

Some understanding of recent inelastic neutron scattering on Ce-pyrochlores: Signatures of electric monopoles

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We revisit the pyrochlore spin liquid materials $\text{Ce}_2\text{Sn}_2\text{O}_7$ and $\text{Ce}_2\text{Zr}_2\text{O}_7$ by examining the existing experiments. We continue to rely on the special properties of the dipole-octupole nature of the Ce^{3+} moment. The inelastic neutron scattering (INS) measurement in the octupolar U(1) spin liquid *selects* the (gapped) spinon continuum, and thus has suppressed spectral weights below the energy threshold of two spinon gaps. This measurement, however, includes all other emergent excitations at lower energies in the dipolar U(1) spin liquid, in particular, the gapless photon and the continuum of the *electric monopoles*. Although the electric monopole continuum is weakly gapped (compared to the larger spinon gap), the energy scale is actually close to the gauge photons, and the spectrum largely overlaps with the photons. Due to the background dual π flux for the electric monopoles, the density of states is enhanced at lower energies, creating peak structures. This can be contrasted with the linearly suppressed spectral weight of the gauge photons at low energies. We propose that the electric monopole continuum should be mostly responsible for the low-energy spectrum in the INS measurement in the dipolar U(1) spin liquid. With these understanding and calculation, we discuss the available experimental results and predict further experiments for $\text{Ce}_2\text{Sn}_2\text{O}_7$ and $\text{Ce}_2\text{Zr}_2\text{O}_7$.

There has been a continuous interest in the pyrochlore quantum spin liquid [1, 2]. It is an exotic quantum state that is understood quite well theoretically and at the same time has a strong experimental connection with the pyrochlore quantum spin ice materials [3–5]. After the intensive exploration of the Tb-pyrochlores, Yb-pyrochlores and Pr-pyrochlore [1, 2], the recent activities were about the Ce-based pyrochlore magnets, where the Ce^{3+} local moments were identified as the dipole-octupole (DO) doublets and related to distinct symmetry-enriched spin liquids [6, 7]. There are now three different Ce-based pyrochlore compounds, $\text{Ce}_2\text{Sn}_2\text{O}_7$, $\text{Ce}_2\text{Zr}_2\text{O}_7$ and $\text{Ce}_2\text{Hf}_2\text{O}_7$ [8–17]. Most experiments were performed on the first two compounds so far. In this work, we review some of the existing experiments, mostly the inelastic neutron scattering (INS) results, and aim to provide some insights from the theoretical perspective to the understanding of the experiments.

Since our major point is about an important corner of the low-energy excitation spectra that distinguishes the dipolar U(1) spin liquid and the octupolar U(1) spin liquid, we simply state the results here. In short, the INS measurement could detect all the emergent excitations for the dipolar U(1) spin liquid. In particular, due to the intrinsic π dual U(1) gauge flux, the electric monopole continuum stands out at low energies and creates the peak structures on top of the suppressed photon intensity. In reality, the whole INS spectrum contains the peaks from the electric monopole continuum at low energies and the peaks from the spinon continuum at higher energies, in addition to the linearly suppressed photon modes [18, 19] at low energies. The presence of the electric monopole continuum for the dipolar U(1) spin liquid

for the DO doublets in the INS measurement were actually claimed in our previous work [20, 21]. Because many results were scattered in different parts in the previous works, the relation with the later experimental results has not yet been fully synthesized and clearly made. To better serve the community, we fulfill this task by reviewing the existing experiments and performing expanded calculations based on the existing experimental progress. Arriving these results requires the combination of the microscopic nature of the DO doublet, the effective spin model, and the experimental measurement. To make the presentation of this work self-contained, we will start from the basics before embarking on a bit new outcomes.

We begin with the dipole-octupole (DO) doublet on the pyrochlore lattice. Each state of this doublet is a one-dimensional irreducible representation of the D_{3d} point group, and is transformed into the other by the time reversal operation [6]. Thus, they are special type of Kramers' doublet whose degeneracy is protected by the time reversal symmetry. More specially, their wavefunctions are linear superpositions of the local $J^z = 3m/2$ (with m an odd integer) where the z direction is along the local 111 direction at each magnetic ion. As it is demanded by the point group symmetry, this doublet generally occurs among the crystal field states of all spin-orbit-coupled local- J moments with $J > 1$. If the local moment J happens to be an odd multiplier of $3/2$ such as $J = 9/2$ for Nd^{3+} in $\text{Nd}_2\text{Zr}_2\text{O}_7$ [24, 25] and $\text{Nd}_2\text{Sn}_2\text{O}_7$ [26, 27] and $J = 15/2$ for Dy^{3+} in $\text{Dy}_2\text{Ti}_2\text{O}_7$ [6], an easy-axis anisotropy would favor the DO doublet as the crystal field ground state. For others J 's, this way to obtain the DO doublet as the crystal field ground state does not seem to apply.

We are particularly interested in the Ce-pyrochlores with $J = 5/2$. The possible spin liquid was first experimentally proposed for $\text{Ce}_2\text{Sn}_2\text{O}_7$ [8]. The DO doublet na-

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Different U(1) QSLs	Inelastic neutron scattering measurement
Octupolar U(1) QSL for DO doublets	Gapped spinon continuum
Dipolar U(1) QSL for DO doublets	Gapless gauge photon, <i>gapped electric monopole continuum*</i> and gapped spinon continuum
Dipolar U(1) QSL for non-Kramers doublets	Gapless gauge photon, gapped electric monopole continuum
Dipolar U(1) QSL for usual Kramers doublets	Gapless gauge photon, gapped electric monopole continuum and gapped spinon continuum

TABLE I. List of different U(1) QSLs for their inelastic neutron scattering properties. Here “usual Kramers doublet” refers to the Kramers doublet that differs from a DO doublet. This is an improved version of Tab.I in Ref. 7 by incorporating the understanding of Refs. 21 and 20. The electric monopole was referred as “magnetic monopole” in Refs. 3, 20, and 22. The point at * is what we emphasize and distinguish from the interpretation in Ref. 23.

ture of ground state doublet for the Ce^{3+} ion was clarified in the theoretical work of Ref. 7, as well as the connection to the pyrochlore ice U(1) spin liquids. The key insight of the *selective* spectroscopic measurements for the spinon continuum in the octupolar U(1) spin liquid by the INS spectroscopy was also made over there [7]. Later developments were achieved for both $\text{Ce}_2\text{Zr}_2\text{O}_7$ and $\text{Ce}_2\text{Sn}_2\text{O}_7$, especially $\text{Ce}_2\text{Zr}_2\text{O}_7$ due to the single-crystal samples [8–17, 28]. The relevance of the π -flux U(1) spin liquid (or U(1) $_{\pi}$ spin liquid) was then raised, and the translational symmetry enrichment on top of the point-group symmetry enrichment, and the symmetry fractionalization were discussed [21].

Later theoretical works have also made progress to examine the distinction between different symmetry enriched U(1) spin liquids [29–33]. The recent thermodynamics and model calculation were even able to extract the parameters of the exchange couplings [16, 34]. The results are different for $\text{Ce}_2\text{Zr}_2\text{O}_7$ and $\text{Ce}_2\text{Sn}_2\text{O}_7$ that are likely located in different U(1) spin liquid ground states of the underlying XYZ spin model. More recent INS measurements indeed found different excitation spectra for $\text{Ce}_2\text{Zr}_2\text{O}_7$ and $\text{Ce}_2\text{Sn}_2\text{O}_7$ [23, 28]. Since very low-temperature thermodynamics may contain some uncertainty in the measurements and fitting, we rely mostly on the neutron scattering measurements for the major reasoning. To understand the INS experiments, we start from the spin model for the Ce-pyrochlores [6, 7],

$$H = \sum_{\langle ij \rangle} [J_x \tau_i^x \tau_j^x + J_z \tau_i^z \tau_j^z + J_{xz} (\tau_i^x \tau_j^z + \tau_i^z \tau_j^x) + J_y \tau_i^y \tau_j^y] - \sum_i (\hat{n} \cdot \hat{z}_i) h \tau_i^z, \quad (1)$$

where τ^x and τ^z transform as a magnetic dipole moment, τ^y transforms a magnetic octupole moment, $h\hat{n}$ is the external magnetic field that only couples to τ^z linearly, and \hat{z}_i is the local z direction. The crossing J_{xz} term can be eliminated by a rotation around y direction in the pseudospin- τ space with $S_i^x = \tau_i^x \cos \theta - \tau_i^z \sin \theta$, $S_i^z = \tau_i^x \sin \theta + \tau_i^z \cos \theta$, $S_i^y = \tau_i^y$, and the resulting spin model

is of the XYZ form [6, 7],

$$H_{\text{XYZ}} = \sum_{\langle ij \rangle} \tilde{J}_x S_i^x S_j^x + \tilde{J}_y S_i^y S_j^y + \tilde{J}_z S_i^z S_j^z - \sum_i (\hat{n} \cdot \hat{z}_i) h (\cos \theta S_i^z - \sin \theta S_i^x). \quad (2)$$

where \tilde{J} ’s are related to the J ’s via the pseudospin rotation.

INS spectroscopy.—Based on the qualitative and powerful symmetry reasoning, we know that, the model in Eq. (2) without the field has two symmetry-enriched U(1) spin liquids from the perspective of the point group. For a dominant and antiferromagnetic \tilde{J}_z or \tilde{J}_x , a (conventional) dipolar U(1) spin liquid is realized. In contrast, for a dominant and antiferromagnetic \tilde{J}_y , a novel octupolar U(1) spin liquid is realized. It was observed that [7], because the magnetic field or the neutron spin only couples to the transverse components relative to S^y in the octupolar U(1) spin liquid, the inelastic neutron scattering *only* measures the spinon continuum (see Tab. I). An energy threshold of two-spinon gap is needed in order to observe the spinon continuum in the neutron scattering measurement. With this insight, the spinon continuum spectroscopy for both 0-flux and π -flux octupolar U(1) spin liquids were studied [7, 21]. Moreover, due to this selective linear coupling for the dipolar spin components, a regular unpolarized INS may already function as a polarized INS, and a polarized INS can be a unnecessary luxury.

Let us re-examine the dipolar U(1) spin liquid for a dominant and antiferromagnetic \tilde{J}_x . If J_{xz} is weak (or equivalently, θ is small), the behaviors in the neutron scattering measurement in this limit should not be very different from the octupolar U(1) spin liquid. In this regime, the coupling to the neutron spin is primarily via S^z component, and the coupling to S^x is suppressed by $\sin \theta$. The intensity of the S^x - S^x correlation, that includes the electric monopole continuum and the gauge photon, is suppressed by $\sin^2 \theta$. The INS measurement cannot differentiate this dipolar U(1) spin liquid in this regime from the octupolar U(1) spin liquid, regardless of whether further translational symmetry enrichment with

0 or π flux for the spinons is considered. From the perspective of S^z component, the hidden S^x and S^y components actually have the same effect. This can be understood via a rotation of the spin \mathbf{S}_i around the z direction by $\pi/2$ that switches S^x and S^z . For the same reason, thermodynamics like specific heat and magnetic susceptibility cannot differentiate them. Thus, in the above mentioned measurements, the dipolar U(1) spin liquid for a dominant and antiferromagnetic \tilde{J}_x is not very different from the octupolar U(1) spin liquid. We note that, in Ref. 34, a finite linear Zeeman coupling to τ^x is introduced in Eq. (1), and this would be more-or-less equivalent to a finite θ . The presence of this extra coupling could in principle make the experimental behaviors of the above two phases different. Our understanding here is restricted to the regime with the vanishing τ^x Zeeman coupling of Eq. (1).

About $\text{Ce}_2\text{Sn}_2\text{O}_7$.—The recent INS measurement in the powder sample of $\text{Ce}_2\text{Sn}_2\text{O}_7$ actually observed a large spectral weight at $\sim 0.045\text{meV}$, and the spectral weights are gradually suppressed below this peak energy. From the above explanation, this is either consistent with the expectation from the octupolar U(1) spin liquid, or compatible with the dipolar U(1) spin liquid for a dominant and antiferromagnetic \tilde{J}_x and a small θ . On the positive side, the measured excitation continuum for both cases is interpreted as the spinon continuum for the former, and the predominantly spinon continuum for the latter. The further structures in the momentum and/or energy domain of the spinon continuum could differentiate the distinct translation symmetry enrichments, i.e. the 0 or π gauge flux for the spinons [7, 21, 32]. It is tempting to mention that, recent experiments of Ref. 16 on both single-crystal and powder $\text{Ce}_2\text{Sn}_2\text{O}_7$ samples with a different preparation condition concluded with a dominant and antiferromagnetic \tilde{J}_x and an intermediate θ ($\approx 0.19\pi$). The authors suggested that, $\text{Ce}_2\text{Sn}_2\text{O}_7$ is located in the dipolar spin ice regime, but the ground state is ordered in the all-in all-out state. Even if the ground state is not ordered, an intermediate θ ($\approx 0.19\pi$) for a dipolar U(1) spin liquid with a dominant and antiferromagnetic \tilde{J}_x would imply a significant spectral weight due to the electric monopole continuum and the gauge photon at energies below the two-spinon gap. This is clearly different from the inelastic neutron scattering result in Ref. 28. The discrepancy between Ref. 16 and Ref. 28 might arise from the sample preparation [16] and remains to be resolved.

About $\text{Ce}_2\text{Zr}_2\text{O}_7$.—In contrast, the recent INS measurement in the single-crystal sample of $\text{Ce}_2\text{Zr}_2\text{O}_7$ observed enhanced spectral weights at low energies, and the large spectral intensity extends down to zero energy within the experimental reach [23]. Without invoking any other features from the experiments, we think this piece of experimental information alone can be understood from the following three possibilities.

The first possibility is that, the low-energy spectral weights arise from the spinon continuum with a small

Material	\tilde{J}_x	\tilde{J}_y	\tilde{J}_z	θ	Reference
$\text{Ce}_2\text{Sn}_2\text{O}_7$	-	0.05	-	0	Ref. 11
$\text{Ce}_2\text{Sn}_2\text{O}_7$	-	0.048	-	0	Ref. 28
$\text{Ce}_2\text{Sn}_2\text{O}_7$	0.045	-0.001	-0.012	0.19π	Ref. 16
$\text{Ce}_2\text{Zr}_2\text{O}_7$	0.063	0.064	0.011	0	Ref. 15
$\text{Ce}_2\text{Zr}_2\text{O}_7$	0.0385	0.088	0.020	0	Ref. 34
$\text{Ce}_2\text{Zr}_2\text{O}_7$	0.076	-	-	0.12π	Ref. 23

TABLE II. The fitted exchange parameters from different works. In Ref. 11, the authors only obtained $J^\pm/\tilde{J}_y \approx -0.015$ where J^\pm is defined with respect to the y component with $J^\pm = -(\tilde{J}_z + \tilde{J}_x)/4$. Likewise, Ref. 28 obtained $J^\pm/\tilde{J}_y \approx -0.11$. In Ref. 15, an equal fitting is obtained when the values of \tilde{J}_x and \tilde{J}_y are switched. In Ref. 34, a weak but finite linear coupling to τ_i^x was introduced in Eq. (1). In Ref. 23, the authors only obtained $J^\pm/\tilde{J}_x \approx -0.28$ where J^\pm is defined with respect to the x component with $J^\pm = -(\tilde{J}_y + \tilde{J}_z)/4$. The energy unit is given in meV.

gap, and the system is either in the octupolar U(1) spin liquid or in the dipolar U(1) spin liquid for a dominant and antiferromagnetic \tilde{J}_x and a small θ . If the fitted exchange parameters in Tab. II are used, however, the system is clearly not in the regime with weakly gapped spinon excitations. Given the thermodynamic results were used in the fitting, this possibility may be ruled out.

The second possibility is that, the system is in a dipolar U(1) spin liquid with an intermediate θ such that $\sin^2\theta$ and $\cos^2\theta$ are of the same order. This can allow for a dominant and antiferromagnetic \tilde{J}_z , or a dominant and antiferromagnetic \tilde{J}_x , or both. The third possibility is that, the system is in a dipolar U(1) spin liquid with a dominant and antiferromagnetic \tilde{J}_z and a small θ . For the third possibility, the spinon continuum at high energies should be suppressed from $\sin^2\theta$.

For the second and third possibilities, the low-energy spectral weights below the two-spinon gap arise from the electric monopole continuum and the gauge photon. But since the spectral weight of the gauge photon is intrinsically suppressed at low energies, thus the low-energy spectral weights for these two possibilities are mostly given by the electric monopole continuum. This is the fundamental difference of our interpretation from Ref. 23.

Electric monopole continuum.—We here focus on and work out the details for the low-energy electric monopole continuum that is shared by these two possibilities. Unlike the spinons that hop on the diamond lattice formed by the tetrahedral centers of the pyrochlore lattice, the electric monopoles reside on the dual diamond lattice whose bonds penetrate through the hexagonal plaquette centers of the diamond lattice for the spinons. If one sticks to the context of quantum spin ice, the spinons are excitations out of the ice manifold, while the photon and the electric monopoles are excitations that are built

up from the ice manifold and occur in a much lower energy scale than the spinons (see Fig. 1). The spinon has a classical analogue as the classical magnetic monopole, the photon and the electric monopoles are purely of quantum origin and have no classical analogue. Moreover, the electric monopoles experience a π dual U(1) gauge flux on the dual diamond lattice. Thus, the translation symmetry is fractionalized for the electric monopoles.

The Hamiltonian describing the electric monopole hopping on the dual diamond lattice is given as [20],

$$H_m = -t \sum_{\langle \mathbf{r}\mathbf{r}' \rangle} e^{-2\pi i \alpha_{\mathbf{r}\mathbf{r}'}} \Phi_{\mathbf{r}}^\dagger \Phi_{\mathbf{r}'} - \mu \sum_{\mathbf{r}} \Phi_{\mathbf{r}}^\dagger \Phi_{\mathbf{r}}, \quad (3)$$

where $\Phi_{\mathbf{r}}^\dagger$ ($\Phi_{\mathbf{r}}$) represents the monopole creating (annihilation) operator, and t and μ denote the hopping and chemical potential of the monopoles, respectively. The position \mathbf{r} spans the dual diamond lattice, comprising the A and B sublattices. The dual gauge link $\alpha_{\mathbf{r}\mathbf{r}'}$, capturing the π dual U(1) gauge flux due to the effective spin-1/2 nature of the moment, is chosen as $2\pi \alpha_{\mathbf{r}, \mathbf{r}+\mathbf{e}_\mu} = \xi_\mu \mathbf{Q} \cdot \mathbf{r}$, where $\mathbf{r} \in$ sublattice A and $\mathbf{t} + \mathbf{e}_\mu \in$ sublattice B. Here, \mathbf{e}_μ ($\mu = 0, 1, 2, 3$) represents the nearest-neighbor vectors connecting the two sublattices (see Appendix A). To fix the gauge, we choose $(\xi_0, \xi_1, \xi_2, \xi_3) = (0, 1, 1, 0)$ and $\mathbf{Q} = 2\pi(1, 0, 0)$. Diagonalizing Eq. (3), we obtain four monopole bands

$$\Omega_\eta^\zeta = t\zeta \left[4 + 2\eta(3 + C_x C_y - C_x C_z + C_y C_z)^{\frac{1}{2}} \right]^{\frac{1}{2}} - \mu, \quad (4)$$

where $C_\mu = \cos q_\mu$ ($\mu = x, y, z$), $\zeta = \pm$, and $\eta = \pm$. Due to the background π dual U(1) gauge flux, the monopole bands are then defined over a “magnetic” Brillouin zone that is half of the crystal Brillouin zone.

As previously discussed, the two-monopole continuum predominantly contributes to the INS signal in the low-energy regime. This signal is expected to correlate with the two-monopole continuum density of states, $\rho^{(2)}(\mathbf{k}, \omega)$, regardless of a specific form factor. With the momentum and energy conservation, the density of states is given by

$$\rho^{(2)}(\mathbf{k}, \omega) \propto \sum_{\mathbf{q}, \zeta_1, \zeta_2, \eta_1, \eta_2} \delta[\omega - \Omega_{\eta_1}^{\zeta_1}(\mathbf{q}) - \Omega_{\eta_2}^{\zeta_2}(\mathbf{k} - \mathbf{q} - \mathbf{Q})], \quad (5)$$

where \mathbf{k} and ω represent the total momentum and energy of the two monopoles, respectively. The offset \mathbf{Q} originates from the spatial periodicity of the gauge potential. Due to the periodicity of $\rho^{(2)}(\mathbf{k}, \omega)$, however, this offset does not influence the spectrum. Specifically, at the Γ point, the momenta of the two monopoles cancel out. Given the symmetry of the monopole bands under reflection with respect to the $\Omega = -\mu$ plane and inversion $\mathbf{q} \rightarrow -\mathbf{q}$, the continuum spectrum features a flat band near $\omega = -2\mu$. Consequently, a δ -function-like peak is observable in the two-monopole continuum density of states at the Γ point (see Fig. 2(a)).

In the following, we set t as the energy unit and $\mu = -3t$ to prevent the electric monopole condensation [22]. We normalize $\rho^{(2)}(\mathbf{k}, \omega)$ by its maximal value.

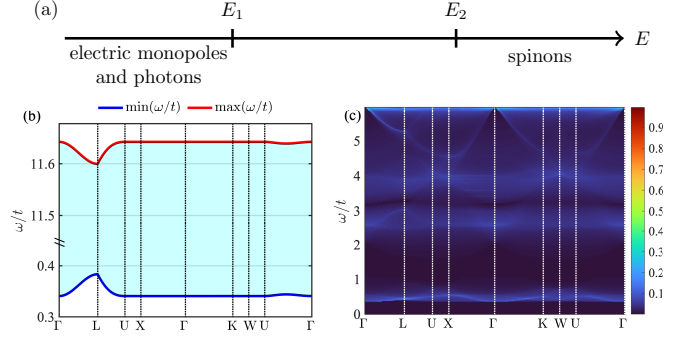


FIG. 1. (a) The schematic energy level diagram of electric monopoles, photons and spinons. (b) and (c) are the structure of the two-monopole continuum along the high-symmetry line in the magnetic Brillouin zone: Γ -L-U-X- Γ -K-W-U- Γ . The chemical potential is set to $\mu = -3t$. (b) Lower edge (blue curve) and upper edge (red curve) of the two-monopole continuum. (c) Normalized two-monopole continuum density of states, $\rho^{(2)}(\mathbf{k}, \omega)$, spanning from the lower edge to half of the continuum width. The brightness indicates the strength of the density of states.

Given the energy overlapping between the four monopole bands, the two-monopole continuum spans continuously from a lower to an upper edge, with a spectrum width of $8\sqrt{2}t$. Fig. 1 illustrates the edge shapes along the high-symmetry line (for the definition of high-symmetry points, see Appendix A). Additionally, Fig. 2(b) and (c) display the density of states at two representative high-symmetry points.

Discussion—We discuss the fitted parameters for $\text{Ce}_2\text{Zr}_2\text{O}_7$ from Tab. II. For the parameters in Ref. 15, since $\theta \sim 0$, as we have discussed, one cannot distinguish the octupolar U(1) spin liquid with a dominant \tilde{J}_y and the dipolar U(1) spin liquid with a dominant \tilde{J}_x . For both phases, the INS measurement can only probe the spinon continuum at high energies. Thus, this set of parameters is inconsistent with the neutron results in Ref. 23.

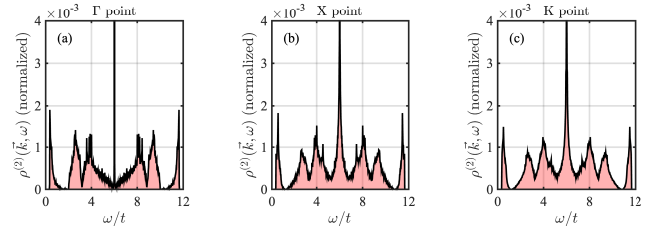


FIG. 2. Normalized two-monopole continuum density of states, $\rho^{(2)}(\mathbf{k}, \omega)$, at three high-symmetry points: Γ , X, and K. (a) At the Γ point, the density of states features a δ -function-like peak at the spectrum center due to the energy-momentum matching of two monopoles. (b) At the X point and (c) K point, the density of states exhibits a similar shape, albeit with weaker peaks at the spectrum center compared to the Γ point.

For the parameters in Ref. 34, even though an additional linear Zeeman coupling to τ^x in Eq. (1) is introduced, the parameters for the exchange simply favor an octupolar U(1) spin liquid, for which the INS measurement can only probe the spinon continuum at high energies. For the parameters in Ref. 23, the system is in the dipolar U(1) spin liquid. As $\theta = 0.12\pi$ is in the intermediate regime, all excitations should show up in the INS measurement, where the low energy part should be mostly contributed by the electric monopole continuum. This is qualitatively compatible with the expectation. On a bit more quantitative side, if one uses the fitted couplings of Ref. 23, the energy scale for the electric monopoles and photons is approximately set by the ring exchange $J_{\text{ring}} = 12(J^\pm)^3/J_x^2 \sim 0.02\text{meV}$, while the spinon energy scale is set by \tilde{J}_x and renormalized by J_\pm . Thus, we expect that, probably below about 0.05meV , the INS spectrum is contributed by photons and electric monopole continuum, and mostly by the electric monopole continuum. The detailed structures of the continuum require a high energy resolution to resolve. Otherwise, the continuous excitations are simply crowded up at low energies (see Appendix B). Roughly above about 0.05meV , the INS spectrum is contributed by spinon continuum combined with other low-energy excitations (photon, electric monopoles).

The above analysis, however, has one caveat. If a linear Zeeman coupling to τ^x is introduced in Eq. (1) as Ref. 34, a dipolar U(1) spin liquid with a dominant \tilde{J}_z or \tilde{J}_x with $\theta \sim 0$ could not be simply ruled out. About the calculation, the caveat is that we treat the electric monopole at the quadratic level and the dual U(1) gauge field as a static background. According to Ref. 35, the threshold spinon-pair production could be abruptly enhanced due to the coupling to the dynamic gauge field. Similar features would occur for the threshold electric monopole production, giving rise to a large density of states at low energies than our current treatment.

It is interesting to vision the effect of a weak magnetic field for a dipolar U(1) spin liquid if $\text{Ce}_2\text{Zr}_2\text{O}_7$ realizes it. A weak magnetic field polarizes the effective Ising component and modifies the dual gauge flux away from π for the electric monopoles. This has two consequences. First, the electric monopoles develop Hofstadter-like band structures, and the continuum will have more modulations in the energy-momentum domain [22]. The abundance of the electric monopole continuum at low energy should also be accessible through the μSR measurement at zero/weak field [12]. Second, the electric monopole band will have a nontrivial Berry curvature distribution, generating the monopole thermal Hall effect at low temperatures [36]. If $\text{Ce}_2\text{Sn}_2\text{O}_7$ realizes the octupolar U(1) spin liquid, no such effect is expected. In general, the very low energy scale of the monopole physics for the rare-earth magnets makes the detailed measurement difficult, and this calls for the candidate d electron systems with higher energy scales and the numerical calculations [37]. So far, one candidate d

electron magnets for the DO doublets is $\text{Cd}_2\text{Os}_2\text{O}_7$, but it is ordered [38].

In summary, we propose the low-energy continuous excitations in the recent inelastic neutron scattering measurements in $\text{Ce}_2\text{Zr}_2\text{O}_7$ is mostly contributed from the electric monopole continuum. Although the gapless gauge phonon is present, its intensity should be suppressed by at low energies. While the fitted parameters in Ref. 23 are qualitatively compatible with our expectation from the INS scattering measurement, our interpretation is different from Ref. 23. We further demonstrate the structures of the electric monopole continuum in the energy and momentum domains.

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Appendix A: Lattice structure

The physical spin operators live in a Pyrochlore lattice, which is formed by corner-shared tetrahedra. A diamond lattice can be constructed by connecting the centers of the tetrahedra, on which a U(1) lattice gauge theory is defined. A further duality transformation yields the theory of monopoles, which live in the dual lattice of the diamond lattice: the dual diamond lattice. It consists of two sublattices. The relative vector between these two sublattices is $\delta = \frac{1}{2}(1, 1, 1)$. Each sublattice is a face-centered cubic (fcc) lattice. Each site of the dual diamond lattice has four nearest neighbors, whose relative vectors are $\mathbf{e}_0 = \frac{1}{4}(1, 1, 1)$, $\mathbf{e}_1 = \frac{1}{4}(1, -1, -1)$, $\mathbf{e}_2 = \frac{1}{4}(-1, 1, -1)$, and $\mathbf{e}_3 = \frac{1}{4}(-1, -1, 1)$. The lattice vectors of the underlying Bravais lattice are chosen to be $\mathbf{a}_1 = \frac{1}{2}(0, 1, 1)$, $\mathbf{a}_2 = \frac{1}{2}(1, 0, 1)$, and $\mathbf{a}_3 = \frac{1}{2}(1, 1, 0)$. Thus, the three lattice vectors of the reciprocal lattice are defined as $\mathbf{b}_1 = 2\pi(-1, 1, 1)$, $\mathbf{b}_2 = 2\pi(1, -1, 1)$, and $\mathbf{b}_3 = 2\pi(1, 1, -1)$. As described in the main text, monopole bands in the π -flux background has an enhanced periodicity, which leads to a smaller Brillouin zone spanned by $\mathbf{b}'_i = \mathbf{b}_i/2$. To reveal the spectrum, it is nature to consider the high-symmetry points defined in Tab. A1 in the smaller Brillouin zone.

Appendix B: More details on the spectrum

The peak structure of the two-monopole continuum DoS appears to be insensitive to the total momentum. To validate this observation, we compute the momentum-integrated DoS, given by $\int d\mathbf{k} \rho^{(2)}(\mathbf{k}, \omega)$. In Figure A1(b), the momentum-integrated DoS exhibits

high-symmetry point	Cartesian coordinates
Γ	$\pi(0, 0, 0)$
X	$\pi(0, 0, 1)$
L	$\pi(1/2, 1/2, 1/2)$
W	$\pi(0, 1/2, 1)$
U	$\pi(1/4, 1/4, 1)$
K	$\pi(0, 3/4, 3/4)$

TABLE A1. The definition of high-symmetry points of the magnetic Brillouin zone.

seven peaks, with their positions centered around the highest peak located at $\omega = -2\mu = 6$. This peak structure can be elucidated by examining the DoS of the four monopole bands. Each monopole band contributes a peak to the one-monopole DoS, and their collective effect manifests as the observed peak structure in the two-monopole continuum DoS.

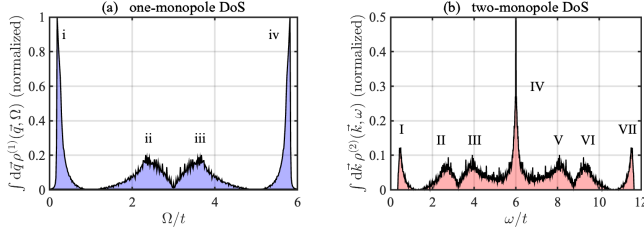


FIG. A1. Illustration of density of states (DoS) for the parameter setting $\mu = -3t$. (a) Normalized momentum-integrated one-monopole DoS, $\int d\mathbf{q} \rho^{(1)}(\mathbf{q}, \Omega)$, exhibiting four peaks labeled i, ii, iii, and iv, corresponding to the four monopole bands. (b) Normalized momentum-integrated two-monopole DoS, $\int d\mathbf{k} \rho^{(2)}(\mathbf{k}, \omega)$, showing seven peaks labeled I, II, III, IV, V, VI, and VII. The enhancement of a sharp peak in the middle is prominently visible.

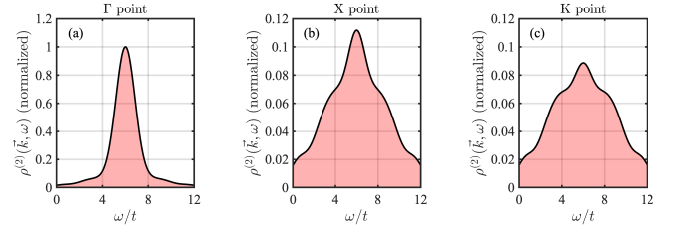


FIG. A2. Normalized two-monopole continuum density of states, $\rho^{(2)}(\mathbf{k}, \omega)$, considering a finite energy resolution at three high-symmetry points: Γ , X, and K.

While the spectrum presents fine structures, these might not be readily discernible in INS experiments. Being excitations within the ice manifold, monopoles operate on a relatively small energy scale, making their spectrum width potentially compatible with the energy resolution of INS. Consequently, a broadened peak is anticipated rather than a sharp one. To account for finite energy resolution, we simulate it using a Gaussian function with a standard deviation equal to 10% of the spectrum width. Figure A2 displays the continuum density of states at three high-symmetry points: Γ , X, and K.

To detect the electric monopole continuum, one potential strategy could involve synthesizing both an octupolar U(1) spin liquid material and a dipolar U(1) spin liquid material with identical parameters. In the case of the octupolar U(1) spin liquid, INS measures only the gapped spinon continuum. By subtracting the spinon continuum signal of the octupolar U(1) spin liquid from the INS signal of the dipolar U(1) spin liquid, the remaining signal would predominantly reflect the monopole continuum and gauge photon. This differential approach may provide a robust method to isolate and detect the electric monopole continuum signal experimentally.

- [1] J. G. Rau and M. J. Gingras, Frustrated quantum rare-earth pyrochlores, *Annual Review of Condensed Matter Physics* **10**, 357–386 (2019).
- [2] L. Savary and L. Balents, Quantum spin liquids: a review, *Reports on Progress in Physics* **80**, 016502 (2016).
- [3] M. Hermele, M. P. A. Fisher, and L. Balents, Pyrochlore photons: The U(1) spin liquid in a $S = \frac{1}{2}$ three-dimensional frustrated magnet, *Phys. Rev. B* **69**, 064404 (2004).
- [4] H. R. Molavian, M. J. P. Gingras, and B. Canals, Dynamically Induced Frustration as a Route to a Quantum Spin Ice State in $\text{Tb}_2\text{Ti}_2\text{O}_7$ via Virtual Crystal Field Excitations and Quantum Many-Body Effects, *Phys. Rev. Lett.* **98**, 157204 (2007).
- [5] K. A. Ross, L. Savary, B. D. Gaulin, and L. Balents, Quantum excitations in quantum spin ice, *Phys. Rev. X* **1**, 021002 (2011).
- [6] Y.-P. Huang, G. Chen, and M. Hermele, Quantum Spin Ices and Topological Phases from Dipolar-Octupolar Doublets on the Pyrochlore Lattice, *Phys. Rev. Lett.* **112**, 167203 (2014).
- [7] Y.-D. Li and G. Chen, Symmetry enriched U(1) topological orders for dipole-octupole doublets on a pyrochlore lattice, *Phys. Rev. B* **95**, 041106 (2017).
- [8] R. Sibille, E. Lhotel, V. Pomjakushin, C. Baines, T. Fennell, and M. Kenzelmann, Candidate Quantum Spin Liquid in the Ce^{3+} Pyrochlore Stannate $\text{Ce}_2\text{Sn}_2\text{O}_7$, *Phys. Rev. Lett.* **115**, 097202 (2015).
- [9] J. Gaudet, E. M. Smith, J. Dudemaine, J. Beare, C. R. C. Buhariwalla, N. P. Butch, M. B. Stone, A. I. Kolesnikov, G. Xu, D. R. Yahne, K. A. Ross, C. A. Marjerrison, J. D. Garrett, G. M. Luke, A. D. Bianchi, and B. D. Gaulin, Quantum Spin Ice Dynamics in the Dipole-Octupole Pyrochlore Magnet $\text{Ce}_2\text{Zr}_2\text{O}_7$, *Phys. Rev. Lett.* **122**, 187201 (2019).

- (2019).
- [10] B. Gao, T. Chen, D. W. Tam, C.-L. Huang, K. Sasmal, D. T. Adroja, F. Ye, H. Cao, G. Sala, M. B. Stone, C. Baines, J. A. T. Verezhak, H. Hu, J.-H. Chung, X. Xu, S.-W. Cheong, M. Nallaiyan, S. Spagna, M. B. Maple, A. H. Nevidomskyy, E. Morosan, G. Chen, and P. Dai, Experimental signatures of a three-dimensional quantum spin liquid in effective spin-1/2 $\text{Ce}_2\text{Zr}_2\text{O}_7$ pyrochlore, *Nature Physics* **15**, 1052 (2019).
 - [11] R. Sibille, N. Gauthier, E. Lhotel, V. Porée, V. Pomjakushin, R. A. Ewings, T. G. Perring, J. Ollivier, A. Wildes, C. Ritter, T. C. Hansen, D. A. Keen, G. J. Nilsen, L. Keller, S. Petit, and T. Fennell, A quantum liquid of magnetic octupoles on the pyrochlore lattice, *Nature Physics* **16**, 546 (2020).
 - [12] J. Beare, E. M. Smith, J. Dudemaine, R. Schäfer, M. R. Rutherford, S. Sharma, A. Fitterman, C. A. Marjerrison, T. J. Williams, A. A. Aczel, S. R. Dunsiger, A. D. Bianchi, B. D. Gaulin, and G. M. Luke, μ SR study of the dipole-octupole quantum spin ice candidate $\text{Ce}_2\text{Zr}_2\text{O}_7$, *Phys. Rev. B* **108**, 174411 (2023).
 - [13] E. M. Smith, J. Dudemaine, B. Placke, R. Schäfer, D. R. Yahne, T. DeLazzer, A. Fitterman, J. Beare, J. Gaudet, C. R. C. Buhariwalla, A. Podlesnyak, G. Xu, J. P. Clancy, R. Movshovich, G. M. Luke, K. A. Ross, R. Moessner, O. Benton, A. D. Bianchi, and B. D. Gaulin, Quantum spin ice response to a magnetic field in the dipole-octupole pyrochlore $\text{Ce}_2\text{Zr}_2\text{O}_7$, *Phys. Rev. B* **108**, 054438 (2023).
 - [14] V. Porée, A. Bhardwaj, E. Lhotel, S. Petit, N. Gauthier, H. Yan, V. Pomjakushin, J. Ollivier, J. A. Quilliam, A. H. Nevidomskyy, H. J. Changlani, and R. Sibille, Dipolar-octupolar correlations and hierarchy of exchange interactions in $\text{Ce}_2\text{Hf}_2\text{O}_7$, *arXiv e-prints*, [arXiv:2305.08261](https://arxiv.org/abs/2305.08261) (2023), [arXiv:2305.08261 \[cond-mat.str-el\]](https://arxiv.org/abs/2305.08261).
 - [15] E. M. Smith, O. Benton, D. R. Yahne, B. Placke, R. Schäfer, J. Gaudet, J. Dudemaine, A. Fitterman, J. Beare, A. R. Wildes, S. Bhattacharya, T. DeLazzer, C. R. C. Buhariwalla, N. P. Butch, R. Movshovich, J. D. Garrett, C. A. Marjerrison, J. P. Clancy, E. Kermarrec, G. M. Luke, A. D. Bianchi, K. A. Ross, and B. D. Gaulin, Case for a $U(1)_\pi$ Quantum Spin Liquid Ground State in the Dipole-Octupole Pyrochlore $\text{Ce}_2\text{Zr}_2\text{O}_7$, *Physical Review X* **12**, 021015 (2022).
 - [16] D. R. Yahne, B. Placke, R. Schäfer, O. Benton, R. Moessner, M. Powell, J. W. Kolis, C. M. Pasco, A. F. May, M. D. Frontzek, E. M. Smith, B. D. Gaulin, S. Calder, and K. A. Ross, Dipolar Spin Ice Regime Proximate to an All-In-All-Out Néel Ground State in the Dipolar-Octupolar Pyrochlore $\text{Ce}_2\text{Sn}_2\text{O}_7$, *Physical Review X* **14**, 011005 (2024).
 - [17] A. Bhardwaj, V. Porée, H. Yan, N. Gauthier, E. Lhotel, S. Petit, J. A. Quilliam, A. H. Nevidomskyy, R. Sibille, and H. J. Changlani, Thermodynamics of the dipole-octupole pyrochlore magnet $\text{Ce}_2\text{Hf}_2\text{O}_7$ in applied magnetic fields, *arXiv e-prints*, [arXiv:2402.08723](https://arxiv.org/abs/2402.08723) (2024), [arXiv:2402.08723 \[cond-mat.str-el\]](https://arxiv.org/abs/2402.08723).
 - [18] L. Savary and L. Balents, Coulombic quantum liquids in spin-1/2 pyrochlores, *Phys. Rev. Lett.* **108**, 037202 (2012).
 - [19] O. Benton, O. Sikora, and N. Shannon, Seeing the light: Experimental signatures of emergent electromagnetism in a quantum spin ice, *Phys. Rev. B* **86**, 075154 (2012).
 - [20] G. Chen, Dirac's "magnetic monopoles" in pyrochlore ice $U(1)$ spin liquids: Spectrum and classification, *Phys. Rev. B* **96**, 195127 (2017).
 - [21] X.-P. Yao, Y.-D. Li, and G. Chen, Pyrochlore $U(1)$ spin liquid of mixed-symmetry enrichments in magnetic fields, *Phys. Rev. Res.* **2**, 013334 (2020).
 - [22] G. Chen, "Magnetic monopole" condensation of the pyrochlore ice $U(1)$ quantum spin liquid: Application to $\text{Pr}_2\text{Ir}_2\text{O}_7$ and $\text{Yb}_2\text{Ti}_2\text{O}_7$, *Phys. Rev. B* **94**, 205107 (2016).
 - [23] B. Gao, F. Desrochers, D. W. Tam, P. Steffens, A. Hiess, Y. Su, S.-W. Cheong, Y. B. Kim, and P. Dai, Emergent photons and fractionalized excitations in a quantum spin liquid (2024), [arXiv:2404.04207 \[cond-mat.str-el\]](https://arxiv.org/abs/2404.04207).
 - [24] J. Xu, O. Benton, A. T. M. N. Islam, T. Guidi, G. Ehlers, and B. Lake, Order out of a Coulomb Phase and Higgs Transition: Frustrated Transverse Interactions of $\text{Nd}_2\text{Zr}_2\text{O}_7$, *Phys. Rev. Lett.* **124**, 097203 (2020).
 - [25] J. Xu, O. Benton, V. K. Anand, A. T. M. N. Islam, T. Guidi, G. Ehlers, E. Feng, Y. Su, A. Sakai, P. Gegenwart, and B. Lake, Anisotropic exchange Hamiltonian, magnetic phase diagram, and domain inversion of $\text{Nd}_2\text{Zr}_2\text{O}_7$, *Phys. Rev. B* **99**, 144420 (2019).
 - [26] A. Bertin, P. Dalmas de Réotier, B. Fåk, C. Marin, A. Yaouanc, A. Forget, D. Sheptyakov, B. Frick, C. Ritter, A. Amato, C. Baines, and P. J. C. King, $\text{Nd}_2\text{Sn}_2\text{O}_7$: An all-in-all-out pyrochlore magnet with no divergence-free field and anomalously slow paramagnetic spin dynamics, *Phys. Rev. B* **92**, 144423 (2015).
 - [27] G. Chen, Coulombic antiferromagnet in spin-1/2 pyrochlores with dipole-octupole doublets, *Phys. Rev. Res.* **5**, L032027 (2023).
 - [28] V. Porée, H. Yan, F. Desrochers, S. Petit, E. Lhotel, M. Appel, J. Ollivier, Y. B. Kim, A. H. Nevidomskyy, and R. Sibille, Fractional matter coupled to the emergent gauge field in a quantum spin ice (2024), [arXiv:2304.05452 \[cond-mat.str-el\]](https://arxiv.org/abs/2304.05452).
 - [29] A. S. Patri, M. Hosoi, and Y. B. Kim, Distinguishing dipolar and octupolar quantum spin ices using contrasting magnetostriction signatures, *Physical Review Research* **2**, 023253 (2020).
 - [30] F. Desrochers, L. E. Chern, and Y. B. Kim, Competing $U(1)$ and Z_2 dipolar-octupolar quantum spin liquids on the pyrochlore lattice: Application to $\text{Ce}_2\text{Zr}_2\text{O}_7$, *Phys. Rev. B* **105**, 035149 (2022).
 - [31] F. Desrochers, L. E. Chern, and Y. B. Kim, Symmetry fractionalization in the gauge mean-field theory of quantum spin ice, *Phys. Rev. B* **107**, 064404 (2023).
 - [32] F. Desrochers and Y. B. Kim, Spectroscopic Signatures of Fractionalization in Octupolar Quantum Spin Ice, *Phys. Rev. Lett.* **132**, 066502 (2024).
 - [33] F. Desrochers and Y. B. Kim, Finite Temperature Dynamics in 0-Flux and π -Flux Quantum Spin Ice: Self-Consistent Exclusive Boson Approach, *arXiv e-prints* [10.48550/arXiv.2401.09551](https://arxiv.org/abs/2401.09551) (2024), [arXiv:2401.09551](https://arxiv.org/abs/2401.09551).
 - [34] A. Bhardwaj, S. Zhang, H. Yan, R. Moessner, A. H. Nevidomskyy, and H. J. Changlani, Sleuthing out exotic quantum spin liquidity in the pyrochlore magnet $\text{Ce}_2\text{Zr}_2\text{O}_7$, *npj Quantum Materials* **7**, [10.1038/s41535-022-00458-2](https://doi.org/10.1038/s41535-022-00458-2) (2022).
 - [35] S. C. Morampudi, F. Wilczek, and C. R. Laumann, Spectroscopy of Spinons in Coulomb Quantum Spin Liquids, *Phys. Rev. Lett.* **124**, 097204 (2020).
 - [36] X.-T. Zhang, Y. H. Gao, C. Liu, and G. Chen, Topological thermal Hall effect of magnetic monopoles in the

- pyrochlore U(1) spin liquid, [Phys. Rev. Res. **2**, 013066 \(2020\)](#).
- [37] C.-J. Huang, C. Liu, Z. Meng, Y. Yu, Y. Deng, and G. Chen, Extended coulomb liquid of paired hardcore boson model on a pyrochlore lattice, [Phys. Rev. Res. **2**, 042022 \(2020\)](#).
- [38] H. Shinaoka, T. Miyake, and S. Ishibashi, Noncollinear Magnetism and Spin-Orbit Coupling in $5d$ Pyrochlore Oxide $\text{Cd}_2\text{Os}_2\text{O}_7$, [Phys. Rev. Lett. **108**, 247204 \(2012\)](#).