

# Follow-up LOFAR observations of the $\tau$ Boötis exoplanetary system

Jake D. Turner<sup>1,2</sup>, Jean-Mathias Grießmeier<sup>3,4</sup>, Philippe Zarka<sup>4,5</sup>, Xiang Zhang<sup>5</sup>, Emilie Mauduit<sup>5</sup>

<sup>1</sup> Department of Astronomy and Carl Sagan Institute, Cornell University, Ithaca, NY, USA  
e-mail: jaketurner@cornell.edu

<sup>2</sup> NHFP Sagan Fellow

<sup>3</sup> Laboratoire de Physique et Chimie de l'Environnement et de l'Espace (LPC2E) Université d'Orléans/CNRS, Orléans, France

<sup>4</sup> Observatoire Radioastronomique de Nançay (ORN), Observatoire de Paris, PSL Research University, CNRS, Univ. Orléans, OSUC, 18330 Nançay, France

<sup>5</sup> LESIA, Observatoire de Paris, CNRS, PSL, Meudon, France

## ABSTRACT

**Context.** Observing the radio emission from exoplanets is among the most promising methods to detect their magnetic fields and a measurement of an exoplanetary magnetic field will help constrain the planet's interior structure, star-planet interactions, atmospheric escape and dynamics, and habitability. Recently, circularly polarized bursty and slow emission from the  $\tau$  Boötis ( $\tau$  Boo) exoplanetary system was tentatively detected using LOFAR (LOW-Frequency ARray) beamformed observations. If confirmed, this detection will be a major contribution to exoplanet science. However, follow-up observations are required to confirm this detection.

**Aims.** Here, we present such follow-up observations of the  $\tau$  Boo system using LOFAR. These observations cover 70% of the orbital period of  $\tau$  Boo b including the orbital phases of the previous tentative detections.

**Methods.** We used the BOREALIS pipeline to mitigate radio frequency interference and to search for bursty and slowing varying radio signals. BOREALIS was previously used to find the tentative radio signals from  $\tau$  Boo.

**Results.** Our new observations do not show any signs of bursty or slow emission from the  $\tau$  Boötis exoplanetary system.

**Conclusions.** The cause for our non-detection is currently degenerate. It is possible that the tentative radio signals were an unknown instrumental systematic or that we are observing variability in the planetary radio emission due to changes in its host star. More radio data (preferably multi-site) and ancillary observations (e.g. magnetic maps) are required to further investigate the potential radio emission from the  $\tau$  Boötis exoplanetary system.

**Key words.** Planets and satellites: magnetic fields – Radio continuum: planetary systems – Magnetic fields – Astronomical instrumentation, methods and techniques – Planet-star interactions – Planets and satellites: aurorae – planets and satellites: gaseous planets

## 1. Introduction

A confirmed direct detection of exoplanetary radio emission has been elusive for decades. Observing the radio emission from exoplanets is among the most promising methods to detect their magnetic fields (Grießmeier 2015) and a measurement of an exoplanetary magnetic field will help constrain the planet's interior structure, star-planet interactions, and atmospheric escape (e.g., Zarka et al. 2015; Grießmeier 2015; Lazio et al. 2016; Grießmeier 2017; Lazio 2018; Griessmeier 2018; Zarka 2018; Lazio et al. 2019). A magnetic field measurement can help also help break degeneracies in the mass-radius diagram (Lazio et al. 2019). Additionally, atmospheric dynamics may be influenced by the presence of a planetary magnetic field (e.g., Perna et al. 2010a; Rauscher & Menou 2013; Rogers & Komacek 2014; Hindle et al. 2021) and Ohmic dissipation might be one of the factors causing the anomalously large radii of hot Jupiters (e.g. Perna et al. 2010a,b; Knierim et al. 2022). Finally, a magnetic field might be one of the many properties needed on Earth-like exoplanets to sustain their habitability (e.g., Grießmeier et al. 2005, 2015, 2016; Lammer et al. 2009; Kasting 2010; Owen & Adams 2014; Lazio et al. 2010a, 2016; Meadows & Barnes 2018; McIntyre et al. 2019; Gronoff et al. 2020; Green et al. 2021b).

Observations of auroral radio emissions produced via the Cyclotron Maser Instability (CMI) mechanism (Wu & Lee 1979; Zarka 1998; Treumann 2006) have been used to directly measure the magnetic field strengths of some of the magnetized planets in our Solar System (Burke & Franklin 1955; Zarka 1998). CMI radio signals are emitted at the local electron cyclotron frequency in the source region. Hence, the maximum gyrofrequency is directly proportional to the maximum magnetic field near the planetary surface. CMI radio emission is circularly polarized, beamed, and time-variable (e.g., Zarka 1998; Zarka et al. 2004).

Many theoretical and observational studies have focused on studying exoplanetary radio emissions over the past few decades. The reviews by Zarka et al. (2015) and Grießmeier (2015, 2017) cover the subject in great detail. A large body of theoretical work<sup>1</sup> has been published since the foundational

<sup>1</sup> e.g., de Pater 2000; Farrell et al. 2004; Lazio et al. 2004; Stevens 2005; Grießmeier et al. 2005; ?; Grießmeier et al. 2007b; Jardine & Collier Cameron 2008; Vidotto et al. 2010; Hess & Zarka 2011; Nichols 2011, 2012; Vidotto et al. 2012; Saur et al. 2013; See et al. 2015; Vidotto et al. 2015; Nichols & Milan 2016; Fujii et al. 2016; Vidotto & Donati 2017; Weber et al. 2017a,b, 2018; Lynch et al. 2018; Zarka 2018; Wang & Loeb 2019; Kavanagh et al. 2019; Vidotto et al. 2019;

theoretical studies of Zarka et al. (1997), Farrell et al. (1999), Zarka et al. (2001), and Zarka (2007). Some of these studies aim to develop theoretical models that can predict the frequency and intensity of radio emissions observed from Earth. A variety of scaling laws have been developed to estimate the magnetic field strengths of exoplanets (e.g. Farrell et al. 1999; Sánchez-Lavega 2004; Grießmeier et al. 2007b; Christensen et al. 2009; Christensen 2010; Reiners & Christensen 2010). Likewise, the Radio-Magnetic Bode’s law is a common approach for estimating the radio emission intensity of these planets (e.g., Farrell et al. 1999; Lazio et al. 2004; Zarka 2007; Grießmeier et al. 2007b; Grießmeier 2017; Zarka et al. 2018; Zarka 2018).

Over the past few decades, there have been many exoplanet radio observing campaigns, but none of these observations have resulted in an unambiguous detection<sup>2</sup>. There have been a few tentative detections (Lecavelier des Etangs et al. 2013; Sirothia et al. 2014; Vasylieva 2015; Bastian et al. 2018), but none of them have been confirmed by follow-up observations. The many degenerate reasons for the non-detections are discussed in great detail in the review articles Zarka et al. (2015) and Grießmeier (2015, 2017). Some reasons for the non-detections could be that the observations were not sensitive enough (e.g., Bastian et al. 2000), there was no emission at the observed frequencies due to a low magnetic field strength (e.g. Hallinan et al. 2013), Earth was outside the CMI beaming pattern (Hess & Zarka 2011; Ashtari et al. 2022), and/or CMI mechanism was not operating (e.g., plasma frequency of the planet’s ionosphere was greater than the cyclotron frequency; Grießmeier et al. 2007b; Weber et al. 2017a,b; Lamy et al. 2018; Weber et al. 2018). By contrast, advancement has been made in detecting free-floating planets near the brown dwarf boundary (Kao et al. 2016, 2018) and stellar emission potentially created by star-planet interactions (e.g., Vedantham et al. 2020; Callingham et al. 2021; Pérez-Torres et al. 2021; Pineda & Villadsen 2023; Triglio et al. 2023; Blanco-Pozo et al. 2023).

In a recent study, circularly polarized bursty and slow emission from the  $\tau$  Boo ( $\tau$  Boo) exoplanetary system was tentatively detected by Turner et al. (2021, hereafter T21) using LOFAR (Low Frequency ARray; van Haarlem et al. 2013) beamformed observations taken in 2017. The slow and burst emission were seen at orbital phases of 0.65 and 0.8, respectively, in relation to the  $\tau$  Boo b’s periastron (see Figure 1). The bursty emission was seen from 15–21 MHz, while the slower emission was present in the 21–30 MHz range. T21 did not identify any plausible sources of false positives for the bursty emission, but when it came to the slower emission, uncertainty arose due to the possibility of unknown systematic effects causing the signal. Assuming an astronomical origin for the signal, T21 concluded that the signals originated from the  $\tau$  Boo system, with the most likely cause being CMI radio emission from the exoplanet  $\tau$  Boo b. Under the premise of a planetary origin for the radio signals, T21 derived planetary magnetic field constraints

that are consistent with theoretical predictions (Grießmeier et al. 2007b; Grießmeier 2017). More recent calculations (Mauduit et al. 2023; Ashtari et al. 2022) using an updated version of these predictions also yielded magnetic moment and flux density values that aligned with the earlier theoretical estimates, in addition to matching the observed phases and the handedness of the circular polarization (Ashtari et al. 2022). Follow-up observations were highly advocated by T21 to confirm their tentative detections and to search for periodicity in the signal.

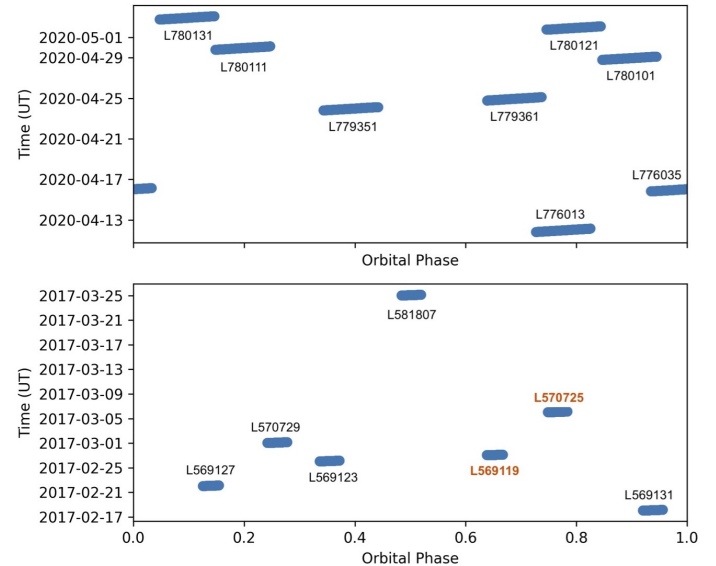


Fig. 1: Orbital phase coverage for the LOFAR observations of  $\tau$  Boo b. The observations from this study and Turner et al. (2021) are shown in the top and bottom panels, respectively. The orbital phase was calculated relative to periastron of  $\tau$  Boo b using the ephemeris (Period = 3.312463 days,  $T(0)_{\text{per}}=2446957.81$  JD) from Wang & Ford (2011). The  $\tau$  Boo LOFAR observations with the tentative detections (L569119 and L570725) from Turner et al. (2021) are displayed in red.

Motivated by the tentative detection of radio emission from  $\tau$  Boo, we performed a large follow-up campaign in 2020. This campaign was coordinated between four different radio telescopes including LOFAR and NenuFAR (New Extension in Nançay Upgrading LOFAR; Zarka et al. 2020). In this paper, we only present the LOFAR data of  $\tau$  Boo since the sensitivity of LOFAR has been well characterized using simulated Jupiter radio emission data (Turner et al. 2019, hereafter T19) and a complete data reduction pipeline is available (T19; T21). A preliminary analysis of the NenuFAR data is presented elsewhere (Turner et al. 2023).

In this paper, Sections 2, 3, and 4 describe the  $\tau$  Boo observations, data processing, and data analysis, respectively. Section 5 discusses the possible reasons for our non-detections and future steps.

## 2. Observations

All of the observations presented here were taken with the Low Band Antenna (LBA) of LOFAR (van Haarlem et al. 2013) using beamformed mode (Stappers et al. 2011). We obtained 8 exoplanet observations (each 8 hours long) for a total of 64 hours, plus 8 observations of the pulsar B0809+74 totaling 82 minutes. Similar to our previous studies (Turner et al. 2017a, T19, T21,

Shiohira et al. 2020; de Gasperin et al. 2020; Green et al. 2021b; Sciola et al. 2021; Narang et al. 2021; Hori 2021; Cendes et al. 2022

<sup>2</sup> e.g., Yantis et al. 1977; Winglee et al. 1986; Zarka et al. 1997; Bastian et al. 2000; Farrell et al. 2003; Ryabov et al. 2004; Shiratori et al. 2006; George & Stevens 2007; Lazio & Farrell 2007; Smith et al. 2009; Lecavelier Des Etangs et al. 2009, 2011; Lazio et al. 2010a,b; Stroe et al. 2012; Hallinan et al. 2013; Murphy et al. 2015; Lynch et al. 2017; Turner et al. 2017a; O’Gorman et al. 2018; Lenc et al. 2018; Lynch et al. 2018; Route 2019; Narang et al. 2020; de Gasperin et al. 2020; Green & Madhusudhan 2021a; Turner et al. 2021; Narang et al. 2021; Cendes et al. 2022; Narang 2022; Narang et al. 2023a,b; Route & Wolszczan 2023; Turner et al. 2023; Shiohira et al. 2024; Narang et al. 2024

Turner et al. 2023), the pulsar observations are used as a calibrator to study systematics in the data and to ensure the reliability of the processing. An example OFF-beam dynamic spectrum is shown in Figure 2. The dates and times of the observations can be found in Table 2; it is similar to that in T21. Most importantly, we ensure we cover the 15–30 MHz range where T21 detected their tentative signals. We only present the Stokes-V<sup>3</sup> in this paper since the tentative detections were seen in Stokes-V (T21).

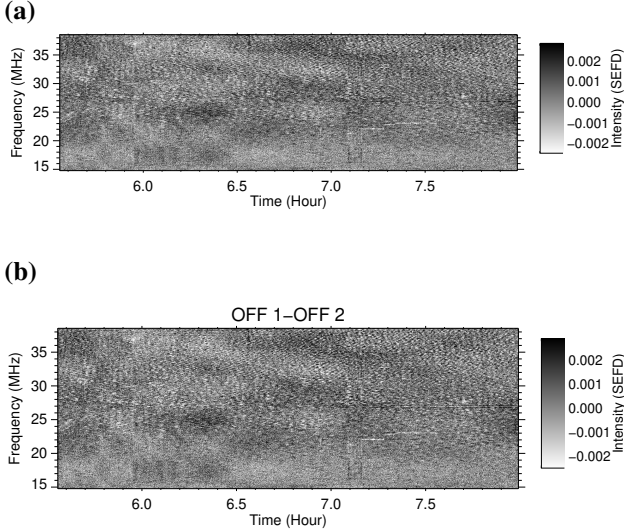


Fig. 2: Example dynamic spectrum of the OFF-beam (panel a) and the difference (panel b) between two OFF-beams in Stokes-V for the L780111 (2020-04-29)  $\tau$  Boo observation. This dynamic spectrum has been processed with BOREALIS.

In this study, we compare our on-target beam (“ON-beam”) to several simultaneous beams pointing to nearby locations in the sky (“OFF-beam 1”, “OFF-beam 2” and “OFF-beam 3”). This is similar to our previous studies (Turner et al. 2017a, 2019, 2021) except we now use three OFF-beams instead of two to further investigate systematic effects in the analysis. This method relies on the assumption that the OFF beams can effectively characterize and account for ionospheric fluctuations, radio frequency interference, and any other systematic errors that may be present. The beamformed LOFAR observations of 55 Cancri,  $\nu$  Andromedae, and  $\tau$  Boo demonstrated that this assumption holds true (T21).

Our new LOFAR observations cover 70% of the orbit of  $\tau$  Boo b as seen in Figure 1. Due to the emission geometry (Hess & Zarka 2011; Ashtari et al. 2022), it is anticipated that the emission will only be directed towards Earth for a limited period. Therefore, the observation windows were selected to ensure the widest possible orbital coverage to maximize the chances to detect beamed emission. Most importantly, we also obtained two new observations around phase 0.65 (2020-04-11 and 2020-05-01) and one new observation around phase 0.8 (2020-04-15) that cover the same phases as the tentative slow emission and burst emission signal from T21, respectively. These phases are most likely to have detectable emission as shown in the modeling of the time-dependent CMI beaming effect by Ashtari et al. (2022).

<sup>3</sup> Similar to T21, we only analyze  $V'$  as defined in equations 9–11 of Section 4.2.1 in T19.

Table 1: Summary of the observations of  $\tau$  Bootis in LOFAR Cycle 13. Column 1: observation number. Column 2: LOFAR observation ID. Column 3: date and start time of the observation (UTC). Column 4: duration ( $\Delta t$ ) of the observation. Column 5:  $\tau$  Boo b orbital phases of the observation relative to periastron and using the ephemeris (Period = 3.312463 days,  $T(0)_{\text{per}}=2446957.81$  JD) from Wang & Ford 2011.

Obs #	LOFAR ID	Date & Start Time (UTC)	$\Delta t$ (hrs)	Phase
1	L776013	2020-04-11 20:00	8	0.73–0.83
2	L776035	2020-04-15 20:00	8	0.93–0.03
3	L779351	2020-04-23 19:30	8	0.34–0.44
4	L779361	2020-04-24 19:00	8	0.64–0.74
5	L780101	2020-04-28 19:00	8	0.85–0.95
6	L780111 <sup>a</sup>	2020-04-29 19:00	8	0.15–0.25
7	L780121	2020-05-01 18:30	8	0.75–0.85
8	L780131	2020-05-02 18:30	8	0.05–0.15
B0809+74 [82 min.]				
1	L776009	2020-04-11 1949	0.17	–
2	L776031	2020-04-15 1949	0.17	–
3	L779347	2020-04-23 1919	0.17	–
4	L779357	2020-04-24 1849	0.17	–
5	L780097	2020-04-28 1849	0.17	–
6	L780107	2020-04-29 1849	0.17	–
7	L780117	2020-05-01 1819	0.17	–
8	L780127	2020-05-02 1819	0.17	–

Notes. <sup>(a)</sup> Only the last 3 hours of this observation are useful due to strong RFI at the start of the observations.

Table 2: Setup of the LOFAR-LBA beamformed observations of  $\tau$  Bootis

Parameter	Value	Units
Array Setup	LOFAR Core	
Number of Stations	24	
Beams	ON & 3 OFF	
Configuration	LBA outer antennas	
Antennas per Station	48	
Minimum Frequency	15	MHz
Maximum Frequency	39	MHz
Subbands Recorded	244	
Subband Width	195	kHz
Channels per Subband	64	
Frequency Resolution (b)	3.05	kHz
Time Resolution ( $\tau$ )	10.5	msec
Beam Diameter	13.8	arcmin
Stokes Parameters	IQUV	

Notes. <sup>(a)</sup> Calculated at 30 MHz (van Haarlem et al. 2013).

### 3. Data pipeline

All observations were processed with the BOREALIS (BeamfOrmed Radio Emission AnaLysis) pipeline (Vasylieva 2015; Turner et al. 2017a, 2019, 2021). BOREALIS is split into two main components: processing and post-processing. The processing part of BOREALIS performs RFI mitigation, corrects for any large-scale time-frequency (t-f) response variations of the telescope (i.e. gain variations), and rebins the RFI mitigated and corrected data. All parts of the analysis were done exactly as in

T21. The RFI mitigation combines four different techniques (Of-  
fringa et al. 2010, 2012; Offringa 2012; Zakharenko et al. 2013;  
Vasylieva 2015) for optimal efficiency and processing time. The  
t-f response of the telescope is empirically determined assuming  
a quadratic dependence in time for each frequency. BOREALIS  
has been extensively tested on pulsar (Turner et al. 2017a) and  
scaled Jupiter radio emission data (Turner et al. 2019) to ensure  
optimal RFI mitigation and t-f response variation corrections  
without sacrificing detection capabilities. See Figure 2 for exam-  
ple OFF-beam dynamic spectrum processed through BOREALIS.

We performed the post-processing of the LOFAR data ex-  
actly as in T21. A very detailed description of the burst and slow  
emission post-processing observables can be found in T19 and  
T21. A brief summary is given in the following paragraph. The  
post-processing is performed on the absolute value of the cor-  
rected Stokes-V data (as defined in T19). For burst emission, we  
only use the Q2 and Q4a-f observable quantities. The Q2 observ-  
able consists of one time series at a time resolution of 1 second  
per beam (ON, OFF1, OFF2, and OFF3). It is obtained by high-  
pass filtering the dynamic spectrum. To search for faint emission,  
the high-pass filtered dynamic spectrum is integrated over a large  
frequency range (e.g., 10 MHz) to produce a time series. Q2 can  
be represented by a “scatter plot” comparing a pair of beams  
(e.g., the ON and one of the OFF beams) and is designed to find  
bursty emission (with time scales  $< \sim 1$  minute). The quantities  
Q4a to Q4f provide statistical measures of the bursts identified  
by the Q2 quantity. When examining Q4a-f, the ON and one of  
the OFF time series are compared to each other; for this, we  
introduce the difference curve  $Q4f_{\text{Diff}} = Q4f(\text{ON}) - Q4f(\text{OFF})$ . We  
then plot this curve against a reference curve computed from  
10000 draws of purely Gaussian noise. The post-processing was  
performed separately over 3 different frequency ranges (15-26  
MHz, 15-39 MHz, and 26-29 MHz).

## 4. Data analysis and results

We searched for excess signal in the ON-beam both by eye and  
using the automated search procedure outlined in T21. To search  
for a possible detection of burst emission we applied the criteria  
(N1 to N12) defined in T19 and T21. We do not find any slow or  
burst emissions in the observations. The time series (Q1a) and  
integrated spectra (Q1b) for all beams can be found in Figure  
3. The top and bottom panels of Figure 3 show the time series  
and integrated spectra, respectively. Unlike in T21, we do  
not see a sinusoidal pattern (that may have been caused by im-  
perfect phasing of the array) in the integrated spectrum (Figure  
3c-d). Next, we examine whether there are any systematic dif-  
ferences in the noise between beams. To do this, we tested for  
the presence of red noise in the difference time series and in-  
tegrated spectra (bottom panels Figure 3a-d) for all the obser-  
vations using the time-averaging method (Pont et al. 2006) and  
wavelet technique<sup>4</sup> (Carter & Winn 2009) as implemented in the  
EXOMOP code by Turner et al. (2016, 2017b), and found none.  
Thus any differences in the signals between the ON and OFF  
beams can be explained purely by Gaussian white noise. The  
burst statistics were also similar between all ON and OFF beams.  
An example of the diagnostics for a non-detection can be found  
in Figure 4, based on the L776035 (2020-04-15) observation,  
analyzed within 15-27 MHz. The ON-OFF and OFF-OFF differ-  
ence curves are very similar to each other. In summary, we did

<sup>4</sup> The wavelet technique assumes that the time series is an additive  
combination of noise with Gaussian white noise and red noise (charac-  
terized as a power spectral density proportional to  $1/f^\alpha$ ).

not detect any slow or bursty signals in our data. In Section 5,  
we will discuss the implications of these non-detections and our  
derived upper-limits.

## 5. Discussion

We do not detect any bursty or slow emission in the new LOFAR  
observations. If the signals in T21 were coming from the planet  
we could have expected to detect them again at the same phases.  
However, this assumes no variability in the radio signal and  
that the signals are from the planet. We discuss these possibil-  
ities more below. Our non-detection is consistent with the non-  
detection of bursty emission from five NenuFAR observations  
(Turner et al. 2023) taken simultaneously with the LOFAR ob-  
servation (see Figure 1 in Turner et al. 2023). Most importantly,  
both NenuFAR and LOFAR don’t observe bursty emission on  
April 15, 2020, which covered the same phase as the T21 bursty  
tentative detection.

### 5.1. Degenerate causes of the non-detection

There are many different degenerate reasons which could explain  
our non-detection of emission from  $\tau$  Boo. We explore the pos-  
sibilities in greater detail below:

1. The first possibility is that the original slow and/or bursty  
signals seen by T21 were caused by statistical anomalies  
or unknown unidentified instrumental systematics. In T21,  
it was suggested that the OFF beam on the slow emission de-  
tection behaved abnormally (the flux on that date was much  
lower than all the other dates) compared to the rest of the ob-  
servations. This possibility can be investigated further with  
the new observations. We find that the OFF beams in the  
L570725 are not abnormal when compared to the OFF beams  
in the new 2020  $\tau$  Boo observations (see Figure 5). This sug-  
gests that the original slow emission signal may have not  
been an unknown instrumental systematic but it doesn’t com-  
pletely rule it out.
2. The planetary emission from  $\tau$  Boo b might be variable. The  
expected intensity of the planetary radio emission can differ  
greatly with stellar rotation (e.g. Grießmeier et al. 2005;  
Grießmeier et al. 2007a; Fares et al. 2010; Vidotto et al.  
2012; See et al. 2015; Strugarek et al. 2022) and phase of the  
stellar magnetic cycle (e.g. Fares et al. 2009; Vidotto et al.  
2012; See et al. 2015; Elekes & Saur 2023). The latter point  
is very much a concern since the host star  $\tau$  Boo is known  
to undergo a rapid magnetic cycle of 120 days (Mittag et al.  
2017; Jeffers et al. 2018). See et al. (2015) showed that the  
expected radio flux of  $\tau$  Boo b can disappear entirely for cer-  
tain stellar rotation phases and also that it can vary (at opti-  
mal stellar longitude) by a factor of  $\sim 5$  along the stellar mag-  
netic cycle. Recently, Elekes & Saur (2023) found using nu-  
merical simulations of the  $\tau$  Boo system that if the magnetic  
field of the star were to reverse and become anti-aligned with  
the magnetic field of the planet, the planetary radio emis-  
sions would be significantly reduced and could fall below  
the detection threshold of LOFAR. Unfortunately, no stellar  
magnetic field measurements exist during the new LOFAR  
observations. Therefore, we cannot rule out planetary radio  
emission variability as the cause of our non-detection.
3. It is possible that one of the conditions to allow the CMI  
mechanism to operate is temporarily violated. For example,  
an extended and evaporating atmosphere may increase the  
local electron plasma frequency such that it is greater than

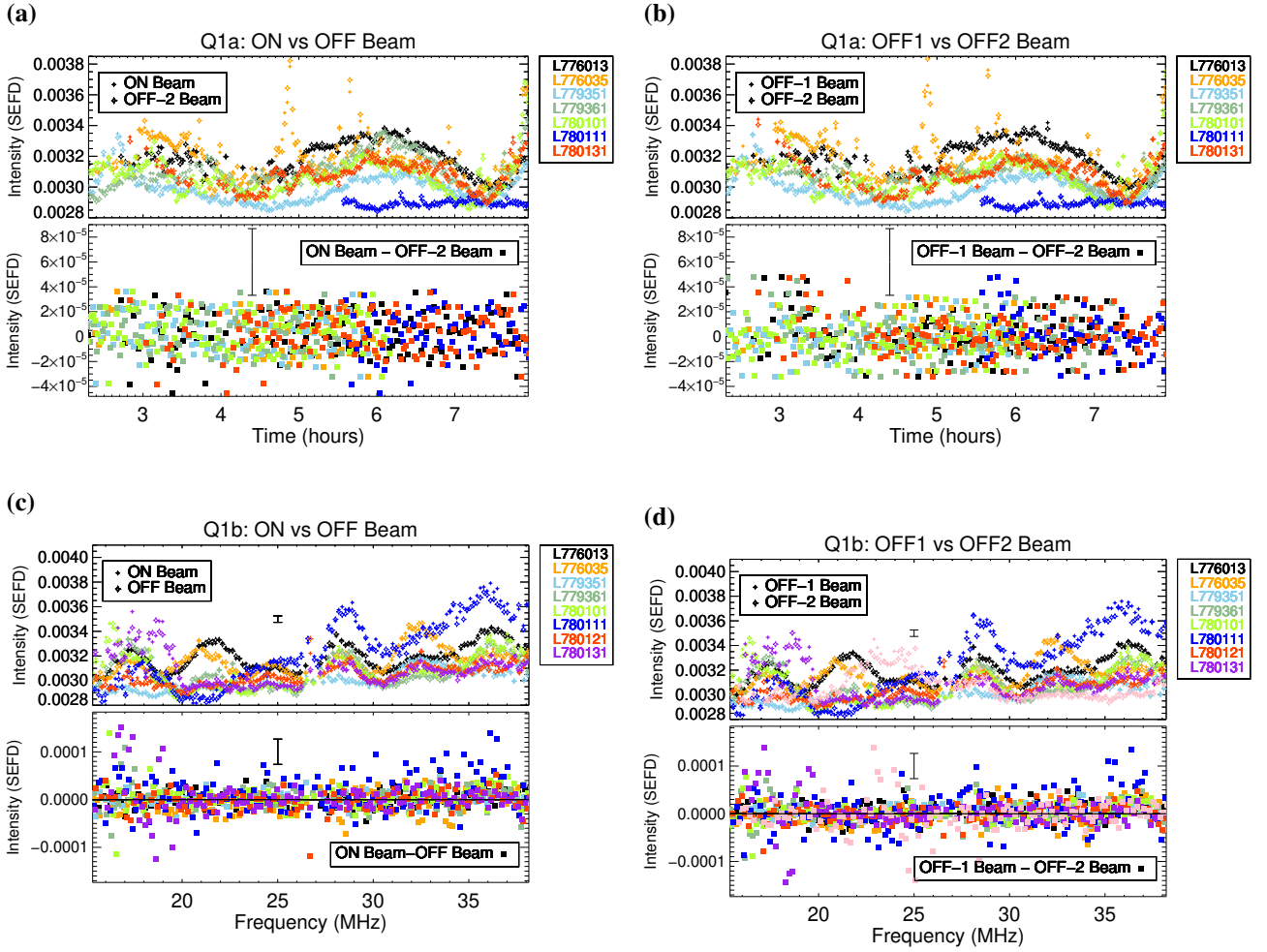


Fig. 3: Time-series (Q1a: panel a and b) and Integrated spectrum (Q1b: panel c and d) for different beams (ON, OFF-1, OFF-2) for the follow-up LOFAR  $\tau$  Bootis observations. The two-sigma error bars are shown in the black brackets. Large-scale features can be seen for all dates, however, they change between observations. No distinct emission in the ON-beam is seen when compared to the OFF-beams. There are no large-scale differences seen between the three beams for any date. The third OFF-beam is indistinguishable from the other OFF-beams within the noise.

the local electron cyclotron frequency,  $\omega_p > \omega_c$ ; (e.g. Weber et al. 2017a,b; Elekes & Saur 2023; Grießmeier et al. 2023). However, for  $\tau$  Boo b this effect is predicted to not be a problem due to its large mass (Weber et al. 2018). On the other hand, temporal variations could lead to circumstances where the emission does not operate continuously if the planetary magnetic field is as low as estimated by Reiners & Christensen (2010). Evidence suggests that the observed field strengths for brown dwarfs (Kao et al. 2016, 2018) cannot be explained by the Reiners & Christensen (2010) energy flux and magnetic energy density scaling relation. However, it is currently unclear if this contradiction extends to  $\tau$  Boo b and other hot Jupiters. Therefore, more work is needed to understand whether the plasma conditions at  $\tau$  Boo b vary such that the conditions for the CMI are not always fulfilled.

4. In contrast to Jupiter’s decametric emissions (Zarka 1998),  $\tau$  Boo b’s radio emission may not exhibit continuous activity. This would suggest that the magnetospheric dynamics of  $\tau$  Boo b could be different from those of Jupiter. For instance, the density of the magnetosphere might vary and at times be temporarily depleted in energetic electrons.

To determine the true cause of the potential radio variability of the  $\tau$  Boo system, it is crucial to conduct an extensive follow-up campaign. To rule out instrumental systematics, multi-site observations are recommended. It is necessary to monitor the planet’s behavior throughout its orbit and the host star’s magnetic cycle to distinguish between conflicting factors. Additionally, measurements of the star’s magnetic field must be taken during the radio follow-up campaign. Currently, there is an ongoing collaborative campaign utilizing NenuFAR to study  $\tau$  Boo radio emission and several telescopes (TBL/Neo-Narval and CHFT/ESPaDONs) to monitor the magnetic field of the host star. The results of this study will be presented in future work.

## 5.2. Upper-limits on the radio flux density

With our non-detection, we can place an upper limit on the radio emission from the  $\tau$  Boo system at the time of the observations. We find a  $3\sigma$  upper limit of 165 mJy from the range 15–39 MHz using the Q1a observable for the slowly varying emission. The noise characteristics of Q1a for the new observations is lower by a factor of 2 than the noise seen in the T21 detection. Using the attenuated Jupiter modeling done in Turner et al. (2019), we



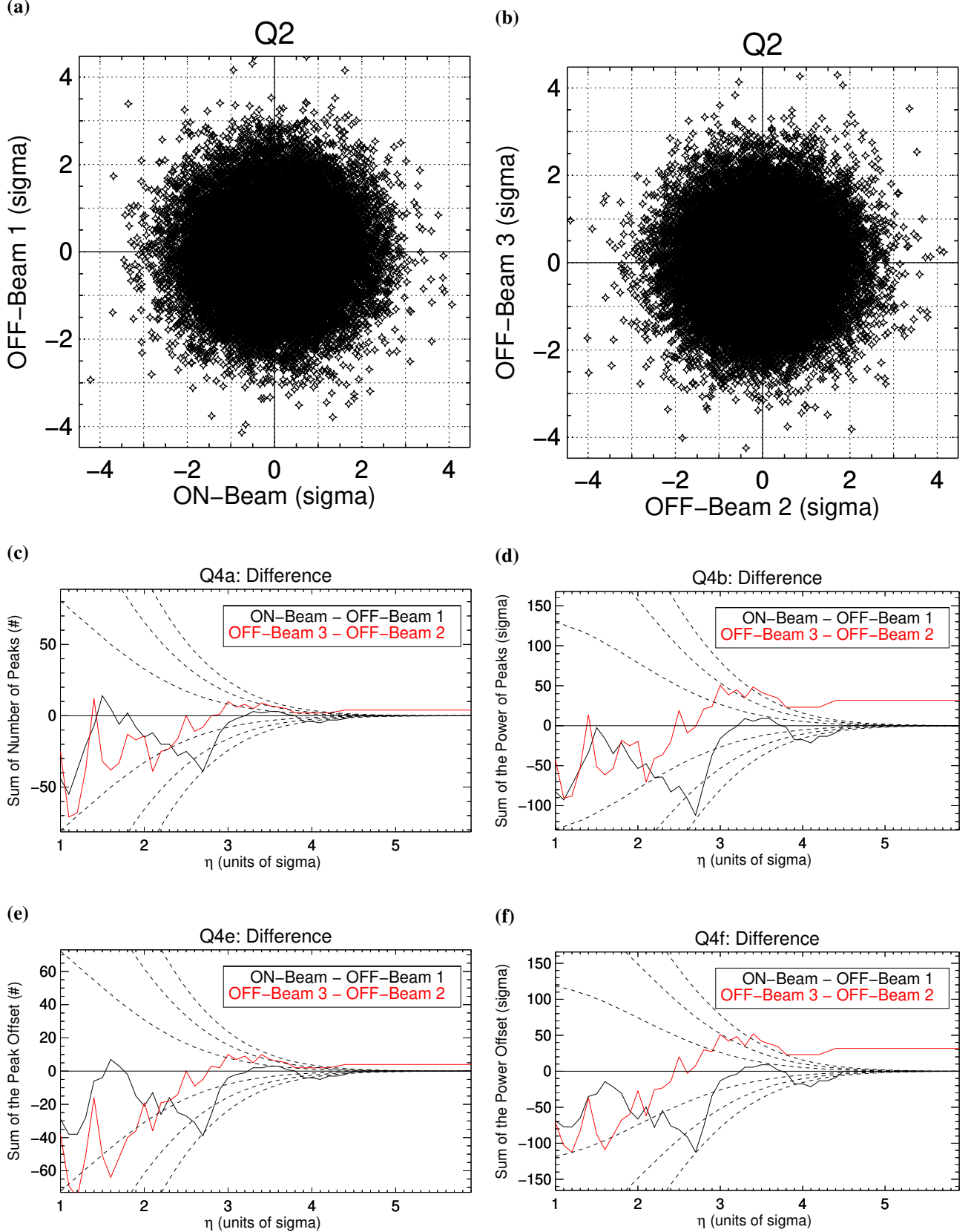


Fig. 4: Non-detection of burst emission for  $\tau$  Boo in the L776035 (2020-04-15) observation between 15-27 MHz in Stokes-V ( $|V'|$ ). Panel a: Q2 for the ON-beam vs the OFF-beam 2. Panel b: Q2 for the OFF-beam 1 vs the OFF-beam 2. Panel c: Q4a (number of peaks). Panel d: Q4b (power of peaks). Panel e: Q4e (peak offset). Panel f: Q4f (peak offset). For panels c to f the black lines are the ON-beam difference with the OFF beam 1 and the red lines are the OFF-beam difference. The dashed lines are statistical limits (1, 2, 3 $\sigma$ ) of the difference between all the Q4 values derived using two different Gaussian distributions (each performed 10000 times). We do not see any excess signal in the ON-beam compared to the OFF-beams. In fact, the OFF-beams show a false-positive weak signal at  $\eta \sim 4$ .

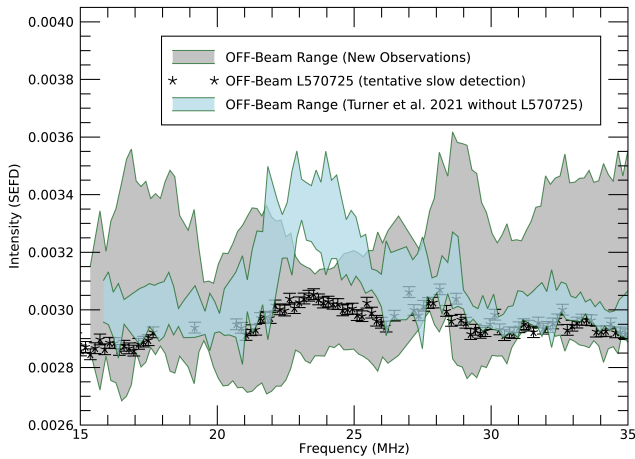


Fig. 5: Comparison of the integrated spectra (Q1b) of the LOFAR  $\tau$  Boötis OFF-beams between the new observations in this paper (gray-shaded area) and those presented (light-blue shaded area) in Turner et al. (2021). We show separately the OFF-beam in the L570725 observation (square black points), where a tentative detection of slow emission was found from Turner et al. (2021). The integrated spectrum for each date in the current observing campaign can be found in Figure 3. While the OFF-beam in L570725 is abnormal when compared to the rest of the Turner et al. (2021) observations, it is consistent with the new observations.

find for the burst emission an upper limit on the flux density that should be less than  $10^5 \times$  the peak flux of Jupiter’s decametric burst emission ( $\sim 5 \times 10^6$  Jy; Zarka et al. 2004).

## 6. Conclusions

In this study, we performed follow-up low-frequency beam-formed radio observations of the  $\tau$  Boötis ( $\tau$  Boo) exoplanetary system with LOFAR. Previous LOFAR observations showed tentative hints of a detection possibly originating from the exoplanetary system (Turner et al. 2021). In the new observations, we do not detect any burst or slow emission from  $\tau$  Boo (Figures 3–4). The current cause of the non-detection is degenerate (see Section 5.1) but could be caused by variability in the planetary radio emission. More radio observations (preferably multi-site) are needed that cover the full orbit of the planet multiple times and at different epochs of the stellar magnetic cycle. Near-simultaneous stellar magnetic maps and stellar monitoring are highly encouraged and needed to disentangle the various competing effects. The newly commissioned NenuFAR and LWA-OVLO are scheduled to observe the  $\tau$  Boo exoplanetary system and will help investigate the cause of the possible variable emission.

## Acknowledgements

J.D. Turner was supported for this work by NASA through the NASA Hubble Fellowship grant #HST-HF2-51495.001-A awarded by the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Incorporated, under NASA contract NAS5-26555.

P. Zarka acknowledges funding from the ERC N° 101020459—Exoradio.

This work was supported by the “Programme National de Planétologie” (PNP) of CNRS/INSU co-funded by CNES and by the “Programme National de Physique Stellaire” (PNPS) of CNRS/INSU co-funded by CEA and CNES.

This research has made use of the Extrasolar Planet Encyclopaedia (exoplanet.eu) maintained by J. Schneider (Schneider et al. 2011), the NASA Exoplanet Archive, which is operated by the California Institute of Technology, under contract with the National Aeronautics and Space Administration under the Exoplanet Exploration Program, and NASA’s Astrophysics Data System Bibliographic Services. This research has also made use of Aladin sky atlas developed at CDS, Strasbourg Observatory, France (Bonnarel et al. 2000; Boch & Fernique 2014).

This paper is based on data obtained with the International LOFAR Telescope (ILT) under project codes LC7\_013 and LC13\_027. LOFAR (van Haarlem et al. 2013) is the Low Frequency Array designed and constructed by ASTRON. It has observing, data processing, and data storage facilities in several countries, that are owned by various parties (each with their own funding sources), and that are collectively operated by the ILT foundation under a joint scientific policy. The ILT resources have benefited from the following recent major funding sources: CNRS-INSU, Observatoire de Paris and Université d’Orléans, France; BMBF, MIWF-NRW, MPG, Germany; Science Foundation Ireland (SFI), Department of Business, Enterprise and Innovation (DBEI), Ireland; NWO, The Netherlands; The Science and Technology Facilities Council, UK.

We thank Mickael Coriat, Cyril Tasse, and Vyacheslav Zakharenko for help with the LC13\_027 LOFAR proposal.

We acknowledge the use of the Nançay Data Center computing facility (CDN - Centre de Données de Nançay). The CDN is hosted by the Nançay Radio Observatory in partnership with Observatoire de Paris, Université d’Orléans, OSUC and the CNRS. The CDN is supported by the Région Centre-Val de Loire, département du Cher.

We thank the ASTRON staff for their help with these observations.

During the process of writing this paper, Jake D. Turner’s young cousin, Evan Turner, passed away unexpectedly. This paper is dedicated to you Evan, keep looking up.

**Facilities:** LOFAR (van Haarlem et al. 2013)

**Software:** BOREALIS (Vasylieva 2015; Turner, Grießmeier, Zarka, & Vasylieva 2017a; Turner, Grießmeier, Zarka, & Vasylieva 2019); IDL Astronomy Users Library (Landsman 1995); Coyote IDL created by David Fanning and now maintained by Paulo Penteado (JPL).

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