV(\partial B; P)$, that is, among partitions $\{S_1, S_2, S_3\}$ of B satisfying

$$\partial S_\ell \cap \partial B = h^{-1}(p_\ell) \quad \text{for } \ell = 1, 2, 3. \quad (3.16)$$

For this upper bound construction we will require a rather complete characterization of minimizers of this partitioning problem, to be found in Theorem 5.1 and Corollary 5.2. Our proof of Proposition 3.3 will also require a sharp lower bound for the energy of a related but somewhat non-standard partitioning problem. To state it, we first observe that given a partition, say $\{S_1, S_2, S_3\}$ of a ball B , its cost as given by E_0 can be equivalently expressed as

$$E_0(S_1, S_2, S_3) = t_1 \mathcal{H}^1(\partial S_1 \cap B) + t_2 \mathcal{H}^1(\partial S_2 \cap B) + t_3 \mathcal{H}^1(\partial S_3 \cap B), \quad (3.17)$$

where the numbers t_1, t_2 and t_3 are the solution to the system

$$t_1 + t_2 = c_{12}, \quad t_1 + t_3 = c_{13}, \quad t_2 + t_3 = c_{23}.$$

Solving, we find

$$t_1 = \frac{1}{2}(c_{12} + c_{13} - c_{23}), \quad t_2 = \frac{1}{2}(c_{12} + c_{23} - c_{13}), \quad t_3 = \frac{1}{2}(c_{13} + c_{23} - c_{12}), \quad (3.18)$$

and so we note that each t_j is positive in light of the assumption (1.5).

Then for any $\delta > 0$ and any $h \in BV(\partial B; P)$ we consider the minimization of E_0 as given by (3.17) among all disjoint subsets $\{S_1, S_2, S_3\}$ of B satisfying the Dirichlet condition (3.16), along with the constraint

$$|B \setminus (\cup_{\ell=1}^3 S_\ell)| \leq \delta. \quad (3.19)$$

We will require a good bound from below for the infimum of E_0 subject to (3.19) and Dirichlet data $h \in BV(\partial B; P)$ in terms of the infimum of E_0 subject to the same Dirichlet condition but for actual partitions of B , that is, with $\delta = 0$ in (3.19). This is presented in Theorem 5.4.

Proof of Prop. 3.3.

The proof of (3.15) will follow by first establishing an upper bound of the form

$$\int_{B_1} RW(U_R) + \frac{1}{2R} |\nabla U_R|^2 dx \leq m_R + \frac{C}{R^\alpha} \quad \text{for some } C > 0 \text{ and } \alpha \in (0, 1), \quad (3.20)$$

where m_R , defined below in (3.38), represents the minimal value of the partitioning problem E_0 subject to a certain Dirichlet condition related to U_R . Then we will utilize Theorem 5.4 to establish a matching lower bound of the form

$$\sqrt{2} \int_{B_1} \sqrt{W(U_R)} |\nabla U_R| dx \geq m_R - \frac{C'}{R^\alpha} \quad \text{for some } C' > 0. \quad (3.21)$$

If we rephrase the desired upper and lower bounds (3.20) and (3.21) in terms of U instead of its blowdowns, then the upper bound we seek takes the form

$$\int_{B_R} W(U) + \frac{1}{2} |\nabla U|^2 dx \leq Rm_R + CR^{1-\alpha}, \quad (3.22)$$

and the lower bound we want takes the form

$$\sqrt{2} \int_{B_R} \sqrt{W(U)} |\nabla U| dx \geq Rm_R - C'R^{1-\alpha}. \quad (3.23)$$

The desired inequality (3.15) then follows by combining (3.22) and (3.23).

Upper bound construction.

Since by Proposition 2.4, U_R minimizes $E_R(\cdot, B_1)$ among competitors sharing its boundary values on ∂B_1 , we can obtain the upper bound through a construction of a low-energy competitor. In essence, this is akin to the recovery sequence construction for vector Allen-Cahn with a multi-well potential, adapted to handle a Dirichlet condition, as in the recent work [16]. The difference is that here this must be made *quantitative* with an error that is $O(R^{-\alpha})$. However, unlike the general recovery sequence construction, here we only need to build it for an E_0 -minimizing partition that yields the value m_R in the problem (3.38) defined below.

To begin the pursuit of an upper bound, we first note that by (2.18) we have for any $R > 0$:

$$\int_R^{2R} E(U, \partial B_r) dr = E(U, B_{2R} \setminus B_R) \leq 2C_1 R.$$

Hence, by the Mean Value Theorem, there exists a value $R' \in (R, 2R)$ such that

$$E(U, \partial B_{R'}) \leq 2C_1. \quad (3.24)$$

If we can establish (3.15) for R' , then replacing C_2 by $2^{1-\alpha}C_2$, we will have established (3.15) for R as well. Thus, with no loss of generality, we may assume that R satisfies (3.24) as well. Phrasing this condition in terms of the blowdowns $\{U_R\}$, the assumed bound takes the form

$$E_R(U_R, \partial B_1) \leq 2C_1. \quad (3.25)$$

The upper bound estimate (3.20) will result from the construction of a low-energy competitor for the minimization of $E_R(\cdot, B_1)$ that agrees with the blown down minimizer U_R on ∂B_1 .

From (3.25), it follows that off of a small set on ∂B_1 , the function U_R must stay near one of the three wells p_1, p_2 or p_3 . We now use this fact to identify a partition of ∂B_1 into three sets.

We note that in light of the non-degeneracy assumption (1.2), there exists a positive number β , depending only on W , such that

$$W \text{ is strictly convex for } |p - p_j| < \beta, \quad j = 1, 2, 3. \quad (3.26)$$

and furthermore,

$$\frac{b}{2} |p_\ell - q|^2 \leq W(q) \leq 2b |p_\ell - q|^2. \quad (3.27)$$

Then let us define the set

$$A_R := \left\{ x \in \partial B_1 : d(U_R(x), P) > \frac{\beta}{2} \right\}. \quad (3.28)$$

This set is necessarily a union of open arcs. If such an arc I possesses a point x such that $d(U_R(x), P) \geq \beta$, then since at the endpoints of I , necessarily U_R is at metric distance $\frac{\beta}{2}$ from P , it must be the case that $E_R(U_R, I) \geq C$, for some positive constant C depending only on W . We define T_R to be the union of all such arcs, and so in light of (3.25), we can assert that T_R consists of a *finite* union of arcs whose total number is bounded by a constant depending only on W . It then follows from $\int_{T_R} RW(U_R) d\mathcal{H}^1 \leq 2C_1$ that

$$\mathcal{H}^1(T_R) \leq CR^{-1}. \quad (3.29)$$

On $\partial B_1 \setminus T_R$ we note that the metric distance from U_R to P is less than β .

Boundary layer construction on the annulus $B_1 \setminus B_{1-\rho}$

We begin with the construction of a boundary layer on $B_1 \setminus B_{1-\rho}$, where ρ will be determined later. However, we will insist that

$$\rho \geq \frac{1}{R}. \quad (3.30)$$

The number of disjoint arcs in T_R is bounded by a constant depending on W only, hence the same is true for the complement of T_R . We split this complement into two sets; S_R and the remainder, the set S_R being the union of arcs having length less than λ , where $\lambda \leq 1$ is another parameter to be determined later. Let us denote the arcs in the remainder by say $\{I_k\}$ for $k = 1, 2, \dots, N_R$, where N_R is bounded by a constant $N_0 = N_0(W)$. Then each I_k will be of length at least λ and can be naturally associated with one of the wells in the sense that U_R remains within a metric distance of β from that well throughout I_k . We now can expand each I_k to a slightly larger arc \tilde{I}_k , absorbing arcs of T_R and of S_R in the process, so as to form a partition of ∂B_1 , where in light of (3.29), we know that

$$\mathcal{H}^1(\partial B_1 \setminus \cup I_k) = \mathcal{H}^1(\cup \tilde{I}_k \setminus \cup I_k) \leq C(R^{-1} + \lambda). \quad (3.31)$$

We note that there is some ambiguity in terms of the assignment of an element of P to arcs comprising $\cup \tilde{I}_k \setminus \cup I_k$. That is, if say I_k is associated with p_1 and an adjacent arc I_{k+1} is associated with p_2 , then one can either expand I_k into the gap between them and assign the value p_1 to the resulting \tilde{I}_k or expand I_{k+1} into the gap and assign the value p_2 to the resulting \tilde{I}_{k+1} . As well shall see, due to the smallness of these gaps guaranteed by (3.31), it will not matter which choice we make here.

On $\partial B_{1-\rho}$ we define a function $V_R(1-\rho, \theta)$ as follows. If $e^{i\theta} \in \tilde{I}_k$ and $e^{i\theta}$ is at least a distance $\frac{1}{R}$ from the endpoints of \tilde{I}_k , we take $V_R(1-\rho, \theta)$ to be equal to whichever well is associated with I_k . On the rest of $\partial B_{1-\rho}$, we define $V_R(1-\rho, \theta)$ through linear interpolation in θ , so that

$$\left| \frac{\partial V}{\partial \theta}(1-\rho, \theta) \right| \leq CR. \quad (3.32)$$

Now we define V_R in the annulus $\mathcal{A}_{1,1-\rho} := B_1 \setminus B_{1-\rho}$ taking $V_R(r, \theta)$ to linearly interpolate in r between $U_R(1, \theta)$ and $V_R(1-\rho, \theta)$ for each θ . We estimate the energy in this annulus as follows:

We begin with the cost of interpolation from $U_R(1, \theta)$ to $V_R(1-\rho, \theta)$ for $e^{i\theta} \in \cup_{k=1}^{N_R} I_k$. In view of (3.26) and the fact that $V_R(1-\rho, \theta)$ is a constant equal to one of the wells on each I_k , we can invoke the convexity of all terms in the energy to assert that for any $t \in (0, 1)$ one has

$$W(V_R(1-t\rho, \theta)) \leq (1-t)W(U_R(1, \theta)) \text{ and } \left| \frac{\partial V_R}{\partial \theta}(1-t\rho, \theta) \right|^2 \leq (1-t) \left| \frac{\partial U_R}{\partial \theta}(1, \theta) \right|^2. \quad (3.33)$$

Estimating the radial derivative, we find

$$\left| \frac{\partial V_R}{\partial r}(r, \theta) \right|^2 \leq \frac{|U_R(1, \theta) - p_\ell|^2}{\rho^2} \leq C \frac{W(U_R(1, \theta))}{\rho^2} \text{ for some } \ell \in \{1, 2, 3\}, \quad (3.34)$$

for any $r \in (1 - \rho, 1)$. Combining (3.25), (3.33) and (3.34), we integrate over that part of the annulus $\mathcal{A}_{1,1-\rho}$ corresponding to the set of arcs $\cup_{k=1}^{N_R} I_k$, say $A'_{1,1-\rho}$, to obtain

$$E_R(V_R, A'_{1,1-\rho}) \leq C \left(\rho + \frac{1}{R^2 \rho} \right). \quad (3.35)$$

Now we turn to an estimate of the energetic cost in that portion of the annulus, say $A''_{1,1-\rho}$, corresponding to arcs in the complement of $\cup_{k=1}^{N_R} I_k$. Estimating the tangential derivative in $A''_{1,1-\rho}$, we find in view of (3.32) that

$$\left| \frac{\partial V_R}{\partial \theta}(r, \theta) \right|^2 \leq CR^2,$$

and for the normal derivative we have

$$\left| \frac{\partial V_R}{\partial r}(r, \theta) \right|^2 \leq C \frac{1}{\rho^2}.$$

Invoking (3.31), for the potential term we can then estimate that

$$\int_{A''_{1,1-\rho}} RW(V_R) dx \leq CR |A''_{1,1-\rho}| \leq CR \rho (R^{-1} + \lambda),$$

and so

$$\begin{aligned} E_R(V_R, A''_{1,1-\rho}) &\leq C \rho \left(\frac{1}{R} + \lambda \right) \left(R + \frac{1}{R \rho^2} \right) \\ &= C \left(\rho + \frac{1}{R^2 \rho} + \rho \lambda R + \frac{\lambda}{R \rho} \right). \end{aligned} \quad (3.36)$$

Therefore, summing (3.35) and (3.36), we see that, since $\lambda \leq 1$ and $R \geq 1$,

$$E_R(V_R, \mathcal{A}_{1,1-\rho}) \leq C \left(\rho + \frac{1}{R \rho} + \rho \lambda R \right). \quad (3.37)$$

Construction of the competitor in $B_{1-\rho}$

Let us now define V_R in the ball $B_{1-\rho}$. For this, we will introduce the minimization of E_0 subject to the Dirichlet condition $h_R : \partial B_{1-\rho} \rightarrow P$ satisfying $h_R(x) = p_\ell$ if $x/|x| \in \tilde{I}_k$, where I_k is the arc associated with p_ℓ and \tilde{I}_k is its expansion, as described above (3.31):

$$m_R := \inf \{ E_0(u, B_{1-\rho}) : u \in BV(B_{1-\rho}; P), u = h_R \text{ on } \partial B_{1-\rho} \}. \quad (3.38)$$

Let u_0 denote a minimizer for (3.38). By Theorem 5.1 we know that $\{u_0 = p_\ell\}$ is a union of no more than N_0 convex open sets, which we refer to as chambers. We emphasize that the constant $N_0 = N_0(W)$ is independent of R . Furthermore, each chamber is bounded by a finite number of line segments and at least one boundary arc from the collection $\{\tilde{I}_k\}$. Lastly, from Corollary 5.2, the number of triple junctions in the configuration u_0 is bounded by a constant depending only on N_0 , thus again a number independent of R .

Because each chamber contains the convex hull of $(1 - \rho)\tilde{I}_k$ for some k and each \tilde{I}_k has arclength at least λ , the thickness of each chamber is bounded below by $C\lambda^2$, where the thickness is defined as the minimal distance between any pair of parallel supporting planes for the chamber.

The map V_R is defined in each chamber, say Ω , as follows: Consider any segment I of $\partial\Omega \cap B_{1-\rho}$ having length at least 2η , where η is to be determined later, but where we require that

$$\eta \geq \frac{1}{R}. \quad (3.39)$$

We then consider a sub-segment $J \subset I$ of length smaller by $\eta > 0$ on each side of I and consider a rectangle in Ω with base J and height $h > 0$. It is clear that if $h > 0$ is small enough, then these rectangles are disjoint and included in Ω . We now quantify how large h is allowed to be for this property to still hold. We will always assume the bound

$$2h \leq \eta. \quad (3.40)$$

To this aim, assume $h \in (0, \eta/2]$ is the largest height for which it holds that the rectangles are mutually disjoint and included in Ω . For this value of h , if it is different from $\eta/2$, either two rectangles make contact with each other, or one rectangle makes contact with $\partial B_{1-\rho}$. In any case, let p be the projection of the contact point onto J . We denote by a and b the endpoints of I . Necessarily, there exists a point $q \in \partial\Omega \setminus I$ such that $|p - q| \leq 2h$. Moreover, choosing the horizontal axis to be the line L through a and b , since the point p is at distance at least η from both a and b , and since $2h \leq \eta$, we have that the first coordinate of q is between those of a and b .

Bringing in from infinity a line parallel to L from the half-plane containing q until it first touches Ω , we denote this first contact point by q' , and we denote by p' the orthogonal projection of q' onto L . Since $q \in \partial\Omega$, it cannot lie inside the triangle formed by a, b and q' . This implies that the segment $[pq]$ intersects either $[aq']$ or $[bq']$. We assume the former, which implies that q is further away from p than α , the orthogonal projection of p onto the line containing a and q' . So we have

$$|p - \alpha| \leq |p - q| \leq 2h. \quad (3.41)$$

See Figure 4.

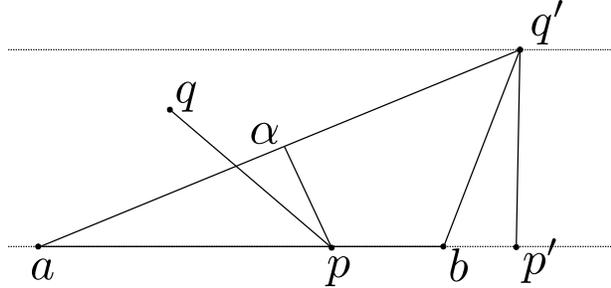


Figure 4: The configuration described above when two rectangles first touch.

The triangles $a\alpha p$ and $aq'p'$ are similar. Therefore

$$\frac{|p - \alpha|}{|a - p|} = \frac{|p' - q'|}{|a - q'|}. \quad (3.42)$$

Now $|p' - q'|$ is at least the thickness of Ω so

$$|p' - q'| \geq C\lambda^2.$$

Furthermore, since a and q' lie in $B_{1-\rho}$ we know $|a - q'| \leq 2$, and since $p \in J$, necessarily $|p - a| \geq \eta$. Combining (3.41), (3.42) with these inequalities, it follows that $h \geq C\eta\lambda^2$. Therefore, the rectangles will be disjoint and included in Ω as long as

$$h < C\eta\lambda^2, \quad (3.43)$$

where this C is $1/4$ of the constant C appearing in the previous display. The case where the segment $[pq]$ intersects $[bq']$ also leads to (3.43) in a similar manner.

Assuming (3.40), (3.43) are satisfied, consider a rectangle \mathcal{R} belonging to a chamber where $u_0 = p_i$ and sharing a boundary segment J with a chamber where $u_0 = p_j$. Then, denoting by s a coordinate orthogonal to J , we take $V_R = V_R(s)$ in \mathcal{R} given by $V_R(s) = \zeta_R(s) := \zeta_{ij}(Rs)$ for $0 \leq s \leq h/2$, where $\zeta_{ij}(0)$ is the midpoint of the heteroclinic ζ_{ij} with respect to the metric d . We take ζ_R to linearly interpolate between $\zeta_{ij}(Rh/2)$ and p_i for $h/2 \leq s \leq h$.

At this point we remark that from the assumption (1.2), it follows that each ζ_{ij} approaches its end-states p_i and p_j at an exponential rate, i.e.

$$|\zeta_{ij}(t) - p_j| \leq Ce^{-c(b)t} \quad \text{as } t \rightarrow \infty \quad (3.44)$$

for some constant $c(b) > 0$, with a similar estimate holding as $t \rightarrow -\infty$. Indeed, writing (1.7) as a first order autonomous system, say $z' = G(z)$, where

$$z = (z_1, z_2, z_3, z_4) = (\zeta_{ij}^{(1)}, \zeta_{ij}^{(2)}, \zeta_{ij}^{(1)'}, \zeta_{ij}^{(2)'}), \quad G(z) = (z_3, z_4, W_{z_1}(z_1, z_2), W_{z_1}(z_1, z_2)),$$

one checks that at any $p_\ell \in P$, the 4×4 matrix $DG(p_\ell)$ has eigenvalues $\pm\sqrt{\mu_1}$, $\pm\sqrt{\mu_2}$ where $\mu_1, \mu_2 \geq b > 0$ are the eigenvalues of $D^2W(p_\ell)$. Thus, each p_ℓ represents a hyperbolic equilibrium point from the perspective of first order ODE theory and from local stable manifold theory the approach of ζ_{ij} to p_i or p_j as $t \rightarrow \pm\infty$ must be at an exponential rate as claimed in (3.44).

In light of (3.44), the modification can be made in such a way that

$$E_R(V_R, \mathcal{R}) \leq \mathcal{H}^1(\partial\mathcal{R} \cap \partial\Omega) \left(\frac{1}{2}c_{ij} + Ce^{-c(b)Rh/2} \right). \quad (3.45)$$

In addition to (3.43), we will insist on a selection of h such that

$$Rh \gg 1, \quad (3.46)$$

so that the exponential term in (3.45) will be negligible.

At this point, we consider an extension of $V_R(1, \theta) = U_R(1, \theta)$ to the annulus $\mathcal{A}_{1+\eta, 1} := B_{1+\eta} \setminus B_1$ that is constant along rays emanating from the origin. In light of (3.25), we have

$$E_R(V_R, \mathcal{A}_{1+\eta, 1}) \leq C\eta. \quad (3.47)$$

Having defined V_R in the rectangles of each chamber, we consider the finite collection of balls of radius 2η centered at every one of the vertices of the polygonal curves $\partial\Omega \cap \bar{B}_{1-\rho}$ for every chamber Ω . Any such vertex either coincides with the location of an endpoint of the arc $(1-\rho)\tilde{I}_k$ on $\partial B_{1-\rho}$, namely a point of discontinuity of h_R , or the location of an interior triple junction. Hence, the number of vertices, say \tilde{N}_R , is bounded by a number \tilde{N} that is independent of R . Referring to the collection of these balls as $\{B_{2\eta}^j\}_{j=1}^{\tilde{N}_R}$, we next define V_R in the part of $B_{1-\rho}$ not belonging to any of the rectangles \mathcal{R} or $\cup_{j=1}^{\tilde{N}_R} B_{2\eta}^j$ by setting $V_R(x) = p_i$, when x lies in a chamber associated with p_i . This is consistent with the boundary values on $\partial B_{1-\rho}$ since by construction $V_R(x)$ is equal to p_i if x is on the boundary arc $\partial B_{1-\rho} \cap \partial\Omega$ and at distance larger than η from the ends of the arc, a condition satisfied when x is not in any $B_{2\eta}^j$.

It remains to define V_R inside the balls $B_{2\eta}^j$. To this end, we first note that for any ball $B_{2\eta}^j \subset B_{1-\rho}$, the boundary values of V_R on $\partial B_{2\eta}^j$ vary between being constant or being given by a scaled heteroclinic, hence the tangential derivative of V_R on this circle is bounded by CR . For any ball $B_{2\eta}^j$ not lying entirely in $B_{1-\rho}$, V_R on $\partial B_{2\eta}^j$ maybe also be given partially by the linear interpolation construction carried out in the annulus $\mathcal{A}_{1, 1-\rho}$ or, if the ball reaches $\mathcal{A}_{1+\eta, 1}$, then it could partially coincide with the extension described above (3.47). However, in light of the assumption (3.30), in all cases the tangential derivative of V_R along $\partial B_{2\eta}^j$ is bounded by CR for some C independent of R .

With this estimate in hand, we proceed to fill in the definition of V_R in $\cup_{j=1}^{\tilde{N}_R} B_{2\eta}^j$ sequentially as follows. Starting with $B_{2\eta}^1$, we take V_R to linearly interpolate between V_R , as previously defined, on $\partial B_{2\eta}^1$ and say, p_1 on $\partial B_{2\eta}^j$. We then take $V_R \equiv p_1$ in the ball $B_{2\eta}^j$. For such an interpolation, in light of assumption (3.39), we have

$$|\nabla V_R| \leq CR \text{ in } B_{2\eta}^1 \quad \text{and so} \quad E_R(V_R, B_{2\eta}^1) \leq CR\eta^2. \quad (3.48)$$

We then proceed to define V_R in $B_{2\eta}^2, B_{2\eta}^3, \dots$ in the same manner. However, it could happen that for some $i < j \in \{2, \dots, N_R\}$ one has $B_{2\eta}^j \cap B_{2\eta}^i \neq \emptyset$. For example, this would occur if the centers of the two balls are vertices constituting endpoints of one side of a polygonal chamber of length less than 4η . When such an intersection occurs, we simply define V_R in the intersection of these two balls using the recipe for V_R in $B_{2\eta}^j$. In light of (3.48), the $O(R)$ gradient bound is preserved through this process so that the $O(R\eta^2)$ energy bound is as well. Given that N_R is uniformly bounded by \bar{N} , a constant independent of R , we can total the energy of V_R inside $\cup_{j=1}^{\tilde{N}_R} B_{2\eta}^j$ to find an energetic contribution bounded by $CR\eta^2$.

Totaling the energetic cost of the construction

Summing the bounds (3.45) and (3.48) over rectangles and balls, we find that for any $h > 0$ satisfying (3.43) we have

$$\begin{aligned} E_R(V_R, B_{1+\eta}) &\leq m_R \left(1 + Ce^{-Rh/C}\right) + CR\eta^2 + E_R(V_R, B_{1+\eta} \setminus B_{1-\rho}) \\ &\leq m_R + C \left(R\eta^2 + e^{-Rh/C} + \rho + \frac{1}{R\rho} + \rho\lambda R + \eta\right), \end{aligned} \quad (3.49)$$

where we have used (3.37), (3.47).

We may choose η, λ, ρ, h — for instance setting $\eta = R^{-2/3}$, $\lambda = R^{-1/8}$, $\rho = R^{-8/9}$ and $h = CR^{-11/12}$ — so that for R large enough the conditions (3.30), (3.39), (3.40), (3.43) and (3.46) are satisfied. For this choice, (3.49) implies that $E_R(V_R, B_{1+\eta}) \leq m_R + CR^{-\alpha}$ with $\alpha = 1/8 - 1/9$.

Finally, recalling that $V_R(1 + \eta, \theta) = U_R(1, \theta)$ we must scale down this construction so that it agrees with U_R on ∂B_1 . Thus, with for example the choice $\eta = R^{-2/3}$ as above, we replace the sequence $V_R : B_{1+R^{-2/3}} \rightarrow \mathbb{R}^2$ with, say, $\bar{V}_R : B_1 \rightarrow \mathbb{R}^2$ given by

$$\bar{V}_R(r, \theta) := V_R\left((1 + R^{-2/3})r, \theta\right).$$

Clearly such a scaling will only affect the energy bound by lower order terms, and since $E_R(U_R, B_1) \leq E_R(\tilde{V}_R, B_1)$ we have established (3.20) for some $\alpha \in (0, 1)$.

Matching lower bound

We turn now to the task of obtaining a matching lower bound, namely (3.21).

Much in the same spirit as was done for the upper bound proof, we will replace the boundary values U_R on ∂B_1 by much simpler boundary values through interpolation. This time, however, rather than interpolating from the boundary values U_R to the much simpler boundary values $V(1 - \rho, \theta)$ on $\partial B_{1-\rho}$ as we did in the argument leading up to (3.35), we now define an *extension*, say \tilde{U}_R , of U_R to a larger ball $B_{1+\rho}$ such that $\tilde{U}_R(1 + \rho, \theta) = V_R(1 - \rho, \theta)$. This amounts to reflecting across ∂B_1 the construction in the annulus $\mathcal{A}_{1,1-\rho}$ used in the upper bound argument to instead obtain an interpolation in the annulus $\mathcal{A}_{1+\rho,1} := B_{1+\rho} \setminus B_1$.

We observe that precisely the same estimate (3.35) for the energetic cost of V_R in the annulus $\mathcal{A}_{1,1-\rho}$ will now hold in the annulus $\mathcal{A}_{1+\rho,1}$ for the extension \tilde{U}_R . Hence, again making the choices taking $\rho = R^{-8/9}$ and $\lambda = R^{-1/8}$ indicated below (3.49), we conclude from (3.37) that

$$\begin{aligned}
& \sqrt{2} \int_{B_1} \sqrt{W(U_R)} |\nabla U_R| dx = \\
& \sqrt{2} \int_{B_{1+\rho}} \sqrt{W(\tilde{U}_R)} |\nabla \tilde{U}_R| dx - \sqrt{2} \int_{\mathcal{A}_{1+\rho,1}} \sqrt{W(\tilde{U}_R)} |\nabla \tilde{U}_R| dx \\
& \geq \sqrt{2} \int_{B_{1+\rho}} \sqrt{W(\tilde{U}_R)} |\nabla \tilde{U}_R| dx - E_R(\tilde{U}_R, \mathcal{A}_{1+\rho,1}) \\
& \geq \sqrt{2} \int_{B_{1+\rho}} \sqrt{W(\tilde{U}_R)} |\nabla \tilde{U}_R| dx - O(R^{-\alpha}), \tag{3.50}
\end{aligned}$$

where, as we did earlier, we have set $\alpha = 1/8 - 1/9$.

We now define three open subsets of $B_{1+\rho}$ via

$$\Omega_\ell^R := \left\{ x \in B_{1+\rho} : d(\tilde{U}_R(x), p_\ell) < t_\ell \right\} \text{ for } \ell = 1, 2 \text{ and } 3,$$

where t_1, t_2 and t_3 are defined in (3.18). These sets are disjoint since, for instance, $d(p, p_1) < t_1$ implies that

$$d(p, p_2) > d(p_1, p_2) - t_1 = c_{12} - t_1 = t_2,$$

in light of (3.18). Then we invoke the property of the metric d that for any fixed ‘base point,’ $p \in \mathbb{R}^2$, one has

$$|\nabla_q d(p, q)| = \sqrt{W(q)} \text{ for all } q \in \mathbb{R}^2, \tag{3.51}$$

cf. e.g. [26], along with the co-area formula to estimate that

$$\begin{aligned}
& \sqrt{2} \int_{B_{1+\rho}} \sqrt{W(\tilde{U}_R)} |\nabla \tilde{U}_R| dx \geq \sum_{\ell=1}^3 \sqrt{2} \int_{\Omega_\ell^R} \sqrt{W(\tilde{U}_R)} |\nabla \tilde{U}_R| dx \\
& = \sum_{\ell=1}^3 \int_{\Omega_\ell^R} |\nabla d(\tilde{U}_R(x), p_\ell)| dx \geq \sum_{\ell=1}^3 \int_{\frac{1}{R^{1/2}}}^{t_\ell} \mathcal{H}^1 \left(\{x : d(\tilde{U}_R(x), p_\ell) = s\} \right) ds \\
& \geq \sum_{\ell=1}^3 \inf_{s \in [\frac{1}{R^{1/2}}, t_\ell]} \mathcal{H}^1 \left(\{x : d(\tilde{U}_R(x), p_\ell) = s\} \right) \left(t_\ell - \frac{1}{R^{1/2}} \right) \\
& \geq \sum_{\ell=1}^3 t_\ell \mathcal{H}^1 \left(\{x : d(\tilde{U}_R(x), p_\ell) = s_\ell^*\} \right) - \frac{C}{R^{1/2}}, \tag{3.52}
\end{aligned}$$

for some numbers $s_\ell^* \in [\frac{1}{R^{1/2}}, t_\ell]$ for $\ell = 1, 2, 3$ and some C independent of R . (The lack of R dependence for C is clear since C represents the minimal length of a level set $\{d(\tilde{U}_R(x), p_\ell) = s\}$ among $s \in [\frac{1}{R^{1/2}}, t_\ell]$ within Ω_ℓ^R .)

With the goal of applying Theorem 5.4 to the triple $\{x : d(\tilde{U}_R(x), p_\ell) < s_\ell^*\}$, $\ell = 1, 2, 3$, we now wish to estimate the measure of the set $B_1 \setminus \cup_{\ell=1}^3 \{x : d(\tilde{U}_R(x), p_\ell) < s_\ell^*\}$. Since $s_\ell^* \geq \frac{1}{R^{1/2}}$ we have that

$$B_1 \setminus \{x : d(\tilde{U}_R(x), P) < s_\ell^*\} \subset B_1 \setminus \{x : d(\tilde{U}_R(x), P) < \frac{1}{R^{1/2}}\}. \tag{3.53}$$

But with an appeal to the nondegeneracy assumption (1.2) we can assert that if

$$d(\tilde{U}_R(x), p_\ell) \geq \frac{1}{R^{1/2}} \text{ for } \ell = 1, 2 \text{ and } 3, \text{ then } W(\tilde{U}_R(x)) \geq \frac{C}{R^{1/2}}, \tag{3.54}$$

for a constant C depending on W . Indeed, through an appeal to (3.27), the definition of the metric d and the convexity of W near P , we see that for $\tilde{U}_R(x)$ in a β -neighborhood of $p_\ell \in P$, one has

$$\begin{aligned}
\frac{1}{R^{1/2}} & \leq d(\tilde{U}_R(x), p_\ell) \leq \int_0^1 \sqrt{W((1-t)p_\ell + t\tilde{U}_R(x))} |\tilde{U}_R(x) - p_\ell| dt \\
& \leq \int_0^1 \sqrt{tW(\tilde{U}_R(x))} |\tilde{U}_R(x) - p_\ell| dt \leq \frac{2}{3} \sqrt{2b} |\tilde{U}_R(x) - p_\ell|^2.
\end{aligned}$$

Then another appeal to (3.27) yields (3.54).

Then from the bound (2.18), it follows that

$$C_1 \geq \int_{B_1 \setminus \{x : d(\tilde{U}_R(x), P) < \frac{1}{R^{1/2}}\}} RW(\tilde{U}_R(x)) dx \geq CR^{1/2} \left| B_1 \setminus \left\{ x : d(\tilde{U}_R(x), P) < \frac{1}{R^{1/2}} \right\} \right|.$$

Hence, in view of (3.53), we find that

$$\left| B_1 \setminus \left\{ x : d(\tilde{U}_R(x), P) < s_\ell^* \right\} \right| \leq \frac{C_0}{R^{1/2}} \quad (3.55)$$

for some $C_0 = C_0(W)$.

With estimate (3.55) in hand, we would like to now apply Theorem 5.4 to the triple $\{x : d(\tilde{U}_R(x), p_\ell) < s_\ell^*\}$, $\ell = 1, 2, 3$ with the Dirichlet condition on $\partial B_{1+\rho}$ given by $h = h_R$, for h_R as defined above (3.38). However, to do so, we must make minor adjustments to these three sets near $\partial B_{1+\rho}$. These adjustments entail adding to or subtracting from these sets small slices of $B_{1+\rho}$ bounded by arcs of $\partial B_{1+\rho}$ and secant lines, so as to ‘fix’ their traces to match those dictated by h_R . By our construction of $V(1-\rho, \theta)$, hence of $\tilde{U}_R(1+\rho, \theta)$, these adjustments occur along N_R arcs that are contained in the set $\cup_{k=1}^{N_R} (1+\rho)\tilde{I}_k \setminus \cup(1+\rho)I_k$, where we recall that N_R is bounded by a constant depending only on W . Therefore, in view of (3.31), one can alter the three sets so as to obtain a triple whose trace on $\partial B_{1+\rho}$ matches h_R exactly, and the extra cost in perimeter will be $O(\lambda) = O(R^{-1/8})$. Furthermore, the estimate (3.55) will still hold since the adjustments are lower order.

With this adjustment in hand, we now return to (3.50) and (3.52), and apply Theorem 5.4 with $\delta = \frac{C_0}{R^{1/2}}$, to obtain

$$\begin{aligned} & \sqrt{2} \int_{B_1} \sqrt{W(U_R)} |\nabla U_R| dx \\ & \geq \sum_{\ell=1}^3 t_\ell \mathcal{H}^1(\{x : d(\tilde{U}_R(x), p_\ell) = s_\ell^*\}) - \frac{C}{R^\alpha} \\ & \geq \min \{E_0(u, B_{1+\rho}) : u = h_R \text{ on } \partial B_{1+\rho}\} - \frac{\sqrt{C_0}\gamma(k)}{R^{1/4}} - \frac{C}{R^{1/2}} - \frac{C}{R^{1/8}} - \frac{C}{R^\alpha} \\ & \geq \min \{E_0(u, B_{1-\rho}) : u = h_R \text{ on } \partial B_{1+\rho}\} - \frac{C}{R^\alpha} = m_R - \frac{C}{R^\alpha}. \end{aligned} \quad (3.56)$$

This is the lower bound (3.21) we were seeking, and so the proof of Proposition 3.3 is complete. \square

3.2 Convergence of the blowdowns to a minimal cone

With the crucial Proposition 3.3 now in hand, we apply (3.15) to (3.14) with $R_1 = R$ and $R_2 \in (R_1, 2R_1]$ to obtain

$$\tilde{W}(R_2) - \tilde{W}(R_1) \geq -C_3 R_1^{-\alpha/2} \quad (3.57)$$

for C_3 depending only on W .

One consequence of (3.57) is:

Lemma 3.4. *Assume $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is a local minimizer of E and define \tilde{W}_R by (3.10). Then the limit $L_0 := \lim_{R \rightarrow \infty} \tilde{W}(R)$ exists.*

Proof. Using (3.57) we will first argue that for any $\eta > 0$ there exists a value $R_0 > 0$ such that

$$\tilde{W}(R') - \tilde{W}(R) > -\eta \quad \text{whenever } R_0 \leq R < R'. \quad (3.58)$$

To see this, let k be the largest integer such that $2^k R < R'$. Then we see that

$$\begin{aligned} \tilde{W}(R') - \tilde{W}(R) &= \tilde{W}(R') - \tilde{W}(2^k R) + \sum_{j=0}^{k-1} \left(\tilde{W}(2^{j+1} R) - \tilde{W}(2^j R) \right) \\ &\geq -C_3 \sum_{j=0}^k (2^j R)^{-\alpha/2} \geq -C_3 \frac{1}{R^{\alpha/2}} \left(\frac{1}{1 - 2^{-\alpha/2}} \right). \end{aligned}$$

Thus, taking R large enough, we obtain (3.58).

Let us now suppose $\lim_{R \rightarrow \infty} \tilde{W}(R)$ does not exist and seek a contradiction. Since by (2.18) we know that $0 < \tilde{W}(R) \leq C_1$ for all $R > 0$, this would imply that there exist sequences $\{R_j\} \rightarrow \infty$ and $\{R_k\} \rightarrow \infty$ such that $\lim_{R_j \rightarrow \infty} \tilde{W}(R_j) < \lim_{R_k \rightarrow \infty} \tilde{W}(R_k)$; say

$$\lim_{R_k \rightarrow \infty} \tilde{W}(R_k) - \lim_{R_j \rightarrow \infty} \tilde{W}(R_j) = L \quad \text{for some } L > 0. \quad (3.59)$$

Hence, there exist J and K such that for $j > J$ and $k > K$ we would have

$$\tilde{W}(R_k) - \tilde{W}(R_j) > \frac{L}{2}.$$

However, by perhaps taking j even larger we may find R_j such that $R_j > R_k$ and then an application of (3.58) with $\eta = \frac{L}{3}$ leads to the condition

$$\tilde{W}(R_j) - \tilde{W}(R_k) \geq -\frac{L}{3},$$

and the contradiction is complete. \square

Another consequence of (3.15) is the following:

Lemma 3.5. *Assume $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is a local minimizer of E . Then for every positive $\lambda_1 < \lambda_2$ we have*

$$\lim_{R \rightarrow \infty} \int_{B_{\lambda_2 R} \setminus B_{\lambda_1 R}} \frac{1}{|x|} |U_\nu|^2 dx = 0, \quad (3.60)$$

where $U_\nu = \nabla U \cdot \frac{x}{|x|}$.

We also have

$$\lim_{R \rightarrow \infty} \int_{B_{\lambda_2 R} \setminus B_{\lambda_1 R}} \frac{1}{2|x|} \left| W(U) - \frac{1}{2} |\nabla U|^2 \right| dx = 0, \quad (3.61)$$

or equivalently,

$$\lim_{R \rightarrow \infty} \int_{B_{\lambda_2} \setminus B_{\lambda_1}} \frac{1}{2|x|} \left| RW(U_R) - \frac{1}{2R} |\nabla U_R|^2 \right| dx = 0. \quad (3.62)$$

Proof. To establish the limit (3.61), we note that with the choices $R_1 = \lambda_1 R$ and $R_2 = \lambda_2 R$, the inequalities (3.12) and (3.13), followed by application of Proposition 3.3 imply that

$$\begin{aligned} & \int_{B_{\lambda_2 R} \setminus B_{\lambda_1 R}} \frac{1}{2|x|} \left| W(U) - \frac{1}{2} |\nabla U|^2 \right| dx \\ & \leq \sqrt{\frac{C_1 \lambda_2}{2\lambda_1}} R^{-1/2} \left\{ \int_{B_{\lambda_2 R} \setminus B_{\lambda_1 R}} \left(\sqrt{W(U)} - \frac{1}{\sqrt{2}} |\nabla U| \right)^2 dx \right\}^{1/2} \\ & \leq \sqrt{\frac{C_1 C_0 \lambda_2}{2\lambda_1}} R^{-1/2} R^{1/2(1-\alpha)} = O(R^{-\alpha/2}). \end{aligned}$$

Then (3.60) follows from (3.11) and (3.61), in light of Lemma 3.4. \square

Now for any $\lambda > 0$ it follows from Lemma 3.4 that

$$\lim_{R \rightarrow \infty} \int_{B_\lambda} RW(U_R) dx = \lambda L_0.$$

Combining this with (3.62) yields that for any $0 < \lambda_1 < \lambda_2$ one has

$$\lim_{R \rightarrow \infty} E_R(U_R, B_{\lambda_2} \setminus B_{\lambda_1}) = 2(\lambda_2 - \lambda_1)L_0. \quad (3.63)$$

We can rephrase (3.60) in terms of the blowdowns as

$$\lim_{R \rightarrow \infty} \frac{1}{R} \int_{B_{\lambda_2} \setminus B_{\lambda_1}} \frac{1}{|x|} \left| \nabla U_R \cdot \frac{x}{|x|} \right|^2 dx = 0 \quad \text{for any } 0 < \lambda_1 < \lambda_2. \quad (3.64)$$

Now we will use this to argue that the limit of blowdowns is necessarily a cone:

Theorem 3.6. *Assume $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is a local minimizer of E . Let $\{R_j\} \rightarrow \infty$ be any sequence and let $\{R_{j_k}\}$ and a function $u_0 \in BV(B_1; P)$ be any subsequence and subsequential L^1 limit guaranteed by Proposition 3.1. If we denote by $\Gamma_{i\ell}$ the phase boundary $\partial\{u_0 = p_i\} \cap \partial\{u_0 = p_\ell\}$, one has*

$$\nu_{i\ell}(x) \cdot x = 0 \quad \text{for every nonzero } x \in B_1 \cap \Gamma_{i\ell}^* \text{ and every } 1 \leq i < \ell \leq 3, \quad (3.65)$$

where $\Gamma_{i\ell}^*$ denotes the reduced boundary of $\Gamma_{i\ell}$ and $\nu_{i\ell}$ denotes a corresponding normal vector.

Proof. We fix any positive number δ and note that for any $\mu \in (0, 1)$ we have

$$\begin{aligned} \int_{B_1 \setminus B_\mu} \sqrt{W(U_R)} \left| \nabla U_R \cdot \frac{x}{|x|} \right| dx &\leq \delta R \int_{B_1} W(U_R) dx + \frac{1}{\delta} \frac{1}{R} \int_{B_1 \setminus B_\mu} \left| \nabla U_R \cdot \frac{x}{|x|} \right|^2 dx \\ &\leq \delta \tilde{W}(R) + \frac{1}{\delta} \frac{1}{R} \int_{B_1 \setminus B_\mu} \frac{1}{|x|} \left| \nabla U_R \cdot \frac{x}{|x|} \right|^2 dx. \end{aligned}$$

Then sending $R \rightarrow \infty$ and invoking Lemma 3.4 and (3.64) we conclude that

$$\limsup_{R \rightarrow \infty} \int_{B_1 \setminus B_\mu} \sqrt{W(U_R)} \left| \nabla U_R \cdot \frac{x}{|x|} \right| dx \leq \delta L_0,$$

and since δ was arbitrary it follows that

$$\lim_{R \rightarrow \infty} \int_{B_1 \setminus B_\mu} \sqrt{W(U_R)} \left| \nabla U_R \cdot \frac{x}{|x|} \right| dx = 0. \quad (3.66)$$

Next, we consider the function $x \mapsto d(p_1, U_R(x))$. Suppressing subsequential notation, the fact that $U_{R_j} \rightarrow u_0$ in L^1 implies through definition (1.4) that

$$d(p_1, U_{R_j}) \xrightarrow{L^1(B_1)} d(p_1, u_0) = \begin{cases} 0 & \text{on } \{u_0 = p_1\} \\ d(p_1, p_2) & \text{on } \{u_0 = p_2\}, \\ d(p_1, p_3) & \text{on } \{u_0 = p_3\} \end{cases}$$

and we note that since $u_0 \in BV(B_1; \mathbb{R}^2)$ one has $x \mapsto d(p_1, u_0(x)) \in BV(B_1)$. Hence, in the sense of distributions, we have

$$\nabla_x d(p_1, U_{R_j}(x)) \rightarrow \nabla_x d(p_1, u_0(x)) = \sum_{1 \leq i < \ell \leq 3} (d(p_1, p_\ell) - d(p_1, p_i)) \nu_{i\ell} \mathcal{H}^1 \llcorner \Gamma_{i\ell}^*, \quad (3.67)$$

where this form of the distributional gradient follows, for example, from [17], Prop. 2.8. Then recalling (3.51), an elementary chain rule calculation shows that

$$\left| \nabla_x d(p_1, U_R(x)) \cdot \frac{x}{|x|} \right| \leq \sqrt{W(U_R(x))} \left| \nabla U_R \cdot \frac{x}{|x|} \right|,$$

and so from (3.66) we may conclude that

$$\lim_{R \rightarrow \infty} \int_{B_1 \setminus B_\mu} \left| \nabla_x d(p_1, U_R(x)) \cdot \frac{x}{|x|} \right| dx = 0. \quad (3.68)$$

Fix now any non-zero point $x_0 \in \Gamma_{i\ell}^*$ for some $1 \leq i < \ell \leq 3$, and fix $\mu > 0$ less than $|x_0|$. Then we take an arbitrary $\phi \in C_0^1(B_r(x_0))$, with r chosen small enough so that $B_r(x_0) \subset B_1 \setminus B_\mu$. It follows from (3.67) and (3.68) that

$$\int_{\Gamma_{i\ell}^* \cap B_r(x_0)} \phi \nu_{i\ell} \cdot \frac{x}{|x|} d\mathcal{H}^1(x) = 0.$$

Since ϕ is arbitrary, we obtain the desired property (3.65). \square

In light of Proposition 3.1, we know that any limit of blowdowns, u_0 , minimizes E_0 subject to its own boundary values. Now that we also know any limit of blowdowns is a cone, it follows immediately from Theorem 5.1 that there are only three possibilities:

Proposition 3.7. *Under the hypothesis and with the notations of Theorem 3.6, either*

$$(i) \quad u_0 \equiv p_i \text{ for some } i \in \{1, 2, 3\}, \quad (3.69)$$

or there exists a half-plane H with ∂H passing through the origin such that

$$(ii) \quad u_0(x) = \begin{cases} p_i & \text{in } H \\ p_\ell & \text{in } \mathbb{R}^2 \setminus H, \end{cases} \quad (3.70)$$

for some $i, \ell \in \{1, 2, 3\}$ with $i \neq \ell$, or

$$(iii) \quad u_0 \text{ takes the form (1.8)} \quad (3.71)$$

with the three sectors S_1, S_2 and S_3 having opening angles α_1, α_2 and α_3 satisfying the condition (1.9).

Remark 3.8. *It was recently pointed out to us by Michael Novack that there is a different avenue available to reach the conclusion that any subsequential limit u_0 of blowdowns of a local minimizer of E must satisfy either (3.69), (3.70) and (3.71). This alternative argument utilizes, among other tools, a monotonicity formula for minimizing partitions adapted to this setting, along with available regularity theory such as that described in Theorem 5.1, to show that the only locally minimizing partitions of the plane with respect to the energy E_0 are partitions fitting one of these three descriptions. Then coupled with Proposition 3.1 we would reach the same conclusion as that of Proposition 3.7. See also [2] for a presentation of this property of minimizing partitions.*

4 Eliminating the half-plane

We will now argue that in our setting, neither (3.69) or (3.70) are possible, leaving (3.71) as the only option, thus leading to a proof of our main result, Theorem 1.1.

We will first prove a ‘clearing-out’ type of result, saying that sufficiently low energy in a ball implies uniform nearness to a potential well on a smaller ball. We remark that a result of this type for mere solutions to (2.11) is established in [11], Prop. 6.4, but since the proof is considerably simpler in the setting of local minimizers, we present a proof in this setting below.

Proposition 4.1. *For any $R > 0$, let z_R be a local minimizer of E_R satisfying a gradient bound $|\nabla z_R| \leq C_1 R$ for some $C_1 > 0$. Then there exists a number η depending only on W such that if*

$$E_R(z_R, B_{r_0}) < \eta \quad \text{on some ball } B_{r_0},$$

then there exists a point $p_\ell \in P$ and a value $\bar{R} > 0$ such that for all $x \in B_{r_0/2}$, one has the uniform estimate

$$|z_R(x) - p_\ell| < \sqrt{3} \left(\frac{2}{b} \right)^{1/4} \left(E_R(z_R, B_{r_0}) \right)^{1/2}$$

for all $R \geq \bar{R}$, where b is the constant appearing in (3.27) and $\bar{R} = \bar{R}(b, r_0)$.

Proof. With no loss of generality, we take B_{r_0} to be centered at the origin. For ease of notation, we will write $e_R := E_R(z_R, B_{r_0})$, so that our hypothesis is that $e_R < \eta$ with η to be specified shortly.

We begin with the observation that if $q \in \mathbb{R}^2$ is any point such that $|q - p_\ell| \leq \beta$ for some $p_\ell \in P$, then invoking (1.4) and (3.27), we have

$$\begin{aligned} \sqrt{\frac{b}{2}} \min_{\gamma(0)=p_\ell, \gamma(1)=q} \int_0^1 |\gamma(t) - p_\ell| |\gamma'(t)| dt &\leq d(q, p_\ell) \leq \\ \sqrt{b} \min_{\gamma(0)=p_\ell, \gamma(1)=q} \int_0^1 |\gamma(t) - p_\ell| |\gamma'(t)| dt. \end{aligned}$$

Hence,

$$\sqrt{\frac{b}{2}} |q - p_\ell|^2 \leq d(q, p_\ell) \leq \sqrt{b} |q - p_\ell|^2. \quad (4.1)$$

Applying the Mean Value Theorem, the assumption $e_R < \eta$ on B_{r_0} allows us to find a radius, say $r^* \in (r_0/2, r_0)$, such that

$$E_R(z_R, \partial B_{r^*}) < 2e_R < 2\eta. \quad (4.2)$$

Consequently, for any two points $x_1, x_2 \in \partial B_{r^*}$ one has

$$d(z_R(x_1), z_R(x_2)) \leq \sqrt{2} \int_{\partial B_{r^*}} \sqrt{W(z_R)} |\nabla z_R| ds < 2e_R.$$

Then the fact that $\int_{\partial B_{r^*}} RW(z_R) ds < 2e_R$ implies that for R large enough there exists a point, say $x_R \in \partial B_{r^*}$, such that $d(z_R(x_R), p_\ell) < e_R$ for some $p_\ell \in P$. Consequently, it follows from the triangle inequality that

$$d(z_R(x), p_\ell) < 3e_R \quad \text{for all } x \in \partial B_{r^*}. \quad (4.3)$$

We now impose the condition

$$\eta \leq \sqrt{\frac{b}{2}} \frac{\beta^2}{3}. \quad (4.4)$$

It follows that

$$|z_R(x) - p_\ell| \leq \beta \quad \text{for all } x \in \partial B_{r^*}. \quad (4.5)$$

Otherwise, for some $x \in \partial B_{r^*}$, we would have

$$d(z_R(x), p_\ell) \geq \min_{\{q: |q-p_\ell|=\beta\}} d(q, p_\ell) \geq \sqrt{\frac{b}{2}} \beta^2,$$

contradicting (4.3), given that $e_R < \eta$.

Having established (4.5), we now appeal to the local minimality of z_R by constructing a competitor, say v_R , in B_{r^*} that linearly interpolates on the annulus $\mathcal{A}_{r^*, r^* - \frac{1}{R}}$ between $z_R(x)$ on ∂B_{r^*} and p_ℓ on $\partial B_{r^* - \frac{1}{R}}$ via the formula

$$v_R(x) := \lambda_R(|x|) z_R\left(r^* \frac{x}{|x|}\right) + (1 - \lambda_R(|x|)) p_\ell, \quad (4.6)$$

for $r^* - \frac{1}{R} \leq |x| \leq r^*$, where $\lambda_R(r) := R(r - r^*) + 1$. We compute that

$$\nabla v_R(x) = \lambda_R(|x|) \nabla z_R\left(r^* \frac{x}{|x|}\right) + R \frac{x}{|x|} \otimes \left(z_R\left(r^* \frac{x}{|x|}\right) - p_\ell\right), \quad (4.7)$$

so that

$$|\nabla v_R(x)|^2 \leq 2 \left| \nabla z_R\left(r^* \frac{x}{|x|}\right) \right|^2 + 2R^2 \left| z_R\left(r^* \frac{x}{|x|}\right) - p_\ell \right|^2. \quad (4.8)$$

Property (4.5) guarantees that this interpolation always yields values inside the ball of radius β about p_ℓ so that v_R takes its values in the region where W

is convex. This convexity allows us to invoke (3.27). Then through the local minimality of z_R , along with (4.2) and (4.8), we find that

$$\begin{aligned}
E_R(z_R, B_{r^*}) &\leq E_R(v_R, B_{r^*}) = \int_{\mathcal{A}_{r^*, r^* - \frac{1}{R}}} RW(v_R) + \frac{1}{2R} |\nabla v_R|^2 dx \\
&\leq \int_{r^* - \frac{1}{R}}^{r^*} \int_{\partial B_r} \left\{ R\lambda(r)W\left(z_R\left(r^* \frac{x}{|x|}\right)\right) + \right. \\
&\quad \left. \frac{1}{2R} \left(2 \left| \nabla z_R\left(r^* \frac{x}{|x|}\right) \right|^2 + \frac{2R^2}{b} W\left(z_R\left(r^* \frac{x}{|x|}\right)\right) \right) \right\} ds dr \leq 2\eta (\max\{2, 1 + \frac{1}{b}\}) \frac{1}{R}.
\end{aligned} \tag{4.9}$$

Next we wish to argue that the maximum of the quantity $|z_R(x) - p_\ell|$ over the set $\overline{B_{r^*}}$ must occur for $x \in \partial B_{r^*}$. We will argue by contradiction. There are two cases to consider:

Case 1: The maximum occurs at a point $x^* \in B_{r^*}$ such that $|z_R(x^*) - p_\ell| < \beta$. We see this is impossible through an appeal to the maximum principle applied to the function $f(x) := \frac{1}{2} |z_R(x) - p_\ell|^2$. Indeed, a simple calculation yields that

$$\Delta f(x) = R^2 \nabla_u W(z_R) \cdot (z_R - p_\ell) + |\nabla z_R|^2 > 0,$$

in light of the strict convexity of $W(q)$ when $|q - p_\ell| \leq \beta$.

Case 2: The maximum occurs at a point $x^* \in B_{r^*}$ such that $|z_R(x^*) - p_\ell| \geq \beta$. Since necessarily any local minimizer satisfies a gradient estimate $|\nabla z_R| \leq C_1 R$ for some constant $C_1 > 0$, it follows that

$$|z_R(x) - p_\ell| > \frac{\beta}{2} \quad \text{for all } x \text{ such that } |x - x^*| < \frac{\beta}{2C_1 R}.$$

Thus, denoting

$$C_\beta := \min_{\{q \in \mathbb{R}^2: \text{dist}(q, P) \geq \frac{\beta}{2}\}} W(q) > 0,$$

we obtain

$$E_R(z_R, B_{r^*}) \geq \int_{\{x: |x - x^*| < \frac{\beta}{2C_1 R}\}} RW(z_R) dx \geq C_\beta R \pi \left(\frac{\beta}{2C_1 R} \right)^2 = \frac{\pi \beta^2 C_\beta}{4C_1^2} \frac{1}{R},$$

which will contradict (4.9), if in addition to (4.4), we insist that η satisfies, say

$$2\eta \max\{2, 1 + \frac{1}{b}\} \leq \frac{\pi \beta^2 C_\beta}{5C_1^2}. \tag{4.10}$$

Hence, assuming η satisfies (4.4) and (4.10), we have argued that the maximum of $|z_R - p_\ell|$ on B_{r^*} must occur on ∂B_{r^*} , and so from (4.1) and (4.3), the conclusion of the Proposition follows. \square

An easy consequence of this result is

Proposition 4.2. *Assume $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is a local minimizer of E and $\{R_j\} \rightarrow \infty$ is such that $\{U_{R_j}\}$ converges in L^1_{loc} to $u_0 \in BV_{\text{loc}}(\mathbb{R}^2; P)$. Then $\{U_{R_j}\}$ converges to u_0 locally uniformly outside the support of ∇u_0 .*

Proof. Assume x_0 does not belong to the support of ∇u_0 . We need to prove that U_{R_j} converges uniformly to u_0 in a neighbourhood of x_0 .

The limit u_0 is identically equal to one of the wells, say p_1 , in a ball $B_r(x_0)$ for some $r > 0$. Using Fatou's Lemma as in the proof of Lemma 2.3 there exists a radius $t \in (0, r)$ and a subsequence still denoted $\{R_j\}$ such that

$$\lim_{R_j \rightarrow \infty} \|U_{R_j} - p_1\|_{L^1(\partial B_t)} = 0, \quad \limsup_{R_j \rightarrow \infty} E_{R_j}(U_{R_j}, \partial B_t) < +\infty.$$

Since U_{R_j} minimizes E_{R_j} and satisfies (2.2) and (2.3) on ∂B_t , we can apply condition (2.7) of Theorem 2.1 to assert that

$$E_{R_j}(U_{R_j}, B_t) \rightarrow E_0^h(u_0, B_t) = 0 \text{ as } R_j \rightarrow \infty,$$

where h is the trace of u_0 on ∂B_t , that is, $h = p_1$. This allows us to apply Proposition 4.1 to conclude that U_{R_j} is converging uniformly to p_1 on $B_{t/2}$. Since the subsequential limit is unique, the whole sequence converges uniformly to p_1 on $B_{t/2}$, proving the proposition. \square

Now we prove

Proposition 4.3. *Assume $U : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is a local minimizer of E in the sense of (1.13) such that $\text{dist}(U(0), \Lambda) > 0$ for Λ given by (2.12). For any sequence $\{R_j\} \rightarrow \infty$, let $\{R_{j_k}\}$ and $u_0 \in BV_{\text{loc}}(\mathbb{R}^2; P)$ be any subsequence and subsequential limit of $\{U_{R_{j_k}}\}$, guaranteed to exist by Proposition 3.1. Then u_0 takes the form (1.8).*

Proof. We need to rule out (3.69) and (3.70).

Suppose first, by way of contradiction, that $u_0 \equiv p_\ell$ for some $\ell \in \{1, 2, 3\}$. Then Proposition 4.2 implies that $U_{R_{j_k}}(0) \rightarrow U(0) = p_\ell$, contradicting $\text{dist}(U(0), \Lambda) > 0$. Thus, possibility (3.69) is eliminated.

Next, we suppose by way of contradiction that $U_{R_{j_k}} \rightarrow u_0$ in L^1_{loc} for u_0 satisfying (3.70) for some $i, \ell \in \{1, 2, 3\}$ with $i \neq \ell$. With no loss of generality we

will take $i = 1$, $\ell = 2$ and the halfplane H to be $\{(x_1, x_2) \in \mathbb{R}^2 : x_2 < 0\}$ so that our contradiction hypothesis takes the form

$$U_{R_{j_k}} \xrightarrow{L_{\text{loc}}^1} u_0 = \begin{cases} p_1 & \text{in } \{(x_1, x_2) : x_2 < 0\} \\ p_2 & \text{in } \{(x_1, x_2) : x_2 > 0\}. \end{cases} \quad (4.11)$$

This possibility is ruled out if we prove Theorem 1.3, since the latter implies that $U(x_1, x_2) = \zeta_{12}(x_2 + \Delta)$ for some $\Delta \in \mathbb{R}$ and therefore $U(0) \in \zeta_{12}(\mathbb{R})$, which contradicts the hypothesis that $\text{dist}(U(0), \Lambda) > 0$ for Λ given by (2.12). \square

Proof of Theorem 1.3. We assume that U is a locally minimizing entire solution and that U_{R_j} converges as $j \rightarrow +\infty$ for some subsequence for $R_j \rightarrow +\infty$ to the function u_0 defined in (1.16). We break the proof that $U(x_1, x_2) = \zeta_{12}(x_2 + \Delta)$ for some $\Delta \in \mathbb{R}$ into several steps.

1. We begin by identifying a circle on which U has “well-controlled” boundary values. To this end, we note that the argument in the proof of Proposition 3.1 leading to (3.2) and (3.3) applies just as well to assert the existence of *two* values, $1 \leq \lambda_1 < \lambda_2 \leq 2$, for which these two properties hold on both ∂B_{λ_1} and ∂B_{λ_2} . We then let $\mathcal{A}_{\lambda_2, \lambda_1}$ denote the annulus $B_{\lambda_2} \setminus B_{\lambda_1}$, and invoke the assumption (4.11). It follows from the Γ -convergence of $\tilde{E}_{R_j}(\cdot, \mathcal{A}_{\lambda_2, \lambda_1})$ to $E_0^h(\cdot, \mathcal{A}_{\lambda_2, \lambda_1})$ with $h = \text{trace of } u_0 \text{ on } \partial \mathcal{A}_{\lambda_2, \lambda_1}$, along with the minimality of U_{R_j} in the annulus, that

$$E_{R_j}(U_{R_j}, \mathcal{A}_{\lambda_2, \lambda_1}) \rightarrow E^h(u_0, \mathcal{A}_{\lambda_2, \lambda_1}) = 2(\lambda_2 - \lambda_1)c_{12} \text{ as } R_j \rightarrow \infty. \quad (4.12)$$

Rewriting this in terms of U we have that

$$E(U, \mathcal{A}_{\lambda_2 R_j, \lambda_1 R_j}) = 2(\lambda_2 - \lambda_1)c_{12}R_j + o(R_j) \text{ as } R_j \rightarrow \infty, \quad (4.13)$$

or equivalently,

$$\frac{1}{(\lambda_2 - \lambda_1)R_j} \int_{\lambda_1 R_j}^{\lambda_2 R_j} E(U, \partial B_r) dr = 2c_{12} + o(1).$$

Thus, by the Mean Value Theorem, there must exist a sequence of radii $\{\rho_j\} \rightarrow \infty$ with $\rho_j \in (\lambda_1 R_j, \lambda_2 R_j)$ such that

$$E(U, \partial B_{\rho_j}) = 2c_{12} + o(1). \quad (4.14)$$

From Proposition 4.2, we have that for any $\tau > 0$

$$\max \{|U(x) - p_1| : x \in \overline{B}_{\rho_j} \cap \{(x_1, x_2) : x_2 \geq \rho_j \tau\}\} \rightarrow 0 \quad (4.15)$$

and

$$\max \{|U(x) - p_2| : x \in \overline{B}_{\rho_j} \cap \{(x_1, x_2) : x_2 \leq -\rho_j \tau\}\} \rightarrow 0 \quad \text{as } \rho_j \rightarrow \infty. \quad (4.16)$$

2. We will next argue that the restriction of U to the circle ∂B_{ρ_j} is approaching two copies of the geodesic ζ_{12} as $\rho_j \rightarrow \infty$. We denote by $\partial B_{\rho_j}^+$ the right half-circle of ∂B_{ρ_j} , and by $\partial B_{\rho_j}^-$ the left half-circle. From (4.15) and (4.16) and the continuity of the metric distance $(p, q) \mapsto d(p, q)$ (cf. (1.4)), it follows that

$$E(U, \partial B_{\rho_j}^+) \geq \sqrt{2} \int_{\partial B_{\rho_j}^+} \sqrt{W(U)} \left| \frac{\partial U}{\partial s} \right| ds \geq d(p_1, p_2) - o(1) = c_{12} - o(1) \quad \text{as } \rho_j \rightarrow \infty \quad (4.17)$$

and similarly,

$$E(U, \partial B_{\rho_j}^-) \geq \sqrt{2} \int_{\partial B_{\rho_j}^-} \sqrt{W(U)} \left| \frac{\partial U}{\partial s} \right| ds \geq d(p_1, p_2) - o(1) = c_{12} - o(1) \quad \text{as } \rho_j \rightarrow \infty \quad (4.18)$$

Combining these last two inequalities with (4.14), we observe that, in fact, we have

$$E(U, \partial B_{\rho_j}^+) \rightarrow c_{12}, \quad \text{as } \rho_j \rightarrow \infty \quad (4.19)$$

and

$$E(U, \partial B_{\rho_j}^-) \rightarrow c_{12} \quad \text{as } \rho_j \rightarrow \infty. \quad (4.20)$$

Now, following the general scheme in [5], sect. 2.3 we fix a positive number d_0 less than say, half the minimal distance between any two of the three wells, and define the energy level W_0 via

$$W_0 := \min \left\{ W(p) : |p - p_\ell| = d_0, \ell = 1, 2, 3 \right\} > 0. \quad (4.21)$$

It follows from (4.15) and (4.16) that there must exist a point $x_j^+ \in \partial B_{\rho_j}$ such that $W(U(x_j^+)) = W_0$. Then for θ denoting the polar angle made with the positive x_1 -axis, we introduce the angle θ_j^+ via $x_j^+ = \rho_j e^{i\theta_j^+}$. We point out that necessarily

$$\theta_j^+ \rightarrow 0 \quad \text{as } j \rightarrow \infty, \quad (4.22)$$

since $W(U(\rho_j e^{i\theta})) \rightarrow 0$ at angles θ bounded away from zero in light of (4.15) and (4.16).

Then we introduce an arclength coordinate s along $\partial B_{\rho_j}^+$ with $s = 0$ corresponding to this x_j^+ via

$$s := \rho_j(\theta - \theta_j^+), \quad (4.23)$$

and define the continuous extension, say $\tilde{U}_j : (-\infty, \infty) \rightarrow \mathbb{R}^2$ of U along $\partial B_{\rho_j}^+$, expressed as a function of arclength variable (4.23), through the formula

$$\tilde{U}_j(s) = \begin{cases} p_1 & \text{for } s \geq \rho_j(\pi/2 - \theta_j^+) + 1 \\ \text{linear} & \text{for } \rho_j(\pi/2 - \theta_j^+) < s < \rho_j(\pi/2 - \theta_j^+) + 1 \\ U \left(\rho_j e^{i(\frac{s}{\rho_j} + \theta_j^+)} \right) & \text{for } \rho_j(-\pi/2 - \theta_j^+) \leq s \leq \rho_j(\pi/2 - \theta_j^+) \\ \text{linear} & \text{for } \rho_j(-\pi/2 - \theta_j^+) - 1 < s < \rho_j(\pi/2 - \theta_j^+) \\ p_2 & \text{for } s \leq \rho_j(-\pi/2 - \theta_j^+) - 1, \end{cases}$$

so that, in particular, we have

$$W(\tilde{U}_j(0)) = W_0. \quad (4.24)$$

Now again appealing to (4.15) and (4.16), we observe that the energy over the intervals of linear interpolation vanish in the limit; that is,

$$\lim_{j \rightarrow \infty} E \left(\tilde{U}_j, [\rho_j(\pi/2 - \theta_j^+), \rho_j(\pi/2 - \theta_j^+) + 1] \right) = 0,$$

with a similar result for the integral over the interval $[\rho_j(-\pi/2 - \theta_j^+) - 1, \rho_j(-\pi/2 - \theta_j^+)]$. Hence, through (4.19), we see that $\{\tilde{U}_j\}$ constitutes a minimizing sequence for (1.6). Invoking (4.24), it is then straightforward to establish that

$$\tilde{U}_j \rightarrow \zeta_{12} \quad \text{in } H_{\text{loc}}^1(\mathbb{R}) \text{ as } j \rightarrow \infty, \quad (4.25)$$

with $W(\zeta_{12}(0)) = W_0$ setting the particular translate of the heteroclinic ζ_{12} . Again we refer to [5] for details. (We recall here that we are assuming uniqueness of the three heteroclinic connections.) We then note that using (4.20), we can apply precisely the same argument along the left half-circle $\partial B_{\rho_j}^-$ to get H_{loc}^1 -convergence to ζ_{12} analogous to (4.25) there as well. In particular, in analogy with (4.22) and (4.24), we note that there exists an angle made with the negative x_1 -axis, which we denote by θ_j^- , that plays the same role as did θ_j^+ ; namely,

$$W \left(U \left(\rho_j e^{i(\pi - \theta_j^-)} \right) \right) = W_0 \quad \text{and} \quad \theta_j^- \rightarrow 0 \text{ as } j \rightarrow \infty. \quad (4.26)$$

Also, referring back to (4.15) and (4.16), with, say $\tau = 1/2$, it follows that

$$\left| U \left(\rho_j e^{i\theta} \right) - p_1 \right| \rightarrow 0 \text{ for } \frac{\pi}{6} \leq \theta \leq \frac{5}{6}\pi \quad \text{and} \quad \left| U \left(\rho_j e^{i\theta} \right) - p_2 \right| \rightarrow 0 \text{ for } -\frac{5}{6}\pi \leq \theta \leq -\frac{\pi}{6}. \quad (4.27)$$

3. Letting L_j denote the line passing through the two points $x_j^+ = \rho_j e^{i\theta_j^+}$ and $x_j^- := \rho_j e^{i(\pi-\theta_j^-)}$, we define the sequence $V_j : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ via

$$V_j(x) := \zeta_{12}(\text{dist}(x, L_j)). \quad (4.28)$$

Our goal in this step is to interpolate between U on ∂B_{ρ_j} and V_j on $\partial B_{\rho_{j-1}}$ so that the energetic cost in the annulus between these two circles is no greater than $2c_{12} + o(1)$. Again we will focus on the right half-annulus, with a similar calculation applying to the left half-annulus.

To this end, we first recall the exponential approach of $\zeta_{12} = \zeta_{12}(t)$ to p_1 for $t \gg 1$ and to p_2 for $t \ll -1$, cf. (3.44). Fixing any $\eta > 0$, it follows that we can find an interval, say $[-a_\eta, a_\eta]$, such that

$$E(\zeta_{12}, [-a_\eta, a_\eta]) \leq c_{12}, \quad (4.29)$$

and such that

$$|\zeta_{12}(t) - p_1| < \eta \text{ for } t > a_\eta, \quad \text{and} \quad |\zeta_{12}(t) - p_2| < \eta \text{ for } t < -a_\eta. \quad (4.30)$$

Then in view of the H^1 -convergence of \tilde{U}_j to ζ_{12} for $s \in [-a_\eta, a_\eta]$ guaranteed by (4.25), we can assert that for ρ_j large enough, one also has

$$c_{12} - \eta \leq E(\tilde{U}_j, [-a_\eta, a_\eta]) \leq c_{12} + \eta. \quad (4.31)$$

As consequence of (4.29) and (4.31), along with (1.6), (4.19) and (4.20), it also follows that

$$E(\zeta_{12}, \mathbb{R} \setminus [-a_\eta, a_\eta]) \leq \eta \quad (4.32)$$

and

$$E(\tilde{U}_j, \mathbb{R} \setminus [-a_\eta, a_\eta]) \leq \eta. \quad (4.33)$$

Of course, the analog of estimates (4.31) and (4.33) hold with the boundary values of U along $\partial B_{\rho_j}^+$ replaced by those along $\partial B_{\rho_j}^-$ as well.

Now we define the linear interpolation in the annulus $\{x : \rho_j - 1 \leq |x| \leq \rho_j\}$ with $\lambda_j(r) := r - \rho_j + 1$, via the formula

$$Z_j(x) := \lambda_j(|x|)U\left(\rho_j \frac{x}{|x|}\right) + (1 - \lambda_j(|x|))V_j(x). \quad (4.34)$$

We will divide up the energy of Z_j in the right half-annulus into two parts as follows:

$$\int \int_{\text{right half-annulus}} W(Z_j) + \frac{1}{2} |\nabla Z_j|^2 dx =$$

$$\begin{aligned}
& \int_{\rho_j-1}^{\rho_j} \int_{\{x \in \partial B_r^+ : \text{dist}(x, L_j) < a_\eta\}} \{ \cdot \} ds dr + \int_{\rho_j-1}^{\rho_j} \int_{\{x \in \partial B_r^+ : \text{dist}(x, L_j) \geq a_\eta\}} \{ \cdot \} ds dr \\
& =: I + II.
\end{aligned} \tag{4.35}$$

Regarding integral I , we note that for any r such that $\rho_j - 1 < r < \rho_j$ and any $x \in \partial B_r$ such that $\text{dist}(x, L_j) < a_\eta$, if we denote by $s(x)$ the arclength along ∂B_r from x to $\partial B_r \cap L_j$, then one easily checks that

$$0 \leq s(x) - \text{dist}(x, L_j) = O\left(\frac{1}{\rho_j}\right). \tag{4.36}$$

Combining this with the uniform convergence of \tilde{U}_j to ζ_{12} on $[-a_\eta, a_\eta]$ guaranteed by (4.25), we obtain that

$$\left| U\left(\rho_j \frac{x}{|x|}\right) - V_j(x) \right| \leq \left| U\left(\rho_j e^{i\left(\frac{s(x)}{\rho_j} + \theta_j^+\right)}\right) - \zeta_{12}(s(x)) \right| + |\zeta_{12}(s(x)) - \zeta_{12}(\text{dist}(x, L_j))| = o(1), \tag{4.37}$$

which then also implies that

$$Z_j(x) = \zeta_{12}(s(x)) + o(1). \tag{4.38}$$

Now

$$\left| \nabla U\left(\rho_j \frac{x}{|x|}\right) \right|^2 = \left| \tilde{U}_j'(s(x)) \right|^2,$$

and for x in the domain of integration of integral I , with an appeal to (4.36), we have

$$|\nabla V_j(x)|^2 = |\zeta'_{12}(\text{dist}(x, L_j))|^2 = |\zeta'_{12}(s)|^2 + o(1). \tag{4.39}$$

Therefore, since

$$\nabla Z_j(x) = \lambda_j(|x|) \nabla U\left(\rho_j \frac{x}{|x|}\right) + (1 - \lambda_j(|x|)) \nabla V_j(x) + \frac{x}{|x|} \otimes \left(U\left(\rho_j \frac{x}{|x|}\right) - V_j(x) \right), \tag{4.40}$$

we may compute that

$$\begin{aligned}
|\nabla Z_j(x)|^2 & \leq \left| \lambda_j(|x|) \nabla U\left(\rho_j \frac{x}{|x|}\right) + (1 - \lambda_j(|x|)) \nabla V_j(x) \right|^2 \\
& + C \left\{ \left| U\left(\rho_j \frac{x}{|x|}\right) - V_j(x) \right|^2 + \left| U\left(\rho_j \frac{x}{|x|}\right) - V_j(x) \right| \left(\left| \nabla U\left(\rho_j \frac{x}{|x|}\right) \right| + |\nabla V_j(x)| \right) \right\} \\
& \leq \lambda_j(|x|) \left| \nabla U\left(\rho_j \frac{x}{|x|}\right) \right|^2 + (1 - \lambda_j(|x|)) |\nabla V_j(x)|^2 + o(1),
\end{aligned} \tag{4.41}$$

where the last inequality follows from the convexity of $|\cdot|^2$, along with the use of (4.36) and (4.37), after noting that both $|\nabla U|$ and $|\nabla V_j|$ are uniformly bounded.

As a consequence of (4.38), (4.41) and another appeal to (4.25) we find that

$$\begin{aligned} & \int_{\rho_j-1}^{\rho_j} \int_{\{x \in \partial B_r^+ : \text{dist}(x, L_j) < a_\eta\}} W(Z_j(x)) + \frac{1}{2} |\nabla Z_j(x)|^2 ds dr \leq \\ & \int_{\rho_j-1}^{\rho_j} \int_{-a_\eta}^{a_\eta} W(\zeta_{12}(s)) + \frac{1}{2} \left\{ \lambda_j(r) |\tilde{U}_j'(s)|^2 + (1 - \lambda_j(r)) |\zeta'_{12}(s)|^2 \right\} ds dr + o(1) \\ & = \int_{\rho_j-1}^{\rho_j} \int_{-a_\eta}^{a_\eta} W(\zeta_{12}(s)) + \frac{1}{2} |\zeta'_{12}(s)|^2 ds dr + o(1). \end{aligned}$$

Hence, it follows from (4.29) that

$$I \leq c_{12} + \eta, \quad (4.42)$$

with a corresponding inequality holding for the energy of $\{Z_j\}$ over the region in the left portion of the annulus given by

$$\{x = (x_1, x_2) : \rho_j - 1 < |x| < \rho_j, \text{dist}(x, L_j) < \eta, x_1 < 0\}.$$

It remains to estimate integral II in (4.35). We will argue that this integral is $o(1)$ by relying on (4.32) and (4.33). We first claim that

$$\int_{\rho_j-1}^{\rho_j} \int_{\{x \in \partial B_r^+ : \text{dist}(x, L_j) \geq a_\eta\}} W(Z_j(x)) ds dr = O(\eta). \quad (4.43)$$

With an eye towards appealing to (3.26), we observe that necessarily for any $x \in \partial B_{\rho_j}$ such that $\text{dist}(x, L_j) \geq a_\eta$, one has either $|U(x) - p_1| < \beta$ or $|U(x) - p_2| < \beta$. Otherwise, the energetic cost incurred by transitioning along ∂B_{ρ_j} from $|U(x) - p_\ell| = \beta$ to $U(x) \approx p_\ell$ would be $O(1)$, violating (4.33), a transition that must occur in light of (4.27).

Since (4.30) implies that $V_j(x)$ also takes values in a region of convexity of W for $x \in \partial B_{\rho_j}$ such that $\text{dist}(x, L_j) \geq a_\eta$, we have that

$$W(Z_j(x)) \leq \lambda_j(|x|) W\left(U\left(\rho_j \frac{x}{|x|}\right)\right) + (1 - \lambda_j(|x|)) W(V_j(x)) \quad \text{for such an } x. \quad (4.44)$$

Hence, with an appeal to (4.32) and (4.33), we obtain claim (4.43).

To show that integral II in (4.35) is small, we still must estimate the integral of $|\nabla Z_j|^2$. Here we note from (4.40) that

$$|\nabla Z_j|^2 \leq C \left\{ \left[\lambda_j(|x|) \nabla U\left(\rho_j \frac{x}{|x|}\right) + (1 - \lambda_j(|x|)) \nabla V_j(x) \right]^2 + \left| U\left(\rho_j \frac{x}{|x|}\right) - V_j(x) \right|^2 \right\}$$

$$\leq C \left\{ \lambda_j(|x|) \left| \nabla U \left(\rho_j \frac{x}{|x|} \right) \right|^2 + (1 - \lambda_j(|x|)) |\nabla V_j(x)|^2 + \left| U \left(\rho_j \frac{x}{|x|} \right) - V_j(x) \right|^2 \right\}. \quad (4.45)$$

Integrating this expression over the set

$$\{x = (x_1, x_2) : \rho_j - 1 < |x| < \rho_j, \text{dist}(x, L_j) \geq a_\eta, x_1 > 0\},$$

we can use (4.32) and (4.33) once again to show that the first two terms in (4.45) integrate to $O(\eta)$.

The third term can be handled in the same manner as was done for the integral of $W(Z_j)$. That is, we split the integral into the set where $U \left(\rho_j \frac{x}{|x|} \right)$ is far from both p_1 and p_2 and where it is close to one of these wells. We know from (4.30) that $V_j(x)$ is near p_1 or p_2 for this domain of integration and therefore the measure of the set where $U \left(\rho_j \frac{x}{|x|} \right)$ is far from the wells must be small in order not to violate (4.33). Then, on any set where it is near to p_1 or p_2 , we have

$$\left| U \left(\rho_j \frac{x}{|x|} \right) - V_j(x) \right|^2 \leq 2 \left(\left| U \left(\rho_j \frac{x}{|x|} \right) - p_\ell \right|^2 + |V_j(x) - p_\ell|^2 \right) \quad \text{for either } \ell = 1 \text{ or } 2,$$

and so the quantity $\left| U \left(\rho_j \frac{x}{|x|} \right) - V_j(x) \right|^2$ is controlled by the sum of the integrals of $W \left(U \left(\rho_j \frac{x}{|x|} \right) \right)$ and $W(V_j(x))$. Hence, by (4.32) and (4.33) it must also integrate to $O(\eta)$.

Since the analysis of I leading to (4.42) holds for any $\eta > 0$ provided ρ_j is sufficiently large, as does this just completed analysis of integral II , we finally conclude that the interpolating sequence $\{Z_j\}$ satisfies the bound

$$\int_{\{x: \rho_j - 1 < |x| < \rho_j, x_1 > 0\}} W(Z_j) + \frac{1}{2} |\nabla Z_j|^2 dx \leq c_{12} + o(1).$$

The argument leading to the same estimate for the energy of Z_j taken over the set $\{x : \rho_j - 1 < |x| < \rho_j, x_1 < 0\}$ is identical, and so we arrive at the estimate

$$E(Z_j, B_{\rho_j} \setminus B_{\rho_{j-1}}) \leq 2c_{12} + o(1). \quad (4.46)$$

4. Having interpolated between the boundary values of U on ∂B_{ρ_j} and those of V_j on $\partial B_{\rho_{j-1}}$ with a cost bounded as in (4.46), we can now appeal to the local minimality of U to assert that

$$E(U, B_{\rho_j}) \leq 2c_{12} + E(V_j, B_{\rho_{j-1}}) + o(1),$$

where we recall that V_j is defined through (4.28). If we now consider a coordinate system (z_1, z_2) with z_1 -axis coinciding with the line L_j , z_2 -axis orthogonal to it, and with origin at the midpoint of the line segment $L_j \cap B_{\rho_j}$, then it follows immediately from the definition of ζ_{12} and (4.28) that

$$E(U, B_{\rho_j}) \leq 2c_{12} + \int_{-\frac{1}{2}\mathcal{H}^1(L_j \cap B_{\rho_j-1})}^{\frac{1}{2}\mathcal{H}^1(L_j \cap B_{\rho_j-1})} E(\zeta_{12}, \mathbb{R}) dz_1 + o(1) = c_{12} \mathcal{H}^1(L_j \cap B_{\rho_j}) + o(1). \quad (4.47)$$

Here we are using the fact that

$$\mathcal{H}^1(L_j \cap \{x : \rho_j - 1 < |x| < \rho_j\}) \leq 2 + o(1),$$

since we recall that the line L_j meets ∂B_{ρ_j} at the points $x_j^+ = \rho_j e^{i\theta_j^+}$ and $x_j^- = \rho_j e^{i(\pi-\theta_j^-)}$ with θ_j^+ and θ_j^- both approaching zero by (4.22) and (4.26). See Figure 5.

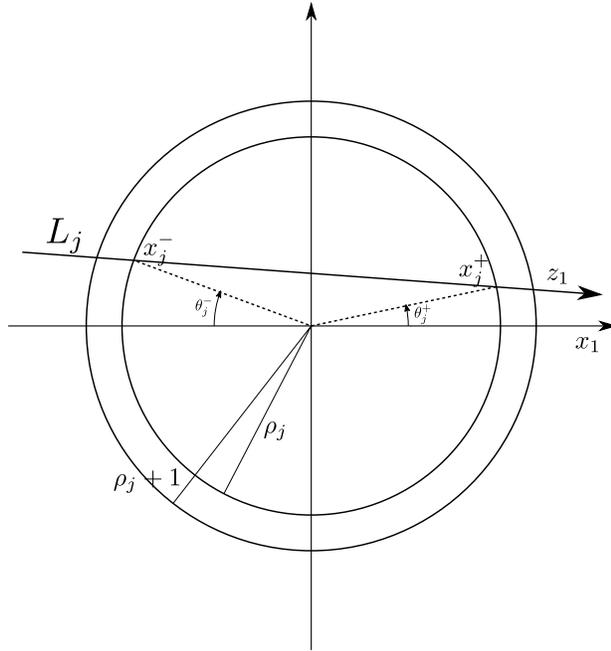


Figure 5: The coordinate system based on the line L_j passing through the points $x_j^+ = \rho_j e^{i\theta_j^+}$ and $x_j^- = \rho_j e^{i(\pi-\theta_j^-)}$ defined in (4.22) and (4.26).

5. We conclude the proof of Proposition 4.3 with a lower bound for the energy of U on B_{ρ_j} that will contradict the upper bound (4.47), thus eliminating the

possibility that (4.11) can occur. To this end, we introduce an interpolation, say \tilde{Z}_j , on the annulus $\{x : \rho_j \leq |x| \leq \rho_j + 1\}$ between the values of U on ∂B_{ρ_j} and those of V_j given by (4.28) on ∂B_{ρ_j+1} . The formula for the sequence $\{\tilde{Z}_j\}$ is identical to that of $\{Z_j\}$ given in (4.34), with the exception that $\lambda_j(r)$ is now replaced by $\tilde{\lambda}_j(r) := \rho_j + 1 - r$. Then the entire argument that led to the energy bound (4.46) applies equally well to $\{\tilde{Z}_j\}$, to establish that

$$E\left(\tilde{Z}_j, B_{\rho_j+1} \setminus B_{\rho_j}\right) \leq 2c_{12} + o(1). \quad (4.48)$$

As a consequence of (4.48), if we now define the sequence $\{\mathcal{U}_j\}$ on B_{ρ_j+1} via

$$\mathcal{U}_j(x) := \begin{cases} U(x) & \text{for } x \in B_{\rho_j}, \\ \tilde{Z}_j(x) & \text{for } x \in B_{\rho_j+1} \setminus B_{\rho_j}, \end{cases}$$

then we have the lower energy bound

$$E(U, B_{\rho_j}) \geq E(\mathcal{U}_j, B_{\rho_j+1}) - 2c_{12} - o(1). \quad (4.49)$$

Next we will compute a lower bound for the integral on the right using a coordinate system (z_1, z_2) very similar to the one introduced above (4.47), where again z_1 is the arc length on the line L_j but now with $z_1 = 0, z_2 = 0$ corresponding to the midpoint of $L_j \cap B_{\rho_j+1}$. For any $z_1 \in \mathbb{R}$ the set $\{z_2 \mid (z_1, z_2) \in B_{\rho_j+1}\}$ is an interval that we denote $(a_j(z_1), b_j(z_1))$.

The line L_j is at a distance $\delta_j = \rho_j \left| \sin\left(\frac{\theta_j^- + \theta_j^+}{2}\right) \right|$ from the origin. Note that $\delta_j \ll \rho_j$ since $|\theta_j^+| + |\theta_j^-| \rightarrow 0$. It follows that B_{ρ_j+1} contains the set of points with coordinates (z_1, z_2) such that

$$-\ell_j < z_1 < \ell_j, \quad -h_j(z_1) < z_2 < h_j(z_1), \quad (4.50)$$

where

$$\begin{aligned} \ell_j &= \frac{1}{2} \mathcal{H}^1(L_j \cap B_{\rho_j+1}) = ((\rho_j + 1)^2 - \delta_j^2)^{1/2} \\ \text{and } h_j(z_1) &= ((\rho_j + 1)^2 - z_1^2)^{1/2} - \delta_j. \end{aligned} \quad (4.51)$$

In particular,

$$h_j(z_1) \leq \min(|a_j(z_1)|, b_j(z_1)). \quad (4.52)$$

Then we can write

$$E(\mathcal{U}_j, B_{\rho_j+1}) \geq \int_{B_{\rho_j+1}} \left| \frac{\partial \mathcal{U}_j}{\partial z} \right|^2 + \int_{-\ell_j}^{\ell_j} I(z_1) dz_1, \quad (4.53)$$

where

$$I(z_1) = E(\mathcal{U}_j(z_1, \cdot), [a_j(z_1), b_j(z_1)]).$$

From the construction of \mathcal{U}_j using \tilde{Z}_j , we have that $\mathcal{U}_j(z_1, z_2) = \zeta_{12}(z_2)$ if $(z_1, z_2) \in \partial B_{\rho_j+1}$, that is if $z_2 = a_j(z_1)$ or $z_2 = b_j(z_1)$. Thus we may extend $\mathcal{U}_j(z_1, \cdot)$ continuously by setting $\mathcal{U}_j(z_1, z_2) = \zeta_{12}(z_2)$ if $z_2 \notin (a_j(z_1), b_j(z_1))$.

Using the minimizing property of ζ_{12} with respect to its boundary conditions on any interval $[a, b]$ and its exponential decay as in (3.44), we have for any $z_1 \in (-\ell_j, \ell_j)$ that

$$\begin{aligned} I(z_1) &= E(\mathcal{U}_j(z_1, \cdot), \mathbb{R}) - E(\mathcal{U}_j(z_1, \cdot), \mathbb{R} \setminus [a_j(z_1), b_j(z_1)]) \\ &= E(\mathcal{U}_j(z_1, \cdot), \mathbb{R}) - E(\zeta_{12}, \mathbb{R} \setminus [a_j(z_1), b_j(z_1)]) \\ &\geq E(\zeta_{12}, \mathbb{R}) - E(\zeta_{12}, \mathbb{R} \setminus [-h_j(z_1), h_j(z_1)]) \geq c_{12} - Ce^{-ch_j(z_1)}, \end{aligned} \tag{4.54}$$

where $c, C > 0$ depend only on W .

We integrate (4.54) over $z_1 \in (-\ell_j, \ell_j)$. In view of (4.53) we find that

$$E(\mathcal{U}_j, B_{\rho_j+1}) \geq \int_{B_{\rho_j+1}} \left| \frac{\partial \mathcal{U}_j}{\partial z_1} \right|^2 + 2c_{12}\ell_j - C \int_{-\ell_j}^{\ell_j} e^{-ch_j(z_1)} dz_1.$$

It is straightforward to show, using (4.51) and the fact that $\delta_j \ll \rho_j$ as $\rho_j \rightarrow +\infty$, that the last integral above is $o(1)$ and then, since

$$2\ell_j = \mathcal{H}^1(L_j \cap B_{\rho_j+1}) = \mathcal{H}^1(L_j \cap B_{\rho_j}) + 2 + o(1),$$

we conclude that

$$E(\mathcal{U}_j, B_{\rho_j+1}) \geq \int_{B_{\rho_j+1}} \left| \frac{\partial \mathcal{U}_j}{\partial z_1} \right|^2 + c_{12} \mathcal{H}^1(L_j \cap B_{\rho_j}) + 2c_{12} + o(1),$$

In view of (4.49) and the upper bound (4.47), we deduce that

$$\int_{B_{\rho_j+1}} \left| \frac{\partial \mathcal{U}_j}{\partial z_1} \right|^2 = o(1). \tag{4.55}$$

6. We may now conclude. Going back to the original coordinates (x_1, x_2) we have that $U = \mathcal{U}_j \circ \varphi_j$ on B_{ρ_j} , where φ_j denotes the $(x_1, x_2) \rightarrow (z_1, z_2)$ map. But, as $j \rightarrow +\infty$, the rotational component of φ_j converges to the identity since θ_j^+ and θ_j^- both tend to 0. Passing to the limit in (4.55) we thus deduce that $U(x_1, x_2)$ does not depend on x_1 .

Then, as a function of x_2 only, U is a minimizing solution on \mathbb{R} , which converges to p_1 as $x_2 \rightarrow -\infty$ and to p_2 as $x_2 \rightarrow +\infty$. Thus there exists $\Delta \in \mathbb{R}$ such that $U(x_1, x_2) = \zeta_{12}(x_2 + \Delta)$. The proof of Theorem 1.3 is complete. \square

5 Two partitioning problems: Regularity and comparison

Our proof of the key estimate (3.15) in Proposition 3.3 relies upon the regularity theory for two partitioning problems, one fairly standard and the other perhaps not so standard. In this section we state the regularity theory for these problems, as established in [24]. Then we state and prove a result relating the infima of these two problems.

For the convenience of the reader, we restate these two partitioning problems here:

Problem 1. Fix a function $h \in BV(\partial B; P)$. For any disjoint sets S_1, S_2 and S_3 of finite perimeter in B such that $|B \setminus \cup S_\ell| = 0$ define

$$m_0 := \inf \{ E_0(S_1, S_2, S_3) : \cup S_\ell = B, \partial S_\ell \cap \partial B = h^{-1}(p_\ell) \text{ for } \ell = 1, 2, 3 \},$$

where E_0 is given by

$$E_0(S_1, S_2, S_3) = t_1 \mathcal{H}^1(\partial^* S_1 \cap B) + t_2 \mathcal{H}^1(\partial^* S_2 \cap B) + t_3 \mathcal{H}^1(\partial^* S_3 \cap B).$$

Problem 2. Fix a number $\delta > 0$. Then for h and E_0 as above, and disjoint sets S_1, S_2 and S_3 of finite perimeter in B define

$$m_0^\delta := \inf \{ E_0(S_1, S_2, S_3) : |B \setminus \cup S_\ell| \leq \delta, \partial S_\ell \cap \partial B = h^{-1}(p_\ell) \text{ for } \ell = 1, 2, 3 \}.$$

Regarding Problem 1, the regularity theory of minimizing planar partitions subject to volume constraints on each phase is developed, for instance, in [20, 22]. For our purposes, however, we require a version valid without volume constraints but subject to a Dirichlet condition, and which has additional properties specific to minimization within a ball. For this we quote the recent work in [24].

Theorem 5.1. (*[24], Thm. 1.6*) *If (S_1^0, S_2^0, S_3^0) be a minimizer of Problem 1, then every connected component of $\partial S_\ell^0 \cap \partial S_m^0 \cap B$ is a line segment terminating at an interior triple junction $x \in \partial S_1^0 \cap \partial S_2^0 \cap \partial S_3^0 \cap B$, at $x \in \partial S_\ell^0 \cap \partial S_m^0 \cap \partial B$ for $\ell \neq m$ which is a point of discontinuity of h , or at a boundary triple junction $x \in \partial S_1^0 \cap \partial S_2^0 \cap S_3^0 \cap \partial B$ which is a point of discontinuity of h . Moreover, there exists angles α_ℓ , $\ell = 1, 2, 3$, satisfying (1.9) such that if $x \in B$ is an interior triple junction, for some $r_x > 0$, $S_\ell^0 \cap B_{r_x}$ for $\ell = 1, 2, 3$ are circular sectors determined by α_ℓ . Finally, every connected component C of S_ℓ^0 is convex and meets ∂B along one or more arcs of $h^{-1}(p_\ell)$.*

The last property, namely that every connected component C of S_ℓ^0 is convex and meets ∂B along one or more arcs of $h^{-1}(p_\ell)$ is not stated in [24], Thm. 1.6, but is immediate since any island of phase could be filled in with a different phase, thereby lowering the total perimeter without disrupting the boundary condition.

Regarding Problem 1, we will also need the following corollary, which follows easily from Theorem 5.1.

Corollary 5.2. *Suppose $h \in BV(\partial B; P)$ has k jump discontinuities for some non-negative integer k . Then there exists an integer $N(k)$ such that the total number of triple junctions appearing in any minimizer of Problem 1 cannot exceed $N(k)$.*

Proof. Denote any minimizer of Problem 1 by u_0 . Since any connected component of $\{u_0 = p_\ell\}$ must meet ∂B along one or more connected components of the set of boundary arcs $h^{-1}(p_\ell)$, and since the number of these arcs is necessarily bounded by k , we can conclude that for $\ell = 1, 2$ and 3 , the number of connected components of $\{u_0 = p_\ell\}$ cannot exceed k as well. Next we note that in light of the convexity of every component of $\{u_0 = p_\ell\}$, any triple of components, one each from $\{u_0 = p_\ell\}$ for $\ell = 1, 2, 3$, can only meet at a triple junction at most once. Counting up all the possible triples, it follows that the number of triple junctions of a minimizer u_0 cannot exceed k^3 . \square

The regularity theory for Problem 2 is more subtle since minimizers will typically develop cusps to replace the triple junctions appearing in the solution of Problem 1, a phenomenon referred to in some literature as a “wetting” of the singularities, see e.g. [12]. See Figure 6. Here we quote the following result:

Theorem 5.3. *([24], Thm. 1.4) Let $(S_1^\delta, S_2^\delta, S_3^\delta)$ be a minimizer of Problem 2, and denote by $G^\delta := B \setminus \cup_{\ell=1}^3 S_\ell^\delta$. Then every connected component of $\partial S_\ell^\delta \cap \partial S_j^\delta \cap B$ is a line segment terminating either on ∂B at a point of discontinuity of h between p_ℓ and p_j or at a point in $\partial S_\ell^\delta \cap \partial S_j^\delta \cap \partial G^\delta \cap \overline{B}$. Referring to those points in $\partial S_\ell^\delta \cap \partial S_m^\delta \cap \partial G^\delta \cap B$ and $\partial S_\ell^\delta \cap \partial S_m^\delta \cap \partial G^\delta \cap \partial B$ as cusp and corner points, respectively, there exist positive κ_ℓ^δ for $\ell = 1, 2, 3$ such that*

$$t_1 \kappa_1^\delta = t_2 \kappa_2^\delta = t_3 \kappa_3^\delta \quad (5.1)$$

and, for $\ell = 1, 2, 3$, $\partial S_\ell^\delta \cap \partial G^\delta$ consists of a union of circular arcs of curvature κ_ℓ^δ , each of whose two endpoints are either a cusp point in B or a corner point in ∂B at a point of discontinuity of h . Furthermore, at cusp points, $\partial S_\ell^\delta \cap \partial G^\delta$ and $\partial S_m^\delta \cap \partial G^\delta$ meet $\partial S_\ell^\delta \cap \partial S_m^\delta$ tangentially. Finally, any connected component C of S_ℓ^δ is convex.

Since any admissible partition of B for Problem 1 is also admissible for Problem 2, it is obvious that $m_0^\delta \leq m_0$. However, an inequality in the reverse direction also holds.

Theorem 5.4. *For any positive integer k , let h be any function in $BV(\partial B; P)$ having no more than k discontinuities. Then for any $\delta > 0$, the infimum m_0 for Problem 1 and the infimum m_0^δ for Problem 2 are related via*

$$m_0^\delta \geq m_0 - \gamma(k) \delta^{1/2}, \quad (5.2)$$

for some constant $\gamma(k)$.

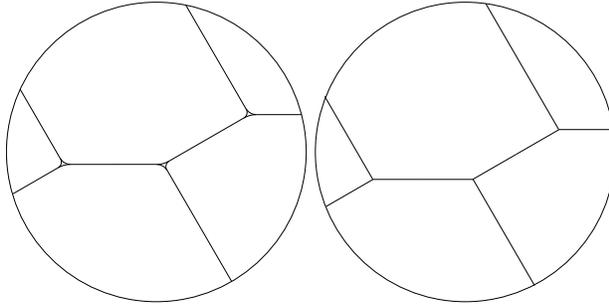


Figure 6: Minimizers of E_0 subject to the same Dirichlet condition for Problem 1 (right) and Problem 2 (left).

Proof of Theorem 5.4. We will view minimization of E_0 subject to the given Dirichlet condition on ∂B and subject to the constraint (3.19) as a problem of coloring most of B with three colors, where for a minimizer $\{S_1^\delta, S_2^\delta, S_3^\delta\}$ we rename S_1^δ as the yellow set Y^δ , S_2^δ as the red set R^δ and S_3^δ as the blue set B^δ . Then the region in B not covered by these sets, namely G^δ , will be referred to as the gray set. Referring back to the formulation (3.17) of the partitioning energy E_0 , let us now write it as

$$E_0(Y^\delta, R^\delta, B^\delta) = c_Y \mathcal{H}^1(\partial Y^\delta \cap B) + c_R \mathcal{H}^1(\partial R^\delta \cap B) + c_B \mathcal{H}^1(\partial B^\delta \cap B), \quad (5.3)$$

where we have changed notation to let $c_Y = t_1$ = the cost of ‘yellow boundary’, $c_R = t_2$ and $c_B = t_3$.

Having fixed the boundary data $h \in BV(\partial B; P)$, we have $\partial Y^\delta \cap \partial B = h^{-1}(p_1)$, $\partial R^\delta \cap \partial B = h^{-1}(p_2)$ and $\partial B^\delta \cap \partial B = h^{-1}(p_3)$, and so ∂B is partitioned into a finite number circular arcs, some yellow, some red and some blue, though we make no assumption that necessarily all three colors are present in the

boundary data. We recall that we are assuming the total number of these arcs does not exceed k . It also follows from the Dirichlet condition that ∂G^δ meets ∂B only at at most k isolated points, if at all.

The main step in the proof consists of arguing that one can bound the number of components of G^δ by a constant depending only on k . Once this is established, the bound (5.2) will follow rather easily.

At this point, we will assume, with no loss of generality, that

$$c_Y \leq \min\{c_R, c_B\}. \quad (5.4)$$

We now proceed with the proof in four steps.

1. We first claim that with no loss of generality we may assume every component of R^δ and every component of B^δ must meet ∂B . That is, there are no islands of red or blue in the interior of B . This follows since any such island could be changed to yellow, either resulting in a new minimizer in the case of equality in (5.4), or else contradicting the minimality of $\{Y^\delta, R^\delta, B^\delta\}$ in the case of strict inequality in (5.4). As a consequence, the number of connected components of R^δ and of B^δ cannot exceed the number of red and blue boundary components dictated by h . In particular, both numbers are bounded by k . For the remainder of the argument, we will denote these components via

$$R^\delta = R_1^\delta \cup R_2^\delta \cup \dots \cup R_{k_1}^\delta \quad \text{and} \quad B^\delta = B_1^\delta \cup B_2^\delta \cup \dots \cup B_{k_2}^\delta \quad (5.5)$$

for some integers $k_1 = k_1(\delta)$ and $k_2 = k_2(\delta)$ such that $k_1 + k_2 \leq k$.

2. Next we claim that we may assume every component of G^δ is simply connected. This follows since by Step 1, any non-simply connected component of G^δ would have one or more components of Y^δ consisting of full disks lying in its interior. We observe as a consequence of Proposition 5.3 that the outer boundary component of any component of G^δ consists of a union of circular arcs of curvature $\kappa_Y^\delta, \kappa_R^\delta$ or κ_B^δ bowing into G^δ , all meeting tangentially at cusp points, where we have renamed κ_1^δ as κ_Y^δ , etc. Then we may shift any interior yellow disk until it touches this outer boundary component at two points without changing the total value of E_0 , that is, creating a new minimizer. The only obstruction to sliding such an interior yellow disk over to the boundary would be that it first hits another yellow disk, but clearly two yellow disks, tangent at a point, is a non-minimizing configuration and so could not occur. In this manner, any minimizing configuration possessing a non-simply connected component of G^δ could be replaced by another minimizer having more components of G^δ than the original, but for which every component of the new G^δ is simply connected.

3. Our next goal is to bound the number of components of Y^δ which touch G^δ , in the sense that their boundaries have nonempty intersection. It suffices to bound the number of those components which do not touch the boundary ∂B since there are at most k components which touch ∂B .

Let then Y_ℓ^δ be a yellow component which touches G^δ . As described in Theorem 5.3, the boundary of Y_ℓ^δ is C^1 and consists of circular arcs and segments separated by points which are cusp singularities of the partition. Moreover, we may assume that there are at least two such cusp points, for otherwise ∂Y_ℓ^δ minus at most a point would be in a gray component and by sliding Y_ℓ^δ in this gray component, we would obtain a minimizing partition where Y_ℓ^δ has at least two cusp points on its boundary.

Let then p_1 and p_2 be cusp points on ∂Y_ℓ^δ , separated by a circular arc $\gamma \subset \partial Y_\ell^\delta \cap G^\delta$ of radius $r = 1/\kappa_Y^\delta$. These points also belong to the boundaries of red or blue components drawn from the collection (5.5); call these two components A_1 and A_2 . We claim that given components A_1 and A_2 , each either red or blue, there can be at most two yellow components separated from the gray area by a boundary arc γ whose endpoints belong to ∂A_1 and ∂A_2 , respectively.

We will argue this by first noting that the completion of any such circular arc γ into a full circle yields a circle of radius $r = 1/\kappa_Y^\delta$. By Theorem 5.3, this circle necessarily meets both ∂A_1 and ∂A_2 tangentially. We will thus rule out the possibility of there existing three or more such arcs γ by showing that there can never exist three balls of the same radius, all exterior to A_1 and A_2 , with all three meeting both ∂A_1 and ∂A_2 tangentially, unless all three balls have collinear centers. This is an elementary property of convex sets, but not being aware of a reference, we provide a proof.

We note that necessarily the center of any such ball must be equidistant from ∂A_1 and ∂A_2 . Therefore, for each $t > 0$, we consider the possible intersections of the curves $\Gamma_1^t := \{x : \text{dist}(x, \partial A_1) = t\}$ and $\Gamma_2^t = \{x : \text{dist}(x, \partial A_2) = t\}$. Since A_1 and A_2 are two disjoint, convex sets that have C^1 boundaries within B , it follows that for all $t > 0$, Γ_1^t and Γ_2^t are convex, closed curves that are also C^1 within B , and furthermore, $\Gamma_1^t \cap \Gamma_2^t$ are disjoint for t small.

Now consider the first time $t_1 > 0$ when $\Gamma_1^{t_1}$ meets $\Gamma_2^{t_1}$. This could happen along a line segment, since we recall that neither curve is necessarily strictly convex. However, for $\Gamma_1^{t_1} \cap \Gamma_2^{t_1}$ to consist of a line segment would mean that ∂A_1 and ∂A_2 must have boundary components that are parallel line segments. In this case, of course there exists a one-parameter family of circles with this tangency property, but necessarily, their centers all lie on the line segment $\Gamma_1^{t_1} \cap \Gamma_2^{t_1}$; that is, they are collinear.

The other possibility is that at time t_1 , the intersection $\Gamma_1^{t_1} \cap \Gamma_2^{t_1}$ consists of one point. Then for $0 < t - t_1 \ll 1$, the convexity of both curves means that

the intersection will consist of two points, which in the context of our yellow components, allows for the possibility that two curves representing boundary components of two distinct elements of Y^δ , say, $\partial Y_{\ell_1}^\delta \cap G^\delta$ and $\partial Y_{\ell_2}^\delta \cap G^\delta$, both meet ∂A_1 and ∂A_2 . This could certainly happen. See Figure 7.

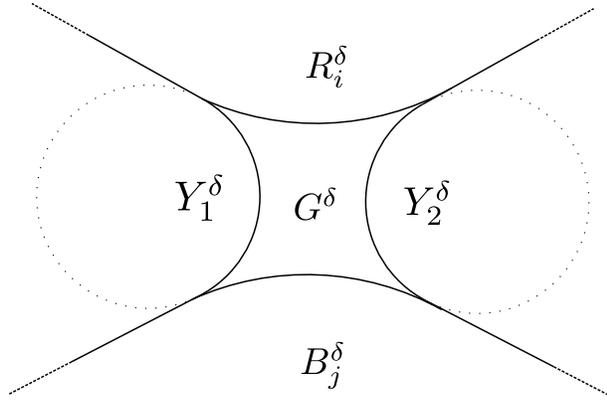


Figure 7: A configuration with two yellow components, both having a circular boundary arc whose endpoints meet the same pair of components from the collection in (5.5), resulting in 4 cusp points.

However, we claim that as t increases there cannot be more than two intersections of the two curves. We argue by contradiction and suppose that there exists a first time $t_2 > t_1$ where $\Gamma_1^{t_2} \cap \Gamma_2^{t_2}$ consists of three points. Two of these points represent the continuous evolution of the original two points that emerged as t passed through $t = t_1$ but at the third point, say $x \in \Gamma_1^{t_2} \cap \Gamma_2^{t_2}$, it must be the case that $\Gamma_1^{t_2}$ and $\Gamma_2^{t_2}$ meet tangentially, this being the first time a third point of intersection emerges. Denoting by \mathcal{L} the line of tangency, it follows from convexity that either both $\{x : \text{dist}(x, \partial A_1) < t_2\}$ and $\{x : \text{dist}(x, \partial A_2) < t_2\}$ lie on the same side of \mathcal{L} or they lie on opposite sides. But if they both lie on the same side then tracing back from x a distance t_2 along the common inner normal to $\Gamma_1^{t_2}$ and $\Gamma_2^{t_2}$, one would arrive at a point in common to ∂A_1 and ∂A_2 , which is impossible given that they are disjoint. If instead one supposes the two sets lie on opposite sides of \mathcal{L} , then that all earlier times, it must have been that $\Gamma_1^t \cap \Gamma_2^t$ was empty, contradicting the fact that the two curves met at the earlier time t_1 .

Returning to the possibility of boundary components of ∂A_1 and ∂A_2 consisting of two parallel line segments, we note that any line segment on the boundary of a component corresponds to one and only one common boundary with a component of a different color, so in this context, it would correspond to only one yellow component meeting A_1 , A_2 and G^δ . This proves the claim that given

components A_1 and A_2 , each either red or blue, there can be at most two yellow components separated from the gray area by a boundary arc γ whose endpoints belong to ∂A_1 and ∂A_2 , respectively. This in turn proves that the number of yellow components touching G^δ but not touching ∂B is at most twice the number of pairs of components of red or blue, chosen from (5.5), which is bounded by $2\binom{k}{2}$.

4. Now we turn to the task of bounding the number of components of G^δ . We will accomplish this by bounding the total number of cusp points. Any component of G^δ has boundary consisting of a union of circular arcs of curvature $\kappa_Y^\delta, \kappa_R^\delta$ or κ_B^δ bowing into G^δ , all meeting tangentially at cusp points. Each of these circular boundary arcs must be a portion of boundary drawn from the collection of sets

$$R_i^\delta \text{ for } i \in \{1, 2, \dots, k_1\}, \quad B_j^\delta \text{ for } j \in \{1, 2, \dots, k_2\}, \text{ or } Y_\ell^\delta \text{ for } \ell \in \{1, 2, \dots, k_3\}, \quad (5.6)$$

with $k_1 + k_2 + k_3 \leq k + 2\binom{k}{2}$ in light of (1) and (3) above.

Now consider any yellow/red, yellow/blue or red/blue pair taken from this collection, say for instance R_1^δ and B_1^δ , and suppose ∂R_1^δ and ∂B_1^δ each have circular arcs bordering the same component of G^δ such that these two arcs meet at a particular cusp point. By the convexity of red and blue components, $\partial R_1^\delta \cap \partial B_1^\delta$ must consist solely of a line segment one of whose endpoints is this cusp point. It follows that ∂R_1^δ and ∂B_1^δ can both meet the boundary of some other gray component at at most one other cusp point, namely at a cusp point sitting at the other endpoint of their one common boundary segment. The same argument could be made between any pairing of a Y_ℓ^δ with any R_i^δ or B_j^δ , provided Y_ℓ^δ is not a full disk. On the other hand, if Y_ℓ^δ is a full disk then its intersection with any R_i^δ or B_j^δ results in only one cusp point due to the convexity of both sets involved. Estimating crudely, the total number of yellow/red, yellow/blue or red/blue pairs drawn from the collection (5.6) is bounded by

$$\binom{k + 2\binom{k}{2}}{2}.$$

Hence, as just argued, the total number of cusp points in a minimizing configuration $\{Y^\delta, R^\delta, B^\delta\}$ cannot exceed twice this number. But since a closed curve comprised of concave circular arcs requires at least three such arcs, it follows that any component of ∂G^δ must have at least three cusp points. Thus we can bound the total number of gray boundary components and hence, the total number of gray components, by

$$C(k) := \frac{2}{3} \binom{k + 2\binom{k}{2}}{2}. \quad (5.7)$$

5. Finally, we are ready to establish inequality (5.2). To this end, we now build out of $\{Y^\delta, R^\delta, B^\delta\}$ a competitor in the problem, denoted by Problem 1 at the outset of this section, of minimizing E_0 among full partitions of the disk B , subject to the Dirichlet condition h , by defining $\tilde{Y}^\delta := Y^\delta \cup G^\delta$. Then $\{\tilde{Y}^\delta, R^\delta, B^\delta\}$ competes with the minimizer of this problem, denoted in the statement of Theorem 5.4 by $\{S_1^0, S_2^0, S_3^0\}$, and so we have

$$E_0(\tilde{Y}^\delta, R^\delta, B^\delta) \geq E_0(S_1^0, S_2^0, S_3^0), \quad (5.8)$$

as well as

$$E_0(\tilde{Y}^\delta, R^\delta, B^\delta) \leq E_0(Y^\delta, R^\delta, B^\delta) + c_Y \mathcal{H}^1(\partial G^\delta). \quad (5.9)$$

Now in light of (5.1) and our bound on the total number of cusp points, any boundary of a component of G^δ consists of at most $2 \binom{k+2}{2} \binom{k}{2}$ circular arcs of radius at most

$$\max \left\{ \frac{1}{\kappa_R}, \frac{1}{\kappa_B} \right\} = \frac{\max\{c_R, c_B\}}{c_Y} \frac{1}{\kappa_Y^\delta}.$$

Consequently, bounding the length of any arc by the perimeter of the corresponding full circle, we can assert that

$$\mathcal{H}^1(\partial G^\delta) \leq 2\pi C(k) \frac{\max\{c_R, c_B\}}{c_Y} \frac{1}{\kappa_Y^\delta}. \quad (5.10)$$

Then turning from bounding perimeter to bounding area, we note that the area enclosed by the boundary of any gray component in this minimizing configuration is bounded from below by the area of the smallest possible (simply connected) gray component, namely the one formed by just three arcs arising from the tangential contact of one yellow, one red and one blue arc. This number could of course be computed precisely but for our purposes, it suffices to observe through another appeal to (5.1) that it is given by $\alpha \frac{1}{(\kappa_Y^\delta)^2}$ for some positive constant $\alpha = \alpha(c_R, c_B)$, where we have expressed this minimal area in terms of κ_Y^δ though of course, we could have expressed it in terms of either of the other two curvatures as well. Hence, assuming there exists at least one component of G^δ , we have the following estimate on the curvature κ_Y^δ :

$$\alpha \frac{1}{(\kappa_Y^\delta)^2} \leq |G^\delta| \leq \delta. \quad (5.11)$$

Combining (5.8)-(5.11), we conclude that

$$E_0(Y^\delta, R^\delta, B^\delta) \geq E_0(S_1^0, S_2^0, S_3^0) - \gamma(k) \delta^{1/2} \quad \text{with} \quad \gamma(k) := \frac{2\pi C(k)}{\sqrt{\alpha}} \max\{c_R, c_B\}.$$

□

Acknowledgments. P.S. would like to thank Mihai Ciucu, Francesco Maggi and Frank Morgan for helpful suggestions. The research of P.S. was supported by a Simons Collaboration grant 585520, an NSF grant DMS 2106516 and through a visit to Université Paris-Est Créteil made possible through the Bezout Chair program. He is grateful for the hospitality shown during this visit.

References

- [1] ALAMA, S., BRONSARD, L., AND GUI, C. Stationary layered solutions in \mathbf{R}^2 for an Allen-Cahn system with multiple well potential. *Calc. Var. Partial Differential Equations* 5, 4 (1997), 359–390.
- [2] ALIKAKOS, N. D. On the structure of phase transition maps for three or more coexisting phases. In *Geometric partial differential equations*, vol. 15 of *CRM Series*. Ed. Norm., Pisa, 2013, pp. 1–31.
- [3] ALIKAKOS, N. D., AND FUSCO, G. On the connection problem for potentials with several global minima. *Indiana Univ. Math. J.* 57, 4 (2008), 1871–1906.
- [4] ALIKAKOS, N. D., AND FUSCO, G. Entire solutions to equivariant elliptic systems with variational structure. *Arch. Ration. Mech. Anal.* 202, 2 (2011), 567–597.
- [5] ALIKAKOS, N. D., FUSCO, G., AND SMYRNELIS, P. *Elliptic systems of phase transition type*, vol. 91 of *Progress in Nonlinear Differential Equations and their Applications*. Birkhäuser/Springer, Cham, 2018.
- [6] ALIKAKOS, N. D., AND GENG, Z. On the triple junction problem on general surface tension coefficients. *arXiv 2304.13106* (2023).
- [7] ALIKAKOS, N. D., AND GENG, Z. On the triple junction problem on the plane without symmetry hypothesis. *arXiv 2210.17134* (2023).
- [8] BALDO, S. Minimal interface criterion for phase transitions in mixtures of Cahn-Hilliard fluids. *Ann. Inst. H. Poincaré Anal. Non Linéaire* 7, 2 (1990), 67–90.
- [9] BATES, P., FUSCO, G., AND SMYRNELIS, P. Entire solutions with six-fold junctions to elliptic gradient systems with triangle symmetry. *Adv. Nonlinear Stud.* 13, 1 (2013), 1–11.

- [10] BATES, P. W., FUSCO, G., AND SMYRNELIS, P. Multiphase solutions to the vector Allen-Cahn equation: crystalline and other complex symmetric structures. *Arch. Ration. Mech. Anal.* 225, 2 (2017), 685–715.
- [11] BETHUEL, F. Asymptotics for two-dimensional vectorial Allen-Cahn systems. to appear, *Acta Math.* (2023).
- [12] BRAKKE, K., AND MORGAN, F. Instability of the wet x soap film. *J. Geom. Anal.* 8, 5 (1998), 749–767.
- [13] BRONSARD, L., GUI, C., AND SCHATZMAN, M. A three-layered minimizer in \mathbf{R}^2 for a variational problem with a symmetric three-well potential. *Comm. Pure Appl. Math.* 49, 7 (1996), 677–715.
- [14] FONSECA, I., AND TARTAR, L. The gradient theory of phase transitions for systems with two potential wells. *Proc. Roy. Soc. Edinburgh Sect. A* 111, 1-2 (1989), 89–102.
- [15] FUSCO, G. Minimizing under relaxed symmetry constraints: Triple and N-junctions. *Annali Scuola Normale Superiore, classe di Scienze* (2022).
- [16] GAZOULIS, D. On the γ -convergence of the Allen-Cahn functional with boundary conditions. *arXiv 2301:07458* (2023).
- [17] GIUSTI, E. *Minimal surfaces and functions of bounded variation*, vol. 80 of *Monographs in Mathematics*. Birkhäuser Verlag, Basel, 1984.
- [18] GUI, C., AND SCHATZMAN, M. Symmetric quadruple phase transitions. *Indiana Univ. Math. J.* 57, 2 (2008), 781–836.
- [19] KOHN, R. V., AND STERNBERG, P. Local minimisers and singular perturbations. *Proc. Roy. Soc. Edinburgh Sect. A* 111, 1-2 (1989), 69–84.
- [20] MAGGI, F. *Sets of finite perimeter and geometric variational problems*, vol. 135 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 2012. An introduction to geometric measure theory.
- [21] MONTEIL, A., AND SANTAMBROGIO, F. Metric methods for heteroclinic connections in infinite-dimensional spaces. *Indiana Univ. Math. J.* 69, 4 (2020), 1445–1503.
- [22] MORGAN, F. Immiscible fluid clusters in \mathbf{R}^2 and \mathbf{R}^3 . *Michigan Math. J.* 45, 3 (1998), 441–450.

- [23] NOVACK, M. R. Dimension reduction for the Landau–de Gennes model: the vanishing nematic correlation length limit. *SIAM J. Math. Anal.* *50*, 6 (2018), 6007–6048.
- [24] NOVACK, M. R. Regularity for minimizers of a planar partitioning problem with cusps. *arXiv 2305.11865* (2023).
- [25] OWEN, N. C., RUBINSTEIN, J., AND STERNBERG, P. Minimizers and gradient flows for singularly perturbed bi-stable potentials with a Dirichlet condition. *Proc. Roy. Soc. London Ser. A* *429*, 1877 (1990), 505–532.
- [26] STERNBERG, P. The effect of a singular perturbation on nonconvex variational problems. *Arch. Rational Mech. Anal.* *101*, 3 (1988), 209–260.
- [27] STERNBERG, P. Vector-valued local minimizers of nonconvex variational problems. *Rocky Mountain J. Math.* *21*, 2 (1991), 799–807. Current directions in nonlinear partial differential equations (Provo, UT, 1987).
- [28] STERNBERG, P., AND ZEIMER, W. P. Local minimisers of a three-phase partition problem with triple junctions. *Proc. Roy. Soc. Edinburgh Sect. A* *124*, 6 (1994), 1059–1073.
- [29] ZUNIGA, A., AND STERNBERG, P. On the heteroclinic connection problem for multi-well gradient systems. *J. Differential Equations* *261*, 7 (2016), 3987–4007.